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Shock acceleration efficiency in radio relics

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ABSTRACT

Context. Radio relics in galaxy clusters are giant diffuse synchrotron sources powered in cluster outskirts by merger shocks. Although the relic–shock connection has been consolidated in recent years by a number of observations, the details of the mechanisms leading to the formation of relativistic particles in this environment are still not well understood.

Aims. The diffusive shock acceleration (DSA) theory is a commonly adopted scenario to explain the origin of cosmic rays at astrophysical shocks, including those in radio relics in galaxy clusters. However, in a few specific cases it has been shown that the energy dissipated by cluster shocks is not enough to reproduce the luminosity of the relics via DSA of thermal particles. Studies based on samples of radio relics are required to further address this limitation of the mechanism.

Methods. In this paper, we focus on ten well-studied radio relics with underlying shocks observed in the X-rays and calculate the electron acceleration efficiency of these shocks that is necessary to reproduce the observed radio luminosity of the relics.

Results. We find that in general the standard DSA cannot explain the origin of the relics if electrons are accelerated from the thermal pool with an efficiency significantly smaller than 10%. Our results show that other mechanisms, such as shock re-acceleration of supra-thermal seed electrons or a modification of standard DSA, are required to explain the formation of radio relics.

Key words. acceleration of particles – radiation mechanisms: non-thermal – radiation mechanisms: thermal – galaxies: clusters: intracluster medium – galaxies: clusters: general – shock waves

1. Introduction

Astrophysical shock waves are able to accelerate particles over a broad range of scales, from astronomical units in the Sun heliosphere to megaparsecs in clusters of galaxies. Among the numerous physical processes proposed, the diffusive shock acceleration (DSA) theory provides a general explanation of particle acceleration in several astrophysical environments (e.g. Blandford & Eichler 1987). This process is based on the original idea of Fermi (1949), according to which particles are scattered upstream and downstream of the shock by plasma irregularities, gaining energy at each reflection.

Radio relics in galaxy clusters are giant synchrotron sources that are explained assuming that relativistic particles are accelerated by shocks crossing the intra-cluster medium (ICM) (e.g. Enßlin et al. 1998; Roettiger et al. 1999). Whilst the relic–shock connection is nowadays well consolidated by radio and X-ray observations (see Brunetti & Jones 2014; van Weeren et al. 2019, for reviews), the details of the acceleration mechanisms are still not fully understood.

To date, the acceleration efficiency of cosmic rays (CR) at astrophysical shocks is mainly constrained by studies of supernova remnants (SNR) in our Galaxy, where strong shocks ($\mathcal{M} \sim 10^3$) propagate in a low-beta plasma ($\beta_{pl} = P_{th}/P_B$, i.e., the ratio between the thermal and magnetic pressures) medium and are able to transfer $\sim 10\%$ or more of the energy flux through them into cosmic ray protons (CRp), and a smaller fraction into cosmic ray electrons (CRE; e.g., Jones 2011; Morlino & Caprioli 2012; Caprioli & Spitkovsky 2014; Caprioli et al. 2015; Park

et al. 2015). In contrast, radio relics in the outskirts of merging galaxy clusters probe particle acceleration in action at much weaker shocks ($\mathcal{M} \lesssim 3 - 5$) and in a high- β_{pl} environment such as the ICM, where the thermal pressure dominates over the magnetic pressure. In this case, the acceleration efficiency of CRp is still poorly understood, although current models predict efficiencies that are less than a few percent (e.g., Kang & Jones 2005; Kang & Ryu 2013), in agreement with direct constraints coming from γ -ray observations of galaxy clusters (e.g., Ackermann et al. 2010, 2014, 2016). On the other hand, the observed connection between radio relics and shocks in merging galaxy clusters demonstrates that electron acceleration (or re-acceleration) at these shocks is efficient, implying surprisingly large values of the ratio of CRE/CRp acceleration efficiencies if these particles are extracted from shocks by the same population (e.g., the thermal ICM). This poses fundamental questions on the mechanisms leading to the formation of relativistic particles in radio relics (e.g., Vazza & Brüggen 2014; Vazza et al. 2015, 2016).

