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Limiting magnetic fields in the cosmic web with diffuse radio emission

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ABSTRACT

We set limits on the presence of the synchrotron cosmic web through the cross-correlation of the 2.3 GHz S-PASS survey with a model of the local cosmic web derived from constrained magnetohydrodynamic (MHD) simulations. The MHD simulation assumes cosmologically seeded magnetic fields amplified during large-scale structure formation and a population of relativistic electrons/positrons from proton–proton collisions within the intergalactic medium. We set a model-dependent 3σ upper limit on the synchrotron surface brightness of $0.16 \text{ mJy arcmin}^{-2}$ at 2.3 GHz in filaments. Extrapolating from magnetic field maps created from the simulation, we infer an upper limit (density-weighted) magnetic field of $0.03 \text{ (} 0.13 \text{)} \mu\text{G}$ in filaments at the current epoch and a limit on the primordial magnetic field (PMF) of $B_{\text{PMF}} < 1.0 \text{ nG}$.

Key words: magnetic fields – plasmas – galaxies: clusters: intracluster medium – intergalactic medium – large-scale structure of Universe – radio continuum: general.

1 INTRODUCTION

The evolution of the diffuse intergalactic medium is difficult to observationally constrain owing to its low density and temperature. The $T = 10^5\text{--}10^7 \text{ K}$ warm-hot intergalactic medium (WHIM; Davé et al. 2001; Cen & Ostriker 2006) is extremely faint at most wavelengths, though it contains as much as half the total baryon content of the local Universe (Bregman 2007). Direct detection of this gas is extremely difficult, though some recent claims have been made (Cantalupo et al. 2014; Eckert et al. 2015; Martin et al. 2015). Indirect methods have been proposed to try and infer the presence of the WHIM, such as the pressure balance in Giant Radio Galaxy lobes (e.g. Malarecki et al. 2013) and Fast Radio Burst dispersion measures (Akahori, Ryu & Gaensler 2016; Fujita et al. 2017). Another indirect method is to observe the non-thermal radio synchrotron emission generated by the relativistic plasma be-

lieved to coexist with the WHIM (Keshet, Waxman & Loeb 2004; Pfrommer et al. 2006; Ryu et al. 2008; Skillman et al. 2008; Araya-Melo et al. 2012; Vazza et al. 2014, 2015a,b). If this plasma is detected at radio wavelengths, it can constrain the origin of the large-scale magnetic fields that permeate the intergalactic medium of galaxy clusters (see Brunetti & Jones 2014, for a review), differentiating, e.g. between the interaction of radio and star-forming galaxies with the WHIM (e.g. Donnert et al. 2009) and a primordial magnetic field (Marinacci et al. 2015). Observationally inferring the presence of magnetic (B) fields outside of dense clusters of galaxies is extremely difficult (e.g. Xu et al. 2006), and B field estimates from simulations can vary from $0.2 \mu\text{G}$ (Cho & Ryu 2009) to $0.005 \mu\text{G}$ (Donnert et al. 2009), depending on the models used for seeding and amplifying the fields. Any measurement of these currently unconstrained values would provide a useful constraint for differentiating between the possible models.

The presence of a faint radio ‘Synchrotron Cosmic Web’ (which we will call simply cosmic web for short) that traces the WHIM has yet to be observationally confirmed (Wilcots 2004; Vazza et al. 2015b). The physical mechanisms producing this

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synchrotron emission are the acceleration of relativistic particles by shocks and turbulence, and recent cosmological simulations have now focused on tracing these mechanisms in the filamentary regions of the cosmic web (Araya-Melo et al. 2012; Vazza et al. 2014, 2015a,b), and they predict surface brightnesses up to a few $\mu\text{Jy arcmin}^{-2}$ in the largest filaments at 1.4 GHz, which is currently only attainable through the statistical analysis we propose below.

Besides intrinsically low surface brightness, diffuse emission from the Milky Way presents yet another obstacle for the detection of the synchrotron cosmic web. Synchrotron emission from the Milky Way dominates the cosmic web at all frequencies and scales, though surface-brightness fluctuations decrease at smaller angular scales (La Porta et al. 2008). Unlike the foreground issues in cosmic microwave background (CMB) analysis, this foreground has a similar frequency behaviour to the cosmic web, and thus cannot be removed with template matching at other wavelengths (e.g. Gold et al. 2011). At the present time, statistical detection of the cosmic web may be the only option to get around the problem of confusion from both the Galaxy and background extragalactic sources.

One possible method of statistical detection is cross-correlation. Cross-correlation is a technique to statistically detect emission below the noise level of a map (both instrumental and source noise). The widest use of this technique has been on the CMB, where it has been used to detect the late Integrated Sachs–Wolfe (ISW) effect (e.g. Crittenden & Turok 1996; Boughn & Crittenden 2004; Nolta et al. 2004; McEwen et al. 2007; Liu, Ng & Pen 2011). Cross-correlation can be performed in many ways, and has been attempted successfully in real (e.g. Nolta et al. 2004), harmonic (Afshordi, Loh & Strauss 2004) and wavelet (McEwen et al. 2007) space. For our purposes, cross-correlation can be used to infer the presence and average surface brightness of the synchrotron cosmic web (e.g. Brown, Farnsworth & Rudnick 2010; Brown 2011; Vernstrom et al. 2017), provided that we have a plausible model for the spatial distribution of the diffuse radio emission. This method was first attempted on a small region of the northern sky using single-dish observations at 1.4 GHz (Brown et al. 2010), and more recently utilizing deep Murchison Widefield Array (MWA) data at a frequency of 180 MHz (Vernstrom et al. 2017). Both studies used the distribution of galaxies (binned in redshift) as tracers of large-scale structure. In the case of Vernstrom et al. (2017), they further modified the galaxy distribution (galaxies per pixel) to model the distribution of the diffuse radio emission through convolution at Mpc scales and utilized the scaling relations derived in Vazza et al. (2015b) to derive magnetic field upper limits. One limitation to this approach is the lack of a detailed spatial model of how the radio emission correlates with large-scale structure (beyond linearly), something that the current work attempts to overcome using constrained magnetohydrodynamic simulations of the local Universe.

