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Indication of a Pulsar Wind Nebula in the hard X-ray emission from SN 1987A

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ABSTRACT

Since the day of its explosion, SN 1987A (SN87A) was closely monitored with the aim to study its evolution and to detect its central compact relic. The detection of neutrinos from the supernova strongly supports the formation of a neutron star (NS). However, the constant and fruitless search for this object has led to different hypotheses on its nature. Up to date, the detection in the ALMA data of a feature somehow compatible with the emission arising from a proto Pulsar Wind Nebula (PWN) is the only hint of the existence of such elusive compact object. Here we tackle this 33-years old issue by analyzing archived observations of SN87A performed by *Chandra* and *NuSTAR* in different years. We firmly detect nonthermal emission in the 10 – 20 keV energy band, due to synchrotron radiation. The possible physical mechanism powering such emission is twofold: diffusive shock acceleration (DSA) or emission arising from an absorbed PWN. By relating a state-of-the-art magneto-hydrodynamic simulation of SN87A to the actual data, we reconstruct the absorption pattern of the PWN embedded in the remnant and surrounded by cold ejecta. We found that, even though the DSA scenario cannot be firmly excluded, the most likely scenario that well explains the data is the PWN emission.

Keywords: X-rays: general - supernovae: individual (SN 1987A)

1. INTRODUCTION

SN 1987A (SN87A) in the Large Magellanic Cloud (LMC) was a hydrogen-rich core-collapse supernova (SN) discovered on 1987 February 23 (West et al. 1987). It occurred approximately 51.4 kpc from Earth (Panzagia 1999) and its dynamical evolution is strictly related to the very inhomogeneous circumstellar medium (CSM), composed by a dense ring-like structure within a diffuse HII region (Sugerman et al. 2005). SN87A is the first naked-eye SN exploded since telescopes exist and its evolution has been deeply monitored in various wavelengths (McCray 1993; McCray & Fransson 2016). In particular, the X-ray band is ideal to investigate the interaction of the shock front with the CSM and the emission of the expected central compact leftover of the supernova explosion.

Despite the unique consideration granted with deep and continuous observations, and the neutrinos detection (Bionta et al. 1987) strongly indicating the formation of a neutron star (Vissani 2015), the elusive compact object of SN87A is still undetected. The most likely explanation for this non-detection is ascribable to the absorption due to ejecta, i.e. the dense and cold material ejected by the supernova (Fransson & Chevalier 1987): because of the young age of SN87A, the ejecta are still very dense and the reverse shock generated in the outer shells of the supernova remnant (SNR) has not heated the inner ejecta yet. Thus, photo-electric absorption from this metal-rich material can hide the X-ray emission of a hypothetical compact object. During the last years, many works (Orlando et al. 2015; Alp et al. 2018; Esposito et al. 2018; Page et al. 2020) investigated the upper limit on the luminosity of the putative compact leftover in various wavelengths, considering the case of a neutron star (NS) emitting thermal (black-body) radiation and obscured by the cold ejecta and/or dust. However, the lack of strong constraints on

the absorption pattern in the internal area of SN87A prevented either to further constrain the luminosity of the putative NS or to make predictions about its future detectability.

On the other side, the X-ray emission from a young NS may include a significant nonthermal component: the synchrotron radiation arising from the pulsar wind nebula (PWN) associated with the rotating NS. Recently, *ALMA* images showed a *blob* structure whose emission is somehow compatible with the radio emission of a PWN (Cigan et al. 2019). However, the authors themselves warned that this blob could be associated with other physical processes, e.g. heating due to ^{44}Ti decay. It is natural, then, to look for the high-energy counterpart of the synchrotron radiation in the X-ray band.

In this letter, we report on the analysis of observations of SN87A performed between 2012 and 2014 by *Chandra* and *NuSTAR*. We also take advantage of the state-of-the-art MHD simulation from Orlando et al. (2020) (hereafter Or20) to reconstruct the absorption pattern within SN87A, and link it to the observed spectra. We provide a single model describing the emission in these years from 0.5 to 20 keV, isolating the synchrotron radiation arising from the remnant. We then discuss the possible origin of such emission and show how the presence of a PWN appears to be the most likely scenario.

2. X-RAY DATA ANALYSIS

We used data collected in 2012, 2013 and 2014 with *Chandra* ACIS-S and *NuSTAR* CZT (FPMA and FPMB). We reprocessed *Chandra* and *NuSTAR* data with the standard pipelines available within CIAO v4.12.2 and NuSTARDAS, respectively. For details on the observations and the data reduction see Appendix A.

We extracted spectra from a circular region centered at $\alpha = 5^{\text{h}}35^{\text{m}}28^{\text{s}}$ and $\delta = -69^{\circ}16'10''$ with a radius of $2''$ and $43''$ for *Chandra* and *NuSTAR* data, respectively. These regions enclose all the remnant and were chosen taking into account the PSFs of the different telescopes (see Fig. 1).