In recent years, deep X-ray observations performed with *Chandra*, *XMM-Newton*, and *Suzaku* have led to an increase in the number of shocks detected in merging galaxy clusters (e.g., Akamatsu et al. 2017; Canning et al. 2017; Emery et al. 2017; Botteon et al. 2018; Thölken et al. 2018; Urdampilleta et al. 2018, for recent works). In a few cases, when the shock front is co-spatially located with a radio relic, it has been shown (under reasonable assumptions on the minimum momentum of the accelerated electrons) that DSA is severely challenged by the large acceleration efficiencies required to reproduce the total radio luminosity of the relics if particles are accelerated from the ICM

thermal pool (Botteon et al. 2016a; Eckert et al. 2016; Hoang et al. 2017).

To mitigate the problem of the high acceleration efficiencies implied by cluster shocks, recent theoretical models assume a pre-existing population of CRe at the position of the relic that is re-accelerated by the passage of the shock (e.g., Markevitch et al. 2005; Macario et al. 2011; Kang & Ryu 2011; Kang et al. 2012, 2014; Pinzke et al. 2013). This re-acceleration scenario seems to be supported by the observation of radio galaxies located nearby or within a number of radio relics (e.g., Bonafede et al. 2014; Shimwell et al. 2015; Botteon et al. 2016a; van Weeren et al. 2017; Di Gennaro et al. 2018).

In order to test the scenario of shock acceleration of thermal particles as the origin of radio relics, in this paper, for the first time we evaluate the efficiency of particle acceleration at cluster shocks using a homogeneous approach in a relatively large number of radio relics. We adopt a Λ CDM cosmology with $\Omega_\Lambda = 0.7$, $\Omega_m = 0.3$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ throughout.

2. Theoretical framework

2.1. Computation of the acceleration efficiency in radio relics

The electron acceleration efficiency η_e can be related to the observed synchrotron luminosity that is produced by the shock-accelerated electrons. In particular, η_e is evaluated assuming that a fraction of the kinetic energy flux from a shock with speed V_{sh} , surface A , compression factor C , and upstream mass density ρ_u is converted into CRe acceleration to produce the bolometric ($\geq \nu_0$) synchrotron luminosity of the relic

$$\int_{\nu_0} L(\nu) d\nu \simeq \frac{1}{2} A \rho_u V_{sh}^3 \eta_e \Psi \left(1 - \frac{1}{C^2}\right) \frac{B^2}{B_{cmb}^2 + B^2}, \quad (1)$$

where $\frac{B^2}{B_{cmb}^2 + B^2}$ takes into account that the energy of electrons is radiated away via both synchrotron and inverse Compton emission (B is the intensity of the magnetic field and $B_{cmb} = 3.25(1+z)^2 \mu\text{G}$ is the equivalent magnetic field of the cosmic microwave background at redshift z), the term

$$\Psi = \frac{\int_{p_0} N_{inj}(p) E dp}{\int_{p_{min}} N_{inj}(p) E dp} \quad (2)$$

accounts for the ratio of the energy flux injected in “all” electrons and those visible in the radio band ($\nu \geq \nu_0$), $N_{inj}(p)$ is the electron momentum distribution, $p_0 \simeq 10^3 m_e c \sqrt{\frac{\nu_0 [\text{MHz}]/(1+z)}{4.6 B [\mu\text{G}]}}$ is the momentum of the relativistic electrons emitting the synchrotron frequency ν_0 , and p_{min} is the minimum momentum of accelerated electrons.

We assume that a fraction η_e of the energy flux that can be dissipated at the shock surface¹ is channelled into the accelerated relativistic electrons in the downstream region F_{relic} :

$$\frac{1}{2} V_{sh}^3 \rho_u \left(1 - \frac{1}{C^2}\right) \eta_e = \overbrace{\frac{F_{relic}}{V_d \epsilon_{e,d}}}^{F_{relic}}, \quad (3)$$

¹ That is the change in the kinetic energy flux across the shock; e.g., Finoguenov et al. (2010).