In this paper, we cross-correlate the radio continuum map from the S-Band Polarization All-Sky Survey (S-PASS) with a model of the synchrotron cosmic web obtained through a cosmological simulation of the local Universe. In Section 2, we describe the S-PASS data, the model of the cosmic web and our pre-processing of the images. Section 3 outlines the cross-correlation procedure, results and how we compute upper limits on the surface brightness of the cosmic web, and in Section 4 we discuss the implications of our limits for theories of cosmic magnetism and particle acceleration in large-scale structure, as well as prospects for direct detection with upcoming Square Kilometre Array pathfinders.

2 THE DATA

2.1 S-PASS Radio Data

S-PASS is a 2.3 GHz continuum polarization survey of the southern sky using the Parkes 64 m telescope operated by CSIRO Astronomy & Space Science in NSW, Australia (Carretti et al. 2013b). The current paper only utilizes the total intensity image produced from the survey. The 2.3 GHz system at Parkes has a resolution of 9 arcmin; however, final S-PASS maps have a resolution of 10.75 arcmin, after a convolution with a Gaussian kernel with FWHM = 6 arcmin applied during the map-making process (details will be presented in the forthcoming S-PASS survey paper; Carretti et al., in preparation). With a sensitivity of 1 mJy beam^{-1} , and high sensitivity to large-scale diffuse emission, S-PASS is ideally suited for detecting synchrotron emission within the cosmic web. The point-source confusion limit is 13 mJy beam^{-1} in the total intensity images (Meyers et al. 2017). The survey data were converted into HEALPIX format (Górski et al. 2005) with $N_{\text{side}} = 1024$, which has an approximate pixel size of 3 arcmin (Fig. 1, left-hand side). We have already utilized S-PASS to detect a synchrotron bridge connecting the cluster centre to the major outgoing merger shock in the massive cluster Abell 3667 (Carretti et al. 2013a), which demonstrates the ability of S-PASS to detect faint extragalactic synchrotron emission associated with large-scale structure.

The S-PASS image is dominated by large-scale Galactic emission, so we have subtracted out a 300 arcmin median-weight filter. The filter size was chosen to optimize the removal of the Galactic foreground without subtracting too much of the potential cosmic web emission, which we estimated from our model below. After this filtering, Galactic and extragalactic point sources remain, as well as diffuse emission less than ~ 300 arcmin in extent. In order to remove the point sources, we further applied a circular spatial-scale filter of diameter 27 arcmin, chosen to be $2.5\times$ the beam to remove bright compact sources with <1 per cent residuals (Rudnick 2002). Fig. 1 (left-hand side) shows the resulting map, which has an rms noise of 26 mJy beam^{-1} (mostly confusion from residual compact sources and Galactic diffuse emission) in the clean unmasked regions (see Section 3.1). It should be noted that diffuse emission of any spatial scale *below the noise level of the map* will not get removed by the median or spatial-scale filters (Rudnick 2002). We account for any missing cosmic web flux in our analysis of the upper limits below.

2.2 Synchrotron cosmic web model

In order to spatially resolve filaments of galaxies, which have a typical width of 1–2 Mpc (e.g. Ryu et al. 2003), we restrict our redshift range of interest to $z < 0.048$ in order to have at least 2×10 arcmin S-PASS beams across a 1 Mpc filament¹. We have therefore adopted as a model the simulations of Dolag et al. (2004, 2005) and Donnert et al. (2009), which trace the amplification of magnetic fields, along with the thermal gas, in a constrained simulation that reproduces local large-scale structure. The simulation used initial conditions based on the IRAS 1.2 Jy galaxy survey (Fisher et al. 1995), and reproduced local clusters and super-structures with sizes >3 Mpc (e.g. Virgo, Coma, Centaurus, Hydra, Perseus, A3627), including some smaller structures such as the merging substructure within

¹ Resolving filaments is necessary to distinguish the diffuse filamentary IGM radio emission from compact radio/star-forming galaxies within the filaments.