We simultaneously analyzed *Chandra* and *NuSTAR* spectra for each year considered by adopting a model composed by a galactic absorption component (*TBabs* model in XSPEC), two optically thin isothermal components in non-equilibrium of ionization (*vnei* model) and a constant factor which takes into account cross-calibration between different detectors. The column density n_{H} is fixed to $2.35 \times 10^{21} \text{ cm}^{-2}$ (Park et al. 2006). We found cross-calibration factors $< 2\%$ within different *Chandra* data sets of the same year, and $< 8\%$ between *Chandra* and *NuSTAR*, compatible with the character-

istic corrections between *NuSTAR* detectors and other telescopes. Temperatures, emission measure, ionization age and abundances of O, Ne, Mg, Si and S were left free to vary in the fitting process. All other abundances were kept fixed to those found by Zhekov et al. (2009). We report no significant variations in the chemical abundances in the time range considered (2012 to 2014).

In agreement with previous works (Orlando et al. 2015; Miceli et al. 2019, Or20), we found that the soft (0.5 – 8 keV) X-ray emission consists of thermal X-rays originating in the shocked CSM. Our best-fit values of temperatures, kT and ionization parameters, τ , are compatible with previous measures (Zhekov et al. 2009). However, we found strong residuals in all *NuSTAR* spectra at energies > 10 keV, clearly showing that an additional component must be added to the model to properly describe the hard X-ray emission (see Reynolds et al. 2015). The best-fit model and the residuals are shown in the left panel of Fig. 2. This additional component cannot be associated with thermal emission from the shocked plasma, since otherwise unrealistically high temperatures would be necessary ($kT \sim 20$ keV), one order of magnitude higher than the maximum electron temperature predicted for SN 1987A (Orlando et al. 2015) and never observed in SNRs.

If the emission is unlikely to be thermal, it is plausible to suppose that it is nonthermal. This nonthermal emission may arise from a possible compact object, most likely a NS embedded in its PWN (hereafter PWN87A), which emits synchrotron radiation. The main issue in tackling this scenario is to isolate the radiation coming from this object.

We can estimate the absorbing power of cold ejecta surrounding the PWN by taking into account the abundance and density pattern along the line of sight provided by the MHD model by Or20 (see Appendix B) in the considered years. To include this information in the spectral analysis, we added an absorbed power-law modeled with *vphabs* within XSPEC. Since its parameters were derived from the model they are not free to vary in the fit. With this additional component, we obtained a very good description of the observed spectra over the whole 0.5 – 20 keV energy band (right panel of Fig. 2). The best-fit value of photon index and normalization of the absorbed power-law do not change significantly between the three years considered. We then fitted simultaneously the 2012, 2013 and 2014 spectra in order to decrease the uncertainties in the best-fit parameters of the power-law component. Temperatures, ionization parameters and normalization of thermal emission were free to change in time, while normalization and photon index of nonthermal emission (and chemical abun-

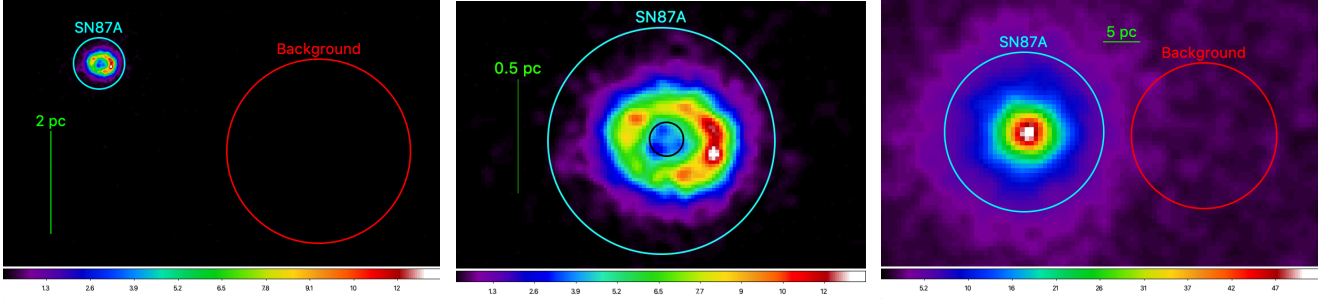


Figure 1. *Chandra* and *NuSTAR* count images of SN87A. All the images are smoothed with a 1.5σ -gaussian. *Left panel.* *Chandra* image in the 0.1-8 keV. The cyan and the red circular regions mark the $2''$ region used to extract the source and the corresponding background spectra for the *Chandra* data, respectively. *Central panel.* Closeup of the left panel. The $0.3''$ radius black ring marks the faint central region of SN87A. *Right panel.* *NuSTAR* image in the 3-30 keV band. The cyan and red regions identify the $43''$ region used to extract the source spectra and the corresponding background for the *NuSTAR* data, respectively. A navigable 3D graphic showing the PWN position in the remnant interior and the comparison with the observations is available at the link <https://skfb.ly/6XZIU>.