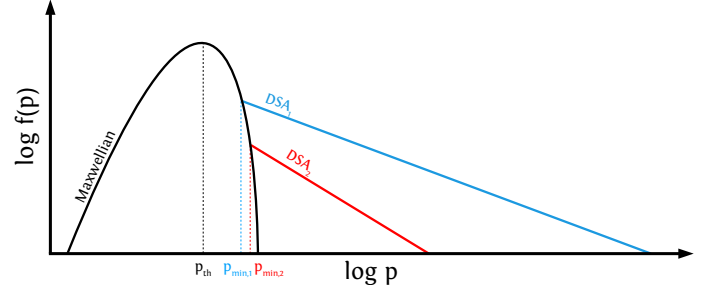


Fig. 1. Schematic representation of the electron momentum distribution in a downstream region. The two power laws show the DSA spectra (Eq. 9) in the case of two Mach numbers M_1 (blue) $>$ M_2 (red).

where

$$V_d = c_{s,u} \frac{M^2 + 3}{4M}, \quad (4)$$

which is the downstream velocity, $c_{s,u}$ is the upstream sound speed, and

$$\epsilon_{e,d} = \int_{p_{min}} N_{inj}(p) E dp, \quad (5)$$

which is the downstream energy density of the accelerated electrons. Although the acceleration efficiency can be computed for any possible electron momentum distribution, in the following we adopt a power-law momentum distribution of CRe in the form

$$N_{inj}(p) = K_e p^{-\delta_{inj}}, \quad (6)$$

which is motivated by the single power laws generally used to describe the integrated radio spectra of radio relics (e.g., van Weeren et al. 2019) and by DSA (see following section). If electrons are accelerated from the thermal pool starting from a minimum momentum as shown in Fig. 1, a relationship between the minimum momentum and the normalization of the spectrum can be derived by matching the number density of nonthermal electrons with that of thermal electrons with momentum p_{min} assuming $K_e p_{min}^{-\delta_{inj}}$. This leads to

$$K_e = \frac{4}{\sqrt{\pi}} n_d \frac{p_{min}^{2+\delta_{inj}}}{p_{th}^3} \exp \left[- \left(\frac{p_{min}}{p_{th}} \right)^2 \right], \quad (7)$$

where $p_{th} = \sqrt{2m_e k T_d}$ is the electron thermal peak momentum in the downstream gas (Fig. 1). At this point, it is important to know what kind of acceleration efficiency (or parameters p_{min} , K_e) is necessary to generate the observed radio properties of radio relics. To address this question we combine Eqs. 1, 3, and 7, and obtain

$$p_{min}^{2+\delta_{inj}} \exp \left[- \left(\frac{p_{min}}{p_{th}} \right)^2 \right] = \frac{\sqrt{\pi}}{4} \frac{\int_{\nu_0} L(\nu) d\nu}{n_d V_d A} \frac{p_{th}^3}{\int_{p_0} p^{-\delta_{inj}} E dp} \frac{B_{cmb}^2 + B^2}{B^2}, \quad (8)$$

which can be used to determine the minimum momentum of the electrons (or the efficiency) that is necessary to generate an observed radio emission for a given set of shock and relic parameters (namely Mach number, density, temperature, surface area,