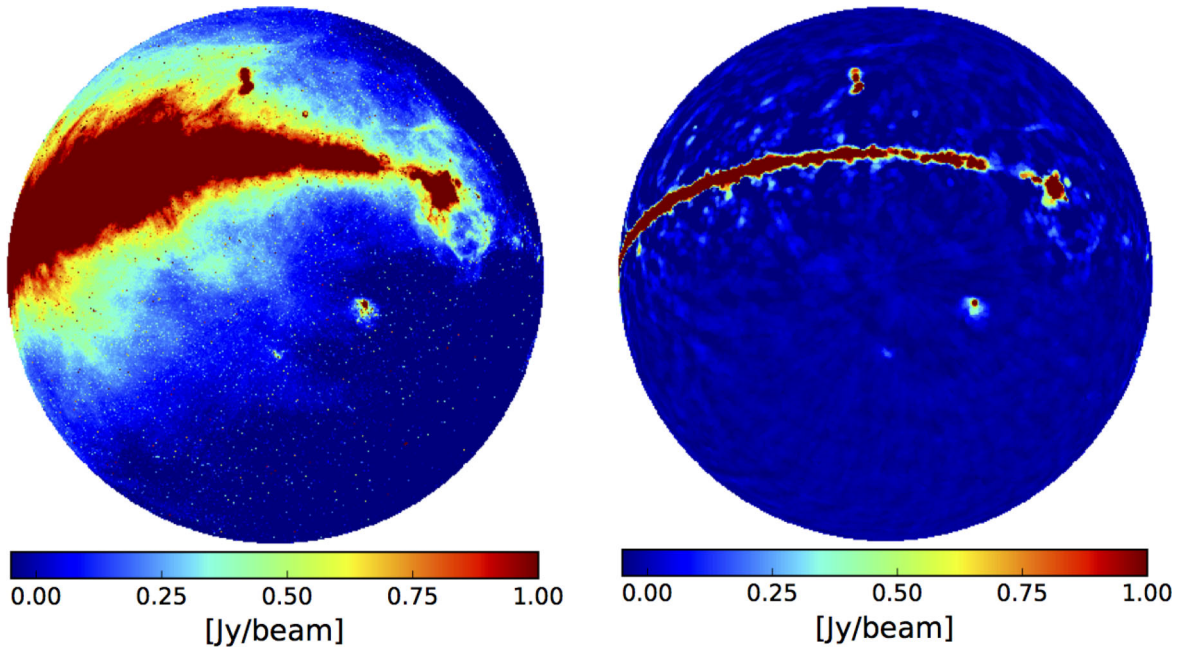


Figure 1. Left-hand side: original S-PASS image orthographic projection (half the sky) centred on the South Celestial Pole. Right-hand side: residual small-scale emission after subtracting a median-weight filtered image (300 arcmin circular filter) and applying a spatial scale filter (Rudnick 2002) that removes emission less than 27 arcmin in size.

the Coma cluster (Donnert et al. 2010). The source of the radiating cosmic ray (CR) electrons in the simulation is secondary electrons/positrons ($\text{CRe}^\pm$) created from collisions between CR protons and the thermal protons of the WHIM (Dennison 1980; Blasi & Colafrancesco 1999). The motivations for using the secondary model are twofold: (1) the CR protons should be shock-accelerated during the build up of large-scale structure (LSS) and live for a Hubble time, making their presence in filaments expected at some level (e.g. Blasi, Gabici & Brunetti 2007), and (2) the number density of these $\text{CRe}^\pm$ should be low in filaments, so using this model should produce a conservative upper limit on the magnetic fields.

Donnert et al. (2009) used different models for seeding the magnetic fields at high redshift and traced the amplification of these fields during the collapse of the local structures. These simulations have been used to model the giant radio halo of the Coma cluster (Donnert et al. 2009), Faraday rotation from the local cosmic web (Stasyszyn et al. 2010, also through cross-correlation) and the origin of fast radio bursts (FRBs; Dolag et al. 2015). Fig. 2 shows the magnetic field strength as a function of density for the various magnetic field models used with the simulation. The magnetic field values in filamentary regions for all models in Fig. 2 are lower than the work of, e.g. Ryu et al. (2008) and Cho & Ryu (2009), which predict $B \sim 0.001 \mu\text{G}$, though there is no discrepancy in dense cluster cores.

The model that we explore in this paper (which we call *cosmological*) is the one presented by Dolag et al. (2004), which uses a uniform cosmological magnetic field of 10^{-11} G seeded at high redshift and evolved to the present epoch (green line in Fig. 2). This primordial magnetic field (PMF) is an order of magnitude lower than the upper limits placed by Planck (Planck Collaboration XIX 2016), though in Fig. 2 we can see that this model predicts the highest values for magnetic fields within filaments, and is the most optimistic model for detection given the limited sensitivity of the S-PASS maps. The other two models in Fig. 2 consider seeding magnetic fields through dipolar galactic outflows at either a single epoch (*galactic outflows*) or multiple epochs (*repeated galactic*

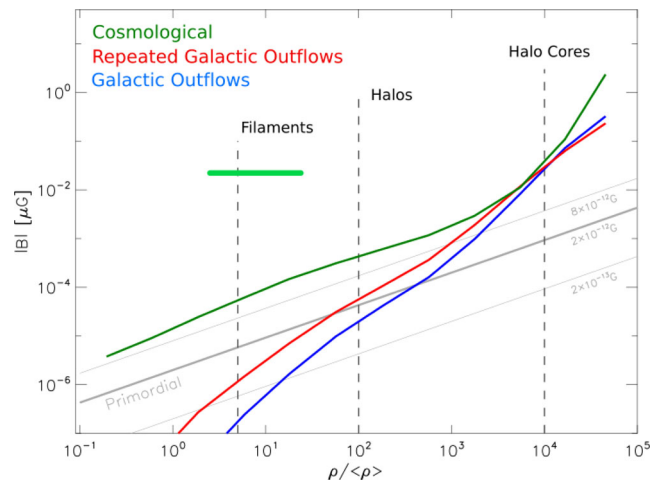


Figure 2. Plot of average magnitude of magnetic field as a function of fractional density for the models used by Dolag et al. (2004, 2005) and Donnert et al. (2009). This work focuses on the *cosmological* model, which predicts the highest values within filaments. The horizontal green bar is the 3σ upper limit derived in this work. $\langle\rho\rangle$ is the average density of the simulation.