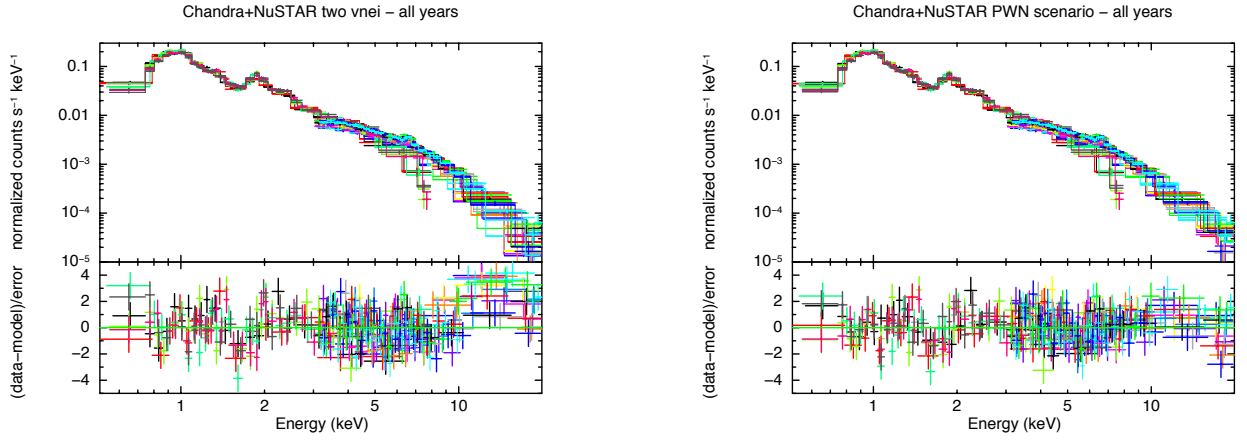


Figure 2. Spectra extracted in the 0.5–20 keV energy band from *Chandra* and *NuSTAR* in various years with the corresponding best-fit model and residuals. A different color is associated with each of the twenty data set. On the left, the best-fit model is composed by two thermal components. On the right, the best-fit model also takes into account the emission coming from a heavily absorbed PWN. The spectra have been rebinned for presentation purposes.

dances) were left free to vary in the fitting procedure, but forced to be the same over the three years.

The resulting PWN best-fit photon index is $\Gamma = 2.5^{+0.3}_{-0.4}$ and the X-ray luminosity in the 1–10 keV band is $L_{1-10} = (2.6 \pm 1.4) \times 10^{35}$ erg/s (Table 1). Our best fit values are compatible with typical values found for PWNe (Sect. 3).

We found an increase in flux in the 8–20 keV band between 2012 and 2014 (as reported by Reynolds et al. 2015). We point out that the emission from the PWN is consistent with being constant, and that the increase in the hard X-ray flux is related to the steadily growing thermal emission arising from the interaction of the remnant with the HII region.

Because of the high ejecta absorption, the PWN flux is strongly suppressed below 8 keV. In particular, we

found that the PWN emission is much dimmer ($\sim 4\%$) than the thermal emission observed by *Chandra* in the faint central region of SN87A (identified by a circular region with radius $R = 0.3''$, well inside the bright X-ray emitting ring, see central panel in Fig. 1).

As an alternative scenario, we also considered the case in which the hard X-ray emission is associated with synchrotron radiation due to DSA occurring in the outer layers of SN87A. We then removed the cold ejecta absorption. The best-fit photon index is $\Gamma_{\text{DSA}} = 2.0 \pm 0.4$, indicating lower steepness than the PWN one but still consistent with it. Analogously to what we obtained for the PWN scenario, we found that photon index and normalization of the nonthermal emission do not change significantly from 2012 to 2014. We then repeated the simultaneous fit of the data from all years obtaining a

good description of the observed spectra ($\chi^2_{\text{DSA}} = 1457$ with 1221 d.o.f. to be compared with $\chi^2_{\text{PWN}} = 1442$ with 1221 d.o.f. of the PWN scenario).

Since from a statistical point of view an improvement of only 1% is not enough to favour the PWN case over the DSA one, in Sect. 3 we present a comparison of the physical implications in both cases.

3. DISCUSSION

In Sect. 2 we showed that a nonthermal component is needed to properly fit the *NuSTAR* data in the 10 – 20 keV band. As already mentioned, the possible origin of the physical mechanism responsible for such emission is twofold: either DSA or emission from an heavily absorbed PWN. Under a spectroscopic perspective, the only difference between the two scenarios is the presence of an additional absorption component in the PWN case, the *vphabs* model, because of the presence of cold ejecta surrounding the putative compact object. The heavy absorption leads to a negligible contribution of the power-law at energies below 6 keV and thus to a steeper Γ . However, the photon index values in the two scenarios are compatible with each other taking into account the 90% confidence error bars. Moreover, the very similar values of χ^2 does not allow us to exclude one of the two possible emission mechanisms under a merely statistical point of view.