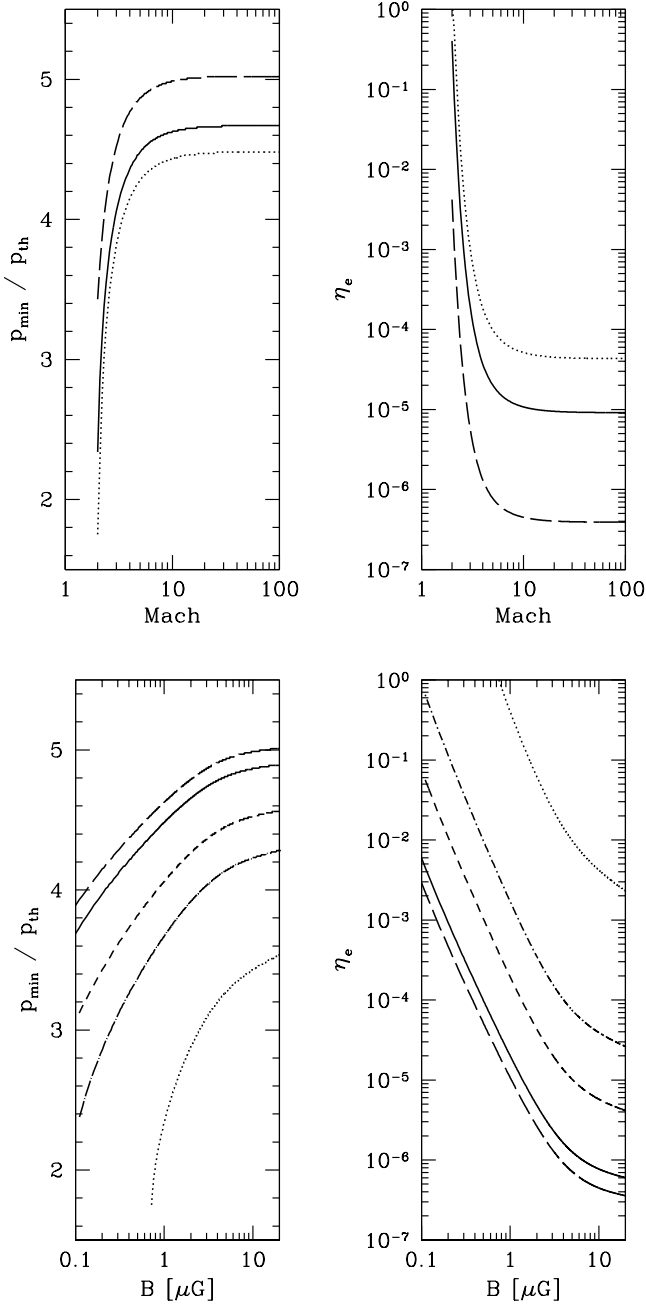


Fig. 2. Values of $p_{\min}/p_{\text{th}}$ and η_e for a mock radio relic (see text) at fixed magnetic field (top) and Mach number (bottom). In the top panels, curves denote different values of B : 0.5 μG (dotted), 1 μG (solid) and 10 μG (dashed). In the bottom panels, curves denote different values of M : 2 (dotted), 2.5 (dot-dashed), 3 (short dashed), 5 (solid) and 10 (long dashed).

and magnetic field). The surface of the shock is assumed to be $A = \pi R^2$, where R is the semi-axis of the relic emission crossed by the shock.

Our knowledge of B in clusters is poor and only a few constraints on the field strength in relics are available in the literature. In particular, the magnetic fields can be boosted at some level by shock compression and/or amplification in these dynamically active regions (Bonafede et al. 2013; Ji et al. 2016), perhaps reaching values up to 5 μG (e.g., van Weeren et al. 2010; Botteon et al. 2016b; Rajpurohit et al. 2018). This is important

to keep in mind because the required acceleration efficiency estimated with our approach is smaller for higher magnetic fields (see below).

2.2. Basic DSA relations

Among the possible mechanisms that can produce a power-law distribution of CRE, DSA is customarily used to describe particle acceleration at astrophysical shocks and radio relics. In this scenario, the slope δ_{inj} of the accelerated electrons (i.e., the injection spectrum) in Eq. 6 is

$$\delta_{inj} = 2 \frac{M^2 + 1}{M^2 - 1}, \quad (9)$$

and it depends only on the shock Mach number (e.g., Blandford & Eichler 1987).