outflows), both predicting significantly weaker magnetic fields in filamentary regions compared to the *cosmological* model. The solid grey lines show the results of simple adiabatic compression of a primordial field, for reference. Fig. 3 shows the electron density and 1.4 GHz radio emission from this model. The radio emission in the cluster cores is roughly what is currently observed in the Coma cluster (e.g. Brown & Rudnick 2011), though the simulation clearly shows more diffuse emission in the filamentary regions; this is the as yet undetected synchrotron cosmic web.

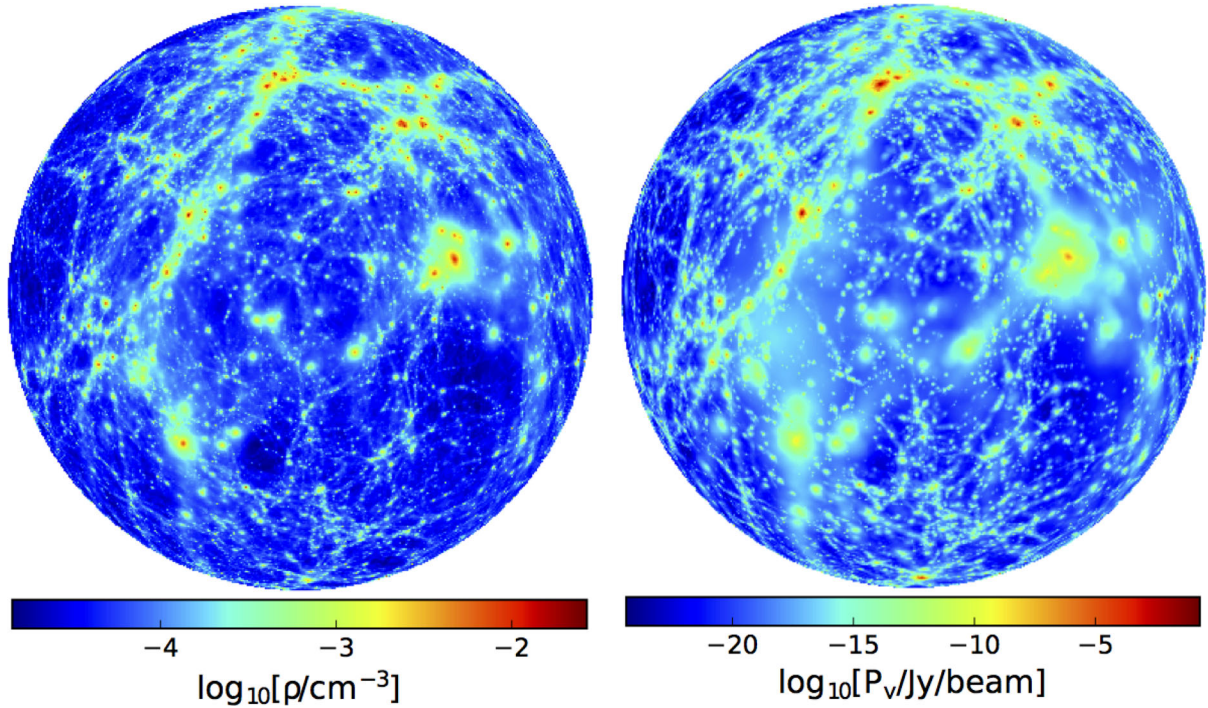


Figure 3. Left-hand side: electron density map ($\log_{10}(n_e)$ [cm^{-3}]) from the *cosmological* model of (Dolag et al. 2004). Right-hand side: radio synchrotron map ($\log_{10}(I)$ [mJy arcmin^{-2}]) at 1.4 GHz from the *cosmological* model.

3 CROSS-CORRELATION ANALYSIS AND UPPER LIMITS

We perform a cross-correlation of the S-PASS and the cosmic web map by using ANAFAST implemented in the PYTHON module HEALPY. We first calculate the pseudo cross-power spectrum (Hivon et al. 2002) through ANAFAST’s spherical harmonic decomposition and then calculate a cross-correlation function (CCF) from the measured C_l values from

$$\text{CCF}(\theta) = \sum_l \frac{2l+1}{4\pi} C_l P_l(\cos\theta), \quad (1)$$

where the P_l are Legendre polynomials. The error bars used below are calculated using cosmic variance only and are given by $\text{Var}(C_l) = 2C_l^2 / f_{\text{sky}}(2l+1)$, which should be regarded as lower limits. Before we perform the cross-correlation, we need to construct the appropriate mask, as well as measure the likelihood of spurious correlations in the data.