In the DSA scenario, the flux is expected to vary with time in the same way both in the X-ray and radio bands. The synchrotron radio flux of SN87A increased by $\sim 15\%$ between 2012 and 2014 (Cendes et al. 2018). Under the DSA hypothesis, this would be at odd with our findings, since we observe a steady X-ray synchrotron emission, and a 15% increase of the nonthermal X-ray flux is discouraged at the 90% confidence level. However, we note that the DSA scenario could allow for spectral variations in time that would make X-ray variations less predictable from radio variations.

To further investigate the emission nature in the DSA scenario, we replaced the power-law component with the XSPEC *srcut* model, which describes synchrotron emission from an exponentially cut off power-law distribution of electrons in the assumption of homogeneous magnetic field (Reynolds 1998). For each of the three years considered, we constrained the normalization of the *srcut* component (i.e., its flux at 1 GHz, $S_{1\text{GHz}}$), from the corresponding values observed at 9 GHz by Cendes et al. (2018), taking into account the radio spectral index of SN87A ($\alpha = 0.74$, Zanardo et al. 2013). We fixed $S_{1\text{GHz}}$ and α in the *srcut* model and left the break frequency, ν_b , free to vary. The radio-to-X-rays spectral index α can be assumed to be constant since the

synchrotron cooling time is much longer than the age of the system, thus no significant cooling is expected (i.e. $\tau_{\text{sync}} \sim 125 \text{ yr } (E/10 \text{ TeV})^{-1} (B/100 \mu\text{G})^{-2}$). We obtained $h\nu_b = 2.4^{+0.3}_{-0.4} \text{ keV}$, $h\nu_b = 2.3^{+0.3}_{-0.4} \text{ keV}$ and $h\nu_b = 1.7^{+0.4}_{-0.2} \text{ keV}$, in 2012, 2013 and 2014 respectively. The break energy is compatible with 2 keV in the three years considered, though the general trend seems to point towards a decrease with time.

The synchrotron emission of electrons peaks at energy $h\nu = 1.8 \times 10^4 E_{100}^2 B_{100}$, where B_{100} is the magnetic field in units of 100 μG and E_{100} is the electron energy in units of 100 TeV. Considering our break energy $h\nu_b \sim 2 \text{ keV}$, we obtain values of E_{100} spanning from ~ 0.14 to ~ 0.33 , for B_{100} ranging from 6 to 1. This maximum electron energy seems to be quite high, especially considering the relatively low shock speed. In fact, the synchrotron radio emission originates in the HII region (Zanardo et al. 2013; Cendes et al. 2018; Orlando et al. 2019), where the shock velocity is of only 2000 km/s (Cendes et al. 2018, Or20).

In the DSA scenario, we can estimate the acceleration time scale as (Parizot et al. 2006) $\tau_{\text{acc}} = 124\eta B_{100}^{-1} V_s^{-2} E_{100}^{\frac{4}{3}} \text{ yr}$, where η is the Bohm factor and V_s is the shock velocity. In the hypothesis of maximum efficiency ($\eta = 1$) and a standard magnetic field $B_{100} = 1$, with $V_s = 2000 \text{ km/s}$, we would need $\tau_{\text{acc}} \sim 390 \text{ yr}$ to accelerate the X-ray emitting electrons up to the observed maximum energy, i.e., much more than the age of SN87A. The observed maximum energy can be obtained in 25 yr only by assuming that the downstream magnetic field is amplified by the SN87A slow shock up to $\sim 600 \mu\text{G}$ (in this case the maximum electron energy would be of $\sim 14 \text{ TeV}$), and only assuming that the acceleration proceeds at the Bohm limit. The aforementioned issues (steady synchrotron flux and extremely large electron energy in a relatively slow shock) concur in making the DSA scenario unconvincing.

On the other hand, the PWN scenario has a strong physical motivation, as we show below. In the absence of a direct identification of the compact object eventually powering PWN87A, the only possible way to constrain its properties is to use the X-ray luminosity obtained in our analysis to locate the putative PWN within the PWNe population.

The properties of a generic PWN can be associated to those of its progenitor SNR and surrounding ISM introducing the characteristic time and luminosity scales (Truelove & McKee 1999):

$$t_{\text{ch}} = E_{\text{sn}}^{-1/2} M_{\text{ej}}^{5/6} \rho_{\text{ism}}^{-1/3}, \quad (1)$$

$$L_{\text{ch}} = E_{\text{sn}}/t_{\text{ch}}, \quad (2)$$

Table 1. Best-fit parameters of the model adopted to describe *Chandra* and *NuSTAR* observations performed in 2012, 2013 and 2014. Chemical abundances and power law parameters are kept constant along the various years. Uncertainties are at 90% confidence level.