Under stationary conditions and assuming that the physical conditions in the downstream regions do not change with distance from the shock, the electron spectrum integrated in the downstream region follows a power law with a slope of $\delta = \delta_{inj} + 1$. Thus, the integrated synchrotron spectrum is connected with the shock Mach number via

$$\alpha = \frac{M^2 + 1}{M^2 - 1} \equiv \alpha_{inj} + \frac{1}{2}. \quad (10)$$

As a consequence of the above relations, DSA predicts that for strong shocks ($M \rightarrow \infty$) the asymptotic behavior of the spectral index is $\alpha \rightarrow 1$ (and $\alpha_{inj} \rightarrow 0.5$), while for weak shocks ($M \lesssim 3 - 5$) it is $\alpha > 1$ (and $\alpha_{inj} > 0.5$).

Following the *thermal leakage injection* scenario for CRp (e.g., Gieseler et al. 2000), we assume here that only electrons with a minimum momentum threshold, $p_{\min}$, can be accelerated. For electron acceleration at weak shocks, the physical details that determine $p_{\min}$ are still poorly known (e.g., Guo et al. 2014; Kang et al. 2019); consequently, we use $p_{\min}$ as a free parameter that is connected with the efficiency (a larger $p_{\min}$ corresponds to lower efficiency). Following the framework described in Section 2.1 and the DSA relations, in Fig. 2 we report an example where we plot the $p_{\min}/p_{\text{th}}$ and η_e for a hypothetical radio relic at $z = 0.1$ with a favorable combination of $kT_d = 10$ keV, $n_d = 1.0 \times 10^{-3} \text{ cm}^{-3}$, $S_{1.4 \text{ GHz}} = 5$ mJy, and $A = 750^2 \pi \text{ kpc}^2$ for different values of the Mach number and the magnetic field strength. Figure 2 immediately identifies the problem: despite the optimistic parameters, these plots already demonstrate that DSA of thermal electrons becomes problematic (i.e., high η_e or large B are required) for weak shocks, that is, those of typically $M \lesssim 2 - 2.5$, and shocks of this strength are relatively common in the ICM and in radio relics. This is because, for weak shocks, an increasingly large fraction of the energy of the accelerated electrons is dumped into subGeV particles (formally, for $M < 2.2$, the majority of the energy is piled up into subrelativistic electrons).

3. Sample of radio relics with underlying shocks

We select a sample of ten radio relics with underlying shocks observed in the X-rays. The clusters are listed in Table 1 and include a few double radio relic systems. The sample is composed of well-studied radio relics with good radio and X-ray data available, which are essential to determine the spectral index of the relics and the properties of the underlying shocks. In particular,

Table 1. Sample of galaxy clusters with radio relics and detected underlying shocks. Reported values of M_{500} (mass within a radius that encloses a mean overdensity of 500 with respect to the critical density at the cluster redshift) and N_H (hydrogen column density) are taken from Planck Collaboration XXIX (2014) and Willingale et al. (2013), respectively. For the Sausage Cluster, M_{500} is from de Gasperin et al. (2014). Redshifts are taken from the NASA/IPAC Extragalactic Database (NED).

Cluster name	RA _{J2000} (h,m,s)	DEC _{J2000} (°,',")	M_{500} ($10^{14} M_\odot$)	z	N_H (10^{20} cm^{-2})	Instrument	Ref.	
							R	X
A2744	00 14 19	−30 23 22	9.56	0.308	1.46	<i>XMM-Newton</i>	1	2
A115	00 55 60	+26 22 41	7.20	0.197	6.38	<i>Chandra</i>	3	3
El Gordo	01 02 53	−49 15 19	8.80	0.870	1.78	<i>Chandra</i>	4	4
A521	04 54 09	−10 14 19	6.90	0.253	6.06	<i>Chandra</i>	5	6
A3376	06 01 45	−39 59 34	2.27	0.046	5.81	<i>XMM-Newton</i>	7	8
Toothbrush Cluster	06 03 13	+42 12 31	11.1	0.225	33.4	<i>Chandra</i>	9	10
Bullet Cluster	06 58 31	−55 56 49	12.4	0.296	6.43	<i>Chandra</i>	11	12
RXC J1314.4-2515	13 14 28	−25 15 41	6.15	0.247	8.91	<i>XMM-Newton</i>	13	14
A2146	15 56 09	+66 21 21	3.85	0.234	3.35	<i>Chandra</i>	15	16
Sausage Cluster	22 42 53	+53 01 05	7.97	0.192	46.4	<i>Chandra</i>	17	18