3.1 The mask

We are only interested in correlations between the two maps that correspond to low-density regions of the cosmic web and must therefore mask out other possible sources of correlation. Our first consideration is the Galactic plane, which needs to be excluded for several reasons. First, the S-PASS map (Fig. 1) shows very large-scale diffuse emission in the plane which is for the most part not extragalactic. In addition, the model of the cosmic web was based on constrained simulations of the local LSS, which used initial conditions based on the IRAS 1.2 Jy galaxy survey (Fisher et al. 1995); these initial conditions could be contaminated by Galactic dust extinction, resulting in an anticorrelation (see Dolag et al. 2004). We therefore applied the Planck CMB mask (Planck Collaboration XIX 2016) which masks out the Galactic plane and many of the

brightest extragalactic point sources. Residual emission remained due to the Galaxy and the bright extragalactic source Centaurus A (Feain et al. 2011), so we further masked regions where the filtered S-PASS map was $> 0.1 \text{ Jy beam}^{-1}$. We have also masked out regions with extinction $E(B - V) > 0.1$ from the Schlegel, Finkbeiner & Davis (1998) maps, as well as galactic latitudes $|b| < 10^\circ$.

As we are interested in only the filamentary regions of the cosmic web, we have also masked out regions where the thermal electron density is greater than 10^{-3} cm^{-2} . The results are shown in Fig. 4. The total unmasked fraction of the sky is $f_{\text{sky}} = 26$ per cent.

3.2 Random correlations

In order to interpret the significance of the result we need to know the range of CCF values expected when there is no real correlation present. To do this, we must cross-correlate the S-PASS image with random maps (which should have no true correlation) that have the same physical characteristics as the LSS simulation, such as the same mean and clustering properties. In order to simulate random maps, while insuring that the statistics are as close as possible to the original, we have created 10 new versions of the *cosmological* model by rotating/flipping the all-sky map by combinations of 45° , 90° and 180° in Galactic latitude and longitude. These rotated maps, which should no longer be correlated with the real radio sky, were also cross-correlated with the S-PASS image. Fig. 5 (top) shows the plot of the CCF as a function of angular scale. The blue line shows the CCF of S-PASS and the original *cosmological* model map. The dashed red line shows the mean CCF of S-PASS and the 10 rotated maps, the shaded blue region shows the 1σ range of the CCFs² and the dashed green line shows the 3σ upper envelope. Note that the

² The 1σ from the mean is obtained by the quadrature sum of the rms of the 10 random maps and the cosmic variance in each angular bin.

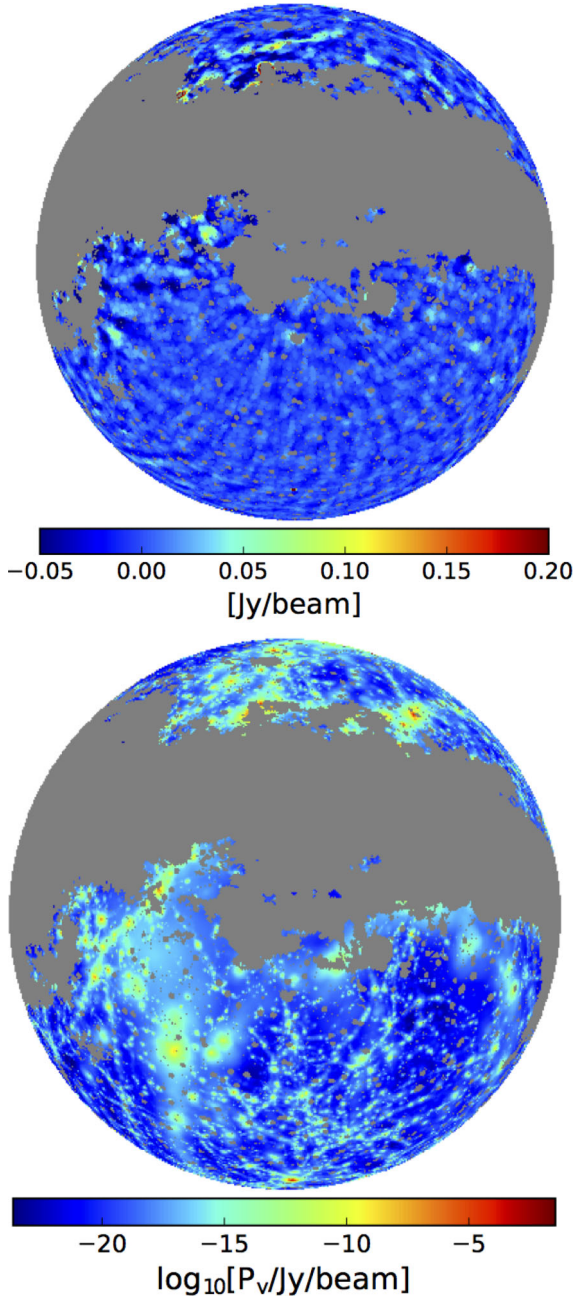


Figure 4. Top: S-PASS image after all the filtering and with the mask applied. Note that the scale has been changed compared to Fig. 1 in order to bring out low-level features. Bottom: the radio emission from the simulation with the same mask applied. The total unmasked fraction of the sky is $f_{\text{sky}} = 26$ per cent.

correlation function is slightly above the mean of the random maps at small angles, but well below even the 1σ significance level, and is considered a non-detection.