Component	Parameter	2012	2013	2014
TBabs	n_H (10^{22} cm^{-2})	0.235 (fixed)		
vnei	kT (keV)	$2.85^{+0.08}_{-0.07}$	2.87 ± 0.07	2.85 ± 0.05
	O		$0.32^{+0.04}_{-0.03}$	
	Ne		0.57 ± 0.02	
	Mg		0.36 ± 0.03	
	Si		$0.39^{+0.03}_{-0.02}$	
	S		0.65 ± 0.06	
	τ (10^{11} s/cm^3)	$2.1^{+0.9}_{-0.5}$	1.7 ± 0.3	$1.8^{+0.4}_{-0.2}$
	EM (10^{58} cm^{-3})	7.0 ± 0.3	7.3 ± 0.3	$8.2^{+0.3}_{-0.2}$
vnei	kT (keV)	0.60 ± 0.04	$0.65^{+0.04}_{-0.03}$	$0.63^{+0.4}_{-0.3}$
	τ (10^{11} s/cm^3)	$1.8^{+0.5}_{-0.3}$	1.5 ± 0.3	1.7 ± 0.3
	EM (10^{58} cm^{-3})	31^{+2}_{-3}	26 ± 2	28 ± 0.2
pow	Photon index Γ	$2.5^{+0.3}_{-0.4}$		
	$L_{0.5-8}^{\text{pwn}}$ (10^{35} erg/s)	$4.1^{+4}_{-2.8}$		
	L_{1-10}^{pwn} (10^{35} erg/s)	2.6 ± 1.4		
	L_{10-20}^{pwn} (10^{35} erg/s)	$0.32^{+0.02}_{-0.01}$		
	Flux $_{0.5-8}$ ($10^{13} \text{ erg/s/cm}^2$)	93 ± 2	92^{+4}_{-2}	95^{+2}_{-1}
	Flux $_{10-20}$ ($10^{13} \text{ erg/s/cm}^2$)	1.5 ± 0.2	$1.6^{+0.6}_{-0.1}$	$1.6^{+0.3}_{-0.1}$
χ^2 (d.o.f.)		1442 (1221)		

where E_{sn} is the supernova explosion energy, usually assumed to be 10^{51} erg, M_{ej} is the mass in the SNR ejecta and ρ_{ism} the mass density of the ISM. These last parameters have been considered to vary uniformly in: $M_{\text{ej}} \in [5 - 20] M_{\odot}$ (Smartt et al. 2009) and $n_{\text{ism}} \in [0.01 - 10] \text{ cm}^{-3}$ (Berkhuijsen 1987; Magnier et al. 1997; Long et al. 2010; Bandiera & Petruk 2010; Asvarov 2014), where $\rho_{\text{ism}} = m_p n_{\text{ism}}$ and m_p is the proton mass. For the pulsar population the best choice is to consider young γ -ray emitting pulsars (Watters & Romani 2011; Johnston et al. 2020), better suited for describing pulsars powering PWNe than the old radio-emitting ones (Kaspi et al. 2006).

Considering the pulsar parameters from that population (namely the initial spin-down time τ_0 and luminosity L_0), with the choice of the canonical dipole braking index $n = 3$, the PWNe population can be then constructed¹ scaling time and luminosity with the characteristic ones defined in Eq. 1-2. In the $(\tau_0/t_{\text{ch}}, L_0/L_{\text{ch}})$ plane the PWNe population appears as an ellipsoidal surface, where each point corresponds to various physical sources with different combinations of τ_0 , L_0 , M_{ej} and ρ_{ism} . We use this PWNe population to discuss the possi-

ble location of PWN87A, by scaling the observed X-ray luminosity for the corresponding characteristic quantities. In particular, the ejecta mass and kinetic energy are $18 M_{\odot}$ and 2×10^{51} erg, respectively (Or20, and references therein). Given the complex structure of the remnant, the density of the material in which the ejecta expand shows large inhomogeneities, varying from $\sim 0.1 \text{ cm}^{-3}$ in the pre-shock blue supergiant wind, $\sim 100 \text{ cm}^{-3}$ in the HII region, up to $10^3 - 10^4 \text{ cm}^{-3}$ in the dense ring (Sugerman et al. 2005, Or20). We consider a value of $\sim 100 \text{ cm}^{-3}$, representative of the equatorial zone of the HII region. Thus the two scaling are:

$$t_{\text{ch}}^{87A} = 807.6 \text{ yr} \left(\frac{E_{\text{sn}}}{2 \times 10^{51} \text{ erg}} \right)^{-1/2} \left(\frac{M_{\text{ej}}}{18 M_{\odot}} \right)^{5/6} \left(\frac{\rho_{\text{ism}}}{100 m_p \text{ cm}^{-3}} \right)^{-1/3}, \quad (3)$$

$$L_{\text{ch}}^{87A} = 7.86 \times 10^{40} \text{ erg/s} \left(\frac{E_{\text{sn}}}{2 \times 10^{51} \text{ erg}} \right). \quad (4)$$