Notes. References (R=radio, X=X-ray): ¹Giacintucci et al. (in preparation); ²Eckert et al. (2016); ³Botteon et al. (2016a); ⁴Botteon et al. (2016b); ⁵Giacintucci et al. (2008); ⁶Bourdin et al. (2013); ⁷Kale et al. (2012); ⁸Urdampilleta et al. (2018); ⁹van Weeren et al. (2012); ¹⁰van Weeren et al. (2016); ¹¹Shimwell et al. (2014); ¹²Shimwell et al. (2015); ¹³Venturi et al. (2013); ¹⁴Mazzotta et al. (2011); ¹⁵Russell et al. (2012); ¹⁶Hoang et al. (2019); ¹⁷van Weeren et al. (2010); ¹⁸Akamatsu et al. (2015).

Table 2. Parameters used to compute the acceleration efficiency. Values of the downstream temperature kT_d and radio flux density S are derived from the works listed in Table 1.

Cluster name	Position	C	n_d	kT_d	A	S_ν	ν
			(cm^{-3})	(keV)	($\pi \text{ kpc}^2$)	(mJy)	(MHz)
A2744 [†]	NE	$1.90^{+0.60}_{-0.40}$	1.8×10^{-4}	12.3	740^2	20	1400
A115	N	$2.15^{+0.16}_{-0.14}$	1.5×10^{-3}	7.9	180^2	34	1400
El Gordo	NW	$2.88^{+0.30}_{-0.25}$	8.5×10^{-4}	17.9	350^2	28	610
A521	SE	$2.41^{+0.71}_{-1.41}$	3.0×10^{-4}	7.0	490^2	42	610
A3376	E	$1.98^{+0.27}_{-0.30}$	9.0×10^{-4}	4.7	260^2	40	1400
Toothbrush Cluster	N	$1.37^{+0.18}_{-0.17}$	5.5×10^{-4}	8.2	300^2	480	610
Bullet Cluster	E	$2.15^{+0.16}_{-0.14}$	5.0×10^{-4}	13.5	250^2	5	2100
RXC J1314.4-2515	W	$1.96^{+0.42}_{-0.36}$	1.0×10^{-3}	13.5	330^2	85	325
A2146	NW	$1.69^{+0.06}_{-0.06}$	3.5×10^{-3}	14.5	160^2	0.8	1500
Sausage Cluster	N	n.a.	3.0×10^{-4}	8.5	900^2	337	610

Notes. [†]Compression factor and downstream density taken from Eckert et al. (2016).

the detection of a shock co-spatially coincident with the relic (or at least a part of it) is necessary to evaluate the particle acceleration efficiency.

We point out that the well-known double radio relic system in A3667 (Röttgering et al. 1997; Johnston-Hollitt 2003) is not considered here because the measured spectral indexes of the two radio relics are ≤ 1 (Hindson et al. 2014; Riseley et al. 2015), and are therefore already in tension with DSA from the thermal pool (which would approach $\alpha = 1$ for very strong shocks; see Eq. 10).

We derive the relevant quantities reported in Table 2 that are necessary to compute the acceleration efficiency of electrons for all the shocks associated with the radio relics in our sample. In particular, we start from downstream quantities (i.e., temperature and density), usually better constrained by X-ray observations, and from the Mach number of the shock derived from the density jump to derive upstream quantities using the Rankine-Hugoniot jump conditions (Landau & Lifshitz 1959). Downstream temperatures are taken from the literature while downstream densities and compression factors are obtained from the re-analysis of the surface brightness profiles extracted across the shocks. Our broken power-law fit confirms previous analyses of the same tar-

gets, providing evidence of a discontinuity coincident with the outer edge of radio emission for nine out of ten relics in the sample. The only case where we use a single power-law model to fit the surface brightness profile is for the Sausage relic, which is known to not exhibit a surface brightness jump across its surface. Details of the analysis including X-ray images and sectors used to extract the surface brightness profiles are given in Appendix A.