UPPER LIMITS: To estimate an upper limit from this non-detection, we must add a fake cosmic-web signal into the S-PASS map (below the noise) and repeat the cross-correlation procedure. To do this, we take one of the rotated cosmic web maps and add it to the S-PASS map iteratively with an increasing multiplicative factor (we have chosen the map that has been rotated by 180° in both Galactic latitude and longitude). For each iteration we repro-

duced the cross-correlation, and terminated the procedure when the cross-correlation of the injected S-PASS map and the corresponding rotated cosmic web map reaches roughly 3σ at any angular scale. The injection takes place before the spatial scale filtering, and we repeat the exact process of cross-correlating this injected map with the rotated model maps. The filtering will remove some of the injected signal, alter the S-PASS map and thus affect the shapes of the correlation functions, all of which should be included when calculating the limit. We show the cross-correlation of the model-injected S-PASS map at the point of termination in Fig. 5 (bottom), where the cross-correlation function with the injected model (times the scale factor) has reached 3σ at an angular scale of $\theta \approx 1^\circ$ (solid red line in Fig. 5). We should note that the correlation persists at a low level out to 5° of our 300 arcmin median-weight filter. From the radio map of the *cosmological* model (Fig. 3, right-hand side), we used the multiplicative factor at convergence (800 ± 10) to create an upper limit map; had the radio emission from the cosmic web been as bright as this map, we would have detected it at the 3σ level. From this map we measured the average surface brightness in filaments (defined as $3 \times 10^{-4} < n_e < 5 \times 10^{-3} \text{ [cm}^{-3}\text{]})$ to be $20.9 \text{ mJy beam}^{-1}$ or $0.16 \text{ mJy arcmin}^{-2}$. The density range corresponds to an overdensity range of $\delta = 3\text{--}50$, typical of filaments and cluster outskirts (e.g. Vazza et al. 2014). Fig. 6 (top) shows the upper limit map, masking out all but the filamentary regions (in the n_e range described above) where the average radio brightness was taken.

4 DISCUSSION AND CONCLUSION

We now investigate what the limit for the *cosmological* model means for our knowledge of the magnetic fields in large-scale structure, as well as the prospects for detection of the cosmic web with upcoming surveys.

4.1 Magnetic fields

To estimate the average magnetic field from the surface brightness, we extracted the magnetic field values from the *cosmological* simulation (Fig. 6) and scaled it relative to the amount that the 3σ upper limit surface brightness is greater than the simulations brightness. In order to do this scaling, we must take into account both the spectral index of the radio emission α and the scaling of the synchrotron emissivity with the magnetic field. In the limit where the magnetic field B is much less than $B_{\text{CMB}}/(1+z)^2 = 3.24 \text{ } \mu\text{G}$ (the equivalent effective magnetic field strength used to parametrize the electron Compton energy losses), the synchrotron emissivity ϵ_r of the simulation goes as

$$\epsilon_r \sim X_p n_e B^{1+\alpha} \nu^{-\alpha}, \quad (2)$$

so we scale the magnetic field from the simulation by a factor of

$$\left(\frac{(\epsilon'_r/\epsilon_r)}{(\nu'/\nu)^{-\alpha}} \right)^{1/(1+\alpha)}, \quad (3)$$

where ϵ'_r is the average surface brightness of our injected signal, $\nu' = 2.3 \text{ GHz}$ (the S-PASS frequency), $\nu = 1.4 \text{ GHz}$ (the simulation frequency), n_e is the number density of thermal electrons (in $\text{[cm}^{-3}\text{]})$ and $\alpha \approx 1.0$ for the secondary electrons assumed in the simulation. The factor X_p is the relative fraction of the energy density of CR protons with respect to the thermal gas, and for the current simulations a constant value of $X_p = 0.01$ was used (Dolag et al. 2004). By taking the average value of the scaled magnetic field map in the same density regions as above, we obtain an upper limit magnetic field value of $0.03 \text{ } \mu\text{G}$. However, given that we are constraining