Taking into account the ejecta absorption, as calculated from our MHD simulation, we derive from the spectral analysis the unabsorbed X-ray luminosity of the central source, finding $L_X^{\text{pwn}} = 4.1^{+4}_{-2.8} \times 10^{35} \text{ erg/s}$ in the $0.5 - 8 \text{ keV}$ band, considering a distance to the source of 51.4 kpc. The X-ray luminosity is related to the spin-down power $\dot{E}$ through the conversion efficiency param-

¹ The PWNe population is liable of changes in the parameters plane, especially if considering a different braking index value than the standard dipole one (Parthasarathy et al. 2020.)

Table 2. Summary of the most probable values for the pulsar spin-down time for different values of the X-ray efficiency η_X , and the associated spin-down luminosity (from the best-fit value). All values are given considering $n_{\text{ism}} = 100 \text{ cm}^{-3}$.

η_X	L_0 (erg/s)	τ_0 (yr)
10^{-5}	4.1×10^{40}	18
10^{-4}	4.1×10^{39}	110
10^{-3}	4.1×10^{38}	480
10^{-2}	4.1×10^{37}	3160
10^{-1}	4.1×10^{36}	25000

ter $\eta_X = L_X / \dot{E}$, that shows very large variations within the population, from $\sim 10^{-5}$ to $\sim 10^{-1}$ (Kargaltsev & Pavlov 2008).

To maintain our analysis as general as possible, here we have considered the entire range of measured η_X . In Fig. 3 we show the population of PWNe as determined associating the γ -ray emitting pulsars with the discussed ranges of parameters for SNRs and ISM (in light blue). The positioning of the putative PWN in SN87A for various η_X is shown as brown lines that intersect horizontally the distribution. The Crab nebula, the Vela nebula and the PWN in Kes 75 are shown for comparison. As it can be easily seen, all the different possibilities lead to a location of PWN87A fully compatible with the population. Assuming an ambient density of 100 cm^{-3} might not be correct for SN 1987A, because of the complex density distribution. Therefore, we relaxed this assumption and, in the right panel of Fig. 3, we show that the putative positions for varying η_X are perfectly consistent with the PWN population, even considering the two extreme density values for the surrounding medium ($0.1 - 10^4 \text{ cm}^{-3}$, Sugerman et al. 2005, Or20).

We verified that our estimate of the PWN luminosity in the X-rays is consistent with the radio luminosity derived by Cigan et al. (2019) under the (very reasonable) assumption that the low energy break frequency of the synchrotron radiation is² $\nu_b \gtrsim 10^{12} \text{ Hz}$, while we expect that the high absorption of the cold ejecta prevents the detectability of the PWN emission in the optical band.

An estimate (purely based on a statistical argument) of the most probable spin-down time for each η_X can be determined by combining the PWNe probability distribution with the estimates of the pulsar spin-down luminosity (Table 2, with $n_{\text{ism}} = 100 \text{ cm}^{-3}$).

In conclusion, a PWN seems to be the most likely source of the synchrotron radiation in hard X-rays of SN87A, even though the DSA scenario cannot be firmly excluded. A more conclusive way to discern between the two scenarios will be provided by future observations. An increase in the hard ($> 10 \text{ keV}$) X-ray flux, similar to that observed in radio would be easily detectable with *NuSTAR*, thus supporting the DSA scenario. A decrease would support the PWN scenario. On the other hand, the rapid ejecta expansion and rarefaction will reduce the soft X-ray absorption of an inner source. In particular, from our MHD simulation, we estimate that the PWN will become detectable by *Chandra* and/or *Lynx*³ (Appendix C) in the 2030s, definitely confirming the PWN scenario.

With the present data and based on our findings, the PWN scenario seems the most likely (and appealing) to account for the nonthermal X-ray emission that we have detected.

Software: CIAO (Fruscione et al. 2006), HEASOFT⁴, XSPEC (Arnaud 1996), NuSTARDAS⁵

Facilities: *Chandra*⁶, *NuSTAR* (Harrison et al. 2013)

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² By assuming a radio photon index $\Gamma_R = 1.5$ and an X-ray photon index $\Gamma_X = 2.5$.

³ <https://www.lynxobservatory.com>

⁴ <https://heasarc.gsfc.nasa.gov/docs/software/heasoft/>

⁵ https://heasarc.gsfc.nasa.gov/docs/nustar/analysis/nustar_swguide.pdf⁶ <https://cxc.harvard.edu/index.html>

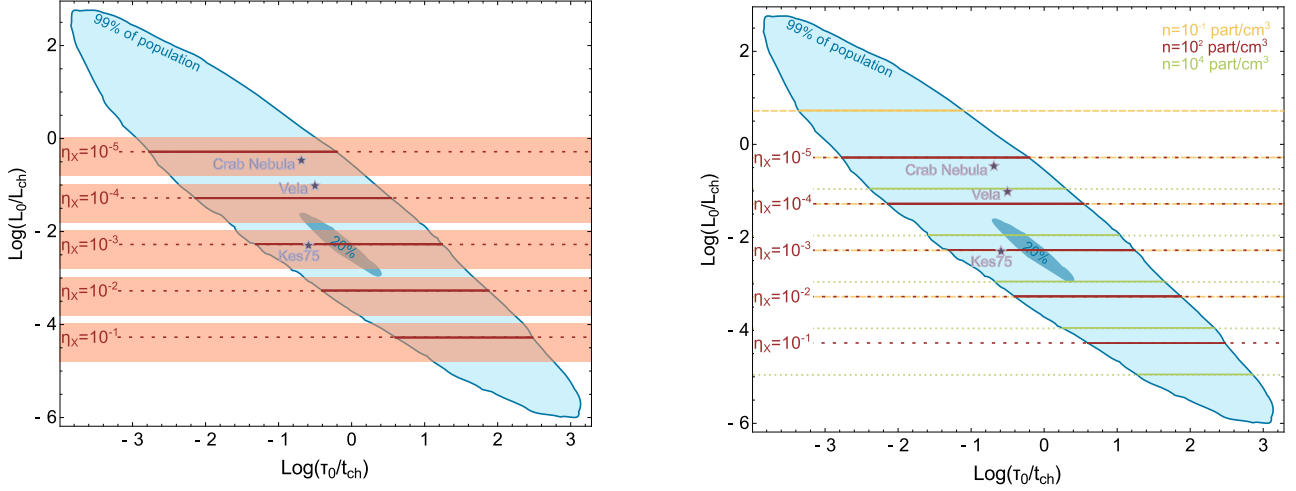


Figure 3. *Left panel.* Positioning of the putative PWN87A (brown lines) within the PWNe population (in light-blue color) in the case of the representative density 100 cm^{-3} of the HII region. The PWNe population is represented with contours enclosing the 99% of the entire population and the 20% (to better mark the center). Different possible positions of the PWN have been obtained using the detected X-ray flux in the 0.5-8 keV band (with the best fit value shown as a solid-dashed line and the entire range of variation as a shaded area) and considering the widest possible variation for the X-ray efficiency η_X (Kargaltsev & Pavlov 2008). *Right panel.* Variation of the predicted location of the PWN87A within the PWNe population, for different values of η_X , considering an extreme variation of the ambient density, from 0.1 cm^{-3} (representative of the blue giant wind, in yellow) to 10^4 cm^{-3} (the maximum density in the dense ring in the HII region, in green).

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APPENDIX

A. DATA REDUCTION

Chandra data were reprocessed with the CIAO v4.12.2 software, using CALDB 4.9.2. We reduced the data through the task *chandra_repro* and we extracted the ACIS-S spectra by using the tool *specextract* which also provided the corresponding arf and rmf files.

NuSTAR data were reprocessed with the standard pipelines provided by the *NuSTAR* data analysis software NuSTARDAS⁷ by using *nupipeline* and *nuproducts*. Details of the observations are reported in Table 3.

Spectral analysis has been performed with XSPEC (v12.11.1, (Arnaud 1996)) in the 0.5 – 8 keV and 3 – 20 keV bands for the *Chandra* and *NuSTAR* data, respectively. All spectra were rebinned adopting the optimal binning procedure described in Kaastra & Bleeker (2016), and the background spectrum to be subtracted was extracted from a nearby region immediately outside of the source. We verified that our results are not affected by the choice of the background regions.

B. X-RAY ABSORPTION FROM COLD EJECTA

We used the 3D MHD simulation of SN87A by Or20 to estimate the absorption pattern of the cold ejecta. The simulation reproduces most of the features observed in the remnant of SN87A in various spectral bands, and links the SNR with the properties of the asymmetric parent SN explosion (Ono et al. 2020) and with the nature of its progenitor star (Urushibata et al. 2018). The model provides all relevant physical quantities in each cell of the 3D spatial domain (MHD variables, plasma composition of the CSM and of the ejecta vs time). The most notable quantities for our purposes are: the electron temperature, the ion density for many chemical species (¹H, ³He, ⁴He, ¹²C, ¹⁴N, ¹⁶O, ²⁰Ne, ²⁴Mg, ²⁸Si, ³²S, ³⁶Ar, ⁴⁰Ca, ⁴⁴Ti, ⁴⁸Cr, ⁵²Fe, ⁵⁴Fe, ⁵⁶Ni, and the decay products, including ⁵⁶Fe), and the ionization age. The model also predicts that the putative NS relic of the supernova explosion has received a kick towards the observer in the north with a lower limit to the kick velocity of ≈ 300 km/s, resulting from a highly asymmetric explosion (Ono et al. 2020, Or20).