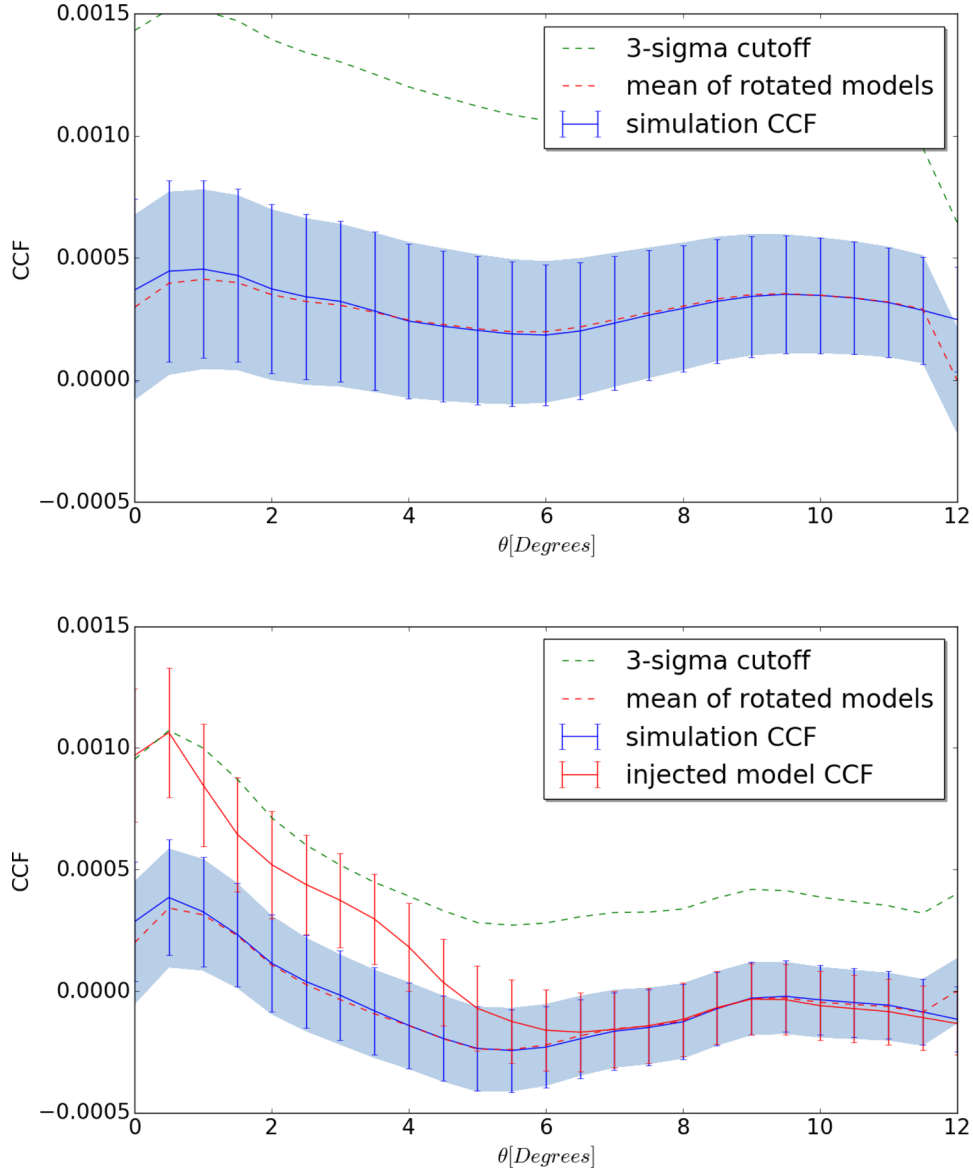


Figure 5. Top: The cross-correlation function for the *cosmological* model with the S-PASS radio map is the blue solid line. The blue shaded region is the $\pm 1\sigma$ from the mean (red dashed line) of the rotated models correlations with the S-PASS map. The green dashed line is the 3σ upper threshold. Bottom: Correlation functions as above after injecting fake cosmic web emission into the S-PASS map. The solid red line is the cross-correlation function of the S-PASS map and the rotated model that was injected.

the magnetic field using synchrotron emission, a density-weighted average of

$$B_{\text{filament}} < 0.13 \mu\text{G} \left(\frac{0.01}{X_p} \right)^{1/(1+\alpha)} \quad (4)$$

over the region is more appropriate, where we have also included the dependency on X_p and α . Fig. 6 (bottom) shows a map of the magnetic field limits in the filamentary overdensity range where the average was taken. We should note that this value is dependent on the assumed model of the local magnetized cosmic web, especially the assumption that the $\text{CRe}^\pm$ were created through hadronic interactions, and does not include any primary CR electrons produced in shocks and/or turbulence (e.g. Vazza et al. 2015a,b). Not including these acceleration mechanisms makes this a true upper limit on the field, as simulations focusing only on primary electrons directly accelerated at structure formation shocks find comparable

or higher densities of $\text{CRe}^\pm$ (e.g. Skillman et al. 2008; Araya-Melo et al. 2012; Vazza et al. 2014, 2015a). Under the assumed model, we can use equations (3) and (4) to infer a limit on the primordial magnetic field from the simulation to obtain

$$B_{\text{PMF}} < 1.0 \text{ nG} \left(\frac{0.01}{X_p} \right)^{1/(1+\alpha)}. \quad (5)$$

To obtain this we take the assumed 10^{-11} G PMF of the simulations, scale it by the factor in equation (3), keeping the model dependence explicit as in equation (4). This is a factor of a few below the limits obtained by Planck ($B_{1\text{Mpc}} < 4.4$ nG; Planck Collaboration XIX 2016) and comparable to the recent limit derived through a joint Planck + South Pole Telescope analysis ($B_{1\text{Mpc}} < 1.5$ nG; Zucca et al. 2017). However, constraints from the Coma cluster restrict $X_p \in [10^{-3}, 10^{-2}]$, which introduces a significant model uncertainty into equations (3) and (4) (e.g. Donnert 2013).