For the present study, we assumed a slightly higher kick velocity of 500 km/s. We checked that the results do not change significantly for values ranging between 300 and 700 km/s. Given the kick velocity and the direction of motion of the NS as predicted by the model (Ono

et al. 2020, Or20), and orienting the modeled remnant as it is observed in the plane of the sky (Or20), we established the position of the NS in the 3D spatial domain of the simulation for each year analyzed in this work. Considering that the extension of the putative radio PWN is of the order of < 1000 AU (Cigan et al. 2019), that the X-ray PWN is expected to be smaller than its radio counterpart, and that the spatial resolution of our MHD simulations is of ≈ 180 AU, we can consider the central source as point-like in our procedure.

We reconstructed the absorption pattern encountered by the synchrotron X-ray emission of the putative PWN through the subsequent absorbing layers along the line of sight. The physical effect responsible for the absorption is the photo-electric effect, since the material surrounding the considered source is cold ($T < 100$ K). We extracted values of temperature, column density and abundances associated to each absorbing layer of the 3D domain of the model and we included these parameters in the spectral analysis through the XSPEC photo-electric absorption model *vphabs*.

For the years considered in this work, the absorption due to cold ejecta is comparable with an equivalent H column density higher than 10^{23} cm⁻². This indicates that potential signature of a PWN emission must be searched in the high energy part ($\gtrsim 10$ keV) of the X-ray spectra, less affected by absorption.

C. SYNTHETIC LYNX SPECTRUM

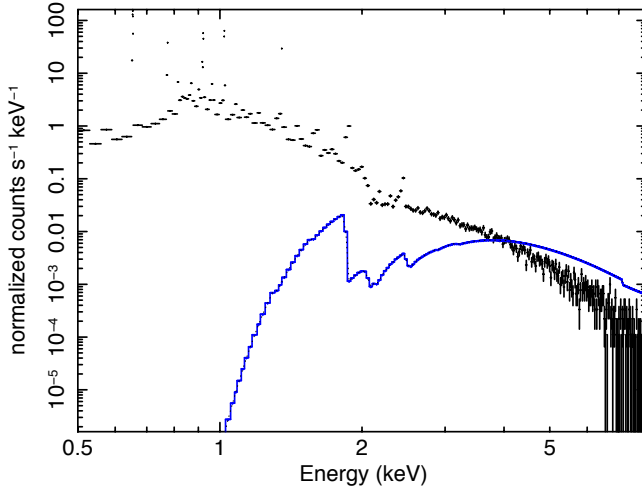
We produced a synthetic LYNX observation of SN87A, as predicted by our MHD model for year 2037. We expect the thermal emission to stay almost constant in the next 15-20 years, though ejecta contribution and/or interaction of the remnant with other inhomogeneities beyond the ring may affect our predictions. Fig. 4 shows the synthetic spectrum extracted from a circular region with radius $R = 0.3''$, well within the bright ring of SN87A (central panel of Fig. 1), assuming an exposure time of 300 ks. If we consider the emission of the PWN absorbed by the ejecta pattern that the model predicts for 2037, we notice that the resulting component would have a flux higher than that associated with thermal emission above 4 keV, thus becoming detectable in the soft x-ray band (see Fig. 4.). Further details about the future detectability of the PWN and on the thermal emission of the putative NS will be described in a forthcoming paper (Greco et al., in prep.).

⁷ https://heasarc.gsfc.nasa.gov/docs/nustar/analysis/nustar_swguide.pdf

Table 3. Summary of the main characteristics of the analyzed observations.

Telescope	OBS ID	PI	Date (yr/month/day)	Exposure time (ks)
<i>Chandra</i>	13735	Burrows	2012/03/28	48
	14417	Burrows	2012/04/01	27
	14697	Burrows	2013/03/21	68
	14698	Burrows	2013/09/28	68
	15809	Burrows	2014/03/19	70
	15810	Burrows	2014/09/20	48
<i>NuSTAR</i>	40001014003	Harrison	2012/09/08	136
	40001014004	Harrison	2012/09/11	200
	40001014007	Harrison	2012/10/21	200
	40001014013	Harrison	2013/06/29	473
	40001014018	Harrison	2014/06/15	200
	40001014020	Harrison	2014/06/19	275
	40001014023	Harrison	2014/08/01	428

Synthetic Lynx spectrum – PWN in 2037

**Figure 4.** Black points: synthetic LYNX spectrum of the emission observed in 2018. Blue line: power-law component absorbed by the ejecta as predicted by the MHD model in 2037.