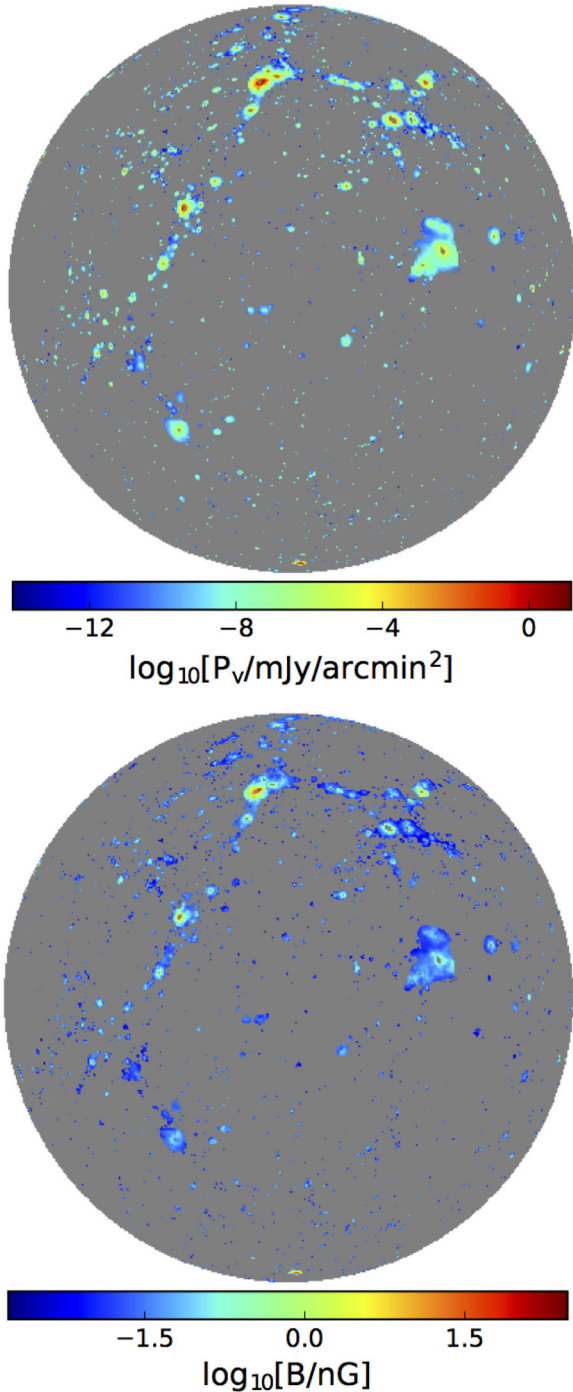


Figure 6. Top: orthographic projection of the upper limit *cosmological* model of the cosmic web with all structure except for the filaments masked out (see the text). The mean synchrotron brightness in the filamentary structures is $0.16 \text{ mJy arcmin}^{-2}$. Bottom: the magnetic field upper limit map derived from the same simulation, with the same mask. The average upper limit (mass-weighted) field is $B < 0.03$ (0.13) μG .

Vernstrom et al. (2017) recently performed a similar cross-correlation experiment using low-frequency MWA data and the distribution of galaxy counts. The approach taken in Vernstrom et al. (2017) uses deeper observations at lower frequency, both of which are advantages when searching for faint, steep spectrum emission (Brown 2011), though over a smaller area of the sky than presented here. Without access to a direct spatial model of the synchrotron

cosmic web at higher redshift, Vernstrom et al. (2017) used a scaling relation derived from Vazza et al. (2015b) to find magnetic field limits of $0.03\text{--}1.98 \mu\text{G}$, depending on the assumed spectral index, which is comparable to the limits derived here. While the Vernstrom et al. (2017) work reached a deeper surface-brightness limit with the MWA, this work was able to achieve similar limits on the magnetic field by referencing a single simulated model of the local Universe. This in turn means that the limit presented here suffers from being highly dependent on the model of Dolag et al. (2004), in particular the source of the radiating $\text{CRe}^\pm$. Future observations that can independently constrain the CR electron production mechanism can make observations of the cosmic web a competitive and complementary method for constraining primordial magnetic fields.

4.2 Future direct detection

The limits on the surface brightness of the synchrotron cosmic web inferred above are deep compared to current radio observations due to the statistical reduction in the noise allowed by cross-correlation. The S-PASS limit derived at 2.3 GHz corresponds to $I_{1.4 \text{ GHz}} < 0.073 \mu\text{Jy arcsec}^{-2}$, while the average brightness of diffuse synchrotron emission in clusters of galaxies is typically around $1 \mu\text{Jy arcsec}^{-2}$ at 1.4 GHz (Brunetti & Jones 2014). This brightness limit is not currently out of reach for deep, interferometric radio observations, though confusion from point sources currently sets the limit on our ability to reach these levels in practice (Vernstrom et al. 2014, 2017). Vernstrom et al. (2017) demonstrated that confusion could be overcome statistically through cross-correlation, but direct detection will require detailed modelling and subtraction of detected compact sources.

Upcoming all-sky surveys will reach sensitivities where direct detection of the brightest portions of the synchrotron cosmic web will be made possible. For example, the Evolutionary Map of the Universe (Norris et al. 2011) survey with the Australian Square Kilometre Array Pathfinder (ASKAP) will have an rms noise on the order of $0.1 \mu\text{Jy arcsec}^{-2}$ and could potentially reach the above limit directly so long as the interferometric nature of ASKAP does not filter out the diffuse emission (Brown 2011). Given the optimistic value of X_p assumed in this work, and the distance between our upper limit and the simulation magnetic field in Fig. 2, the true radio emissivity in the cosmic web might be out of reach for the EMU survey. See Vazza et al. (2015b) for a comparison of upcoming radio surveys (including LOFAR, MeerKAT and SKA) and their ability to probe the synchrotron cosmic web when primary accelerated electrons are considered.

To summarize our results, we performed a cross-correlation between the S-PASS radio survey at 2.3 GHz with a simulated model of the local synchrotron cosmic web and were able to place model-dependent limits on the average surface brightness ($< 0.16 \text{ mJy arcmin}^{-2}$ at 2.3 GHz) and density-weighted magnetic field ($< 0.13 \mu\text{G}$) at the 3σ level. We further placed a limit on the primordial magnetic field ($B_{\text{PMF}} < 1.0 \text{ nG}$) that is competitive to the current CMB limits set by Planck + SPT (Zucca et al. 2017), highlighting the potential power of a direct detection of the synchrotron cosmic web with upcoming radio telescope surveys.

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