4. Results

4.1. Comparison between Mach numbers and spectral indexes

In Table 3 and Fig. 3 we compare the Mach number measured from X-ray observations (M_X) and the relic spectral index estimated from radio observations (α_{radio}) with the expectations from DSA theory derived from Eq. 10 (where M_X and α_{radio} are used to derive α_{DSA} and M_{DSA} , respectively) for the relics in our sample. In the cases of El Gordo, A521, and the Bullet Cluster, the spectral indexes are consistent. However, there is a discrepancy between the observed spectral index of the relic and that implied by DSA in the majority of cases, confirming the findings

Table 3. Observed X-ray Mach number derived from the surface brightness analysis ($\mathcal{M}_X$) and integrated spectral index from the literature (α_{radio}). These were used to compute the expected integrated spectral index (α_{DSA}) and Mach number ($\mathcal{M}_{DSA}$) from DSA (Eq. 10). References for the integrated spectral indexes are also listed.

Cluster name	Position	$\mathcal{M}_X$	$\mathcal{M}_{DSA}$	α_{radio}	α_{DSA}	Reference
A2744	NE	$1.65^{+0.59}_{-0.31}$	$2.69^{+0.42}_{-0.27}$	$1.32^{+0.09}_{-0.09}$	$2.61^{+1.35}_{-0.66}$	Pearce et al. (2017)
A115	N	$1.87^{+0.16}_{-0.13}$	$4.58^{+0.27}_{-2.50}$	$1.10^{+0.50}_{-0.50}$	$1.80^{+0.19}_{-0.16}$	Govoni et al. (2001)
El Gordo	NW	$2.78^{+0.63}_{-0.38}$	$2.53^{+1.04}_{-0.41}$	$1.37^{+0.20}_{-0.20}$	$1.30^{+0.12}_{-0.11}$	Botteon et al. (2016b)
A521	SE	$2.13^{+1.13}_{-1.13}$	$2.33^{+0.05}_{-0.04}$	$1.45^{+0.02}_{-0.02}$	$1.57^{+0.36}_{-0.36}$	Macario et al. (2013)
A3376	E	$1.71^{+0.25}_{-0.24}$	$2.53^{+0.28}_{-0.20}$	$1.37^{+0.08}_{-0.08}$	$2.04^{+0.68}_{-0.34}$	George et al. (2015)
Toothbrush Cluster	N	$1.25^{+0.13}_{-0.12}$	$3.79^{+0.26}_{-0.22}$	$1.15^{+0.02}_{-0.02}$	$4.56^{+3.67}_{-1.34}$	Rajpurohit et al. (2018)
Bullet Cluster	E	$1.87^{+0.16}_{-0.13}$	$2.01^{+0.19}_{-0.14}$	$1.66^{+0.14}_{-0.14}$	$1.80^{+0.19}_{-0.16}$	Shimwell et al. (2015)
RXC J1314.4-2515	W	$1.70^{+0.40}_{-0.28}$	$3.18^{+0.87}_{-0.45}$	$1.22^{+0.09}_{-0.09}$	$2.06^{+0.91}_{-0.47}$	George et al. (2017)
A2146	NW	$1.48^{+0.05}_{-0.05}$	$3.91^{+1.95}_{-0.73}$	$1.14^{+0.08}_{-0.08}$	$2.68^{+0.23}_{-0.19}$	Hoang et al. (2019)
Sausage Cluster	N	n.a.	$4.38^{+1.06}_{-0.59}$	$1.11^{+0.04}_{-0.04}$	n.a.	Hoang et al. (2017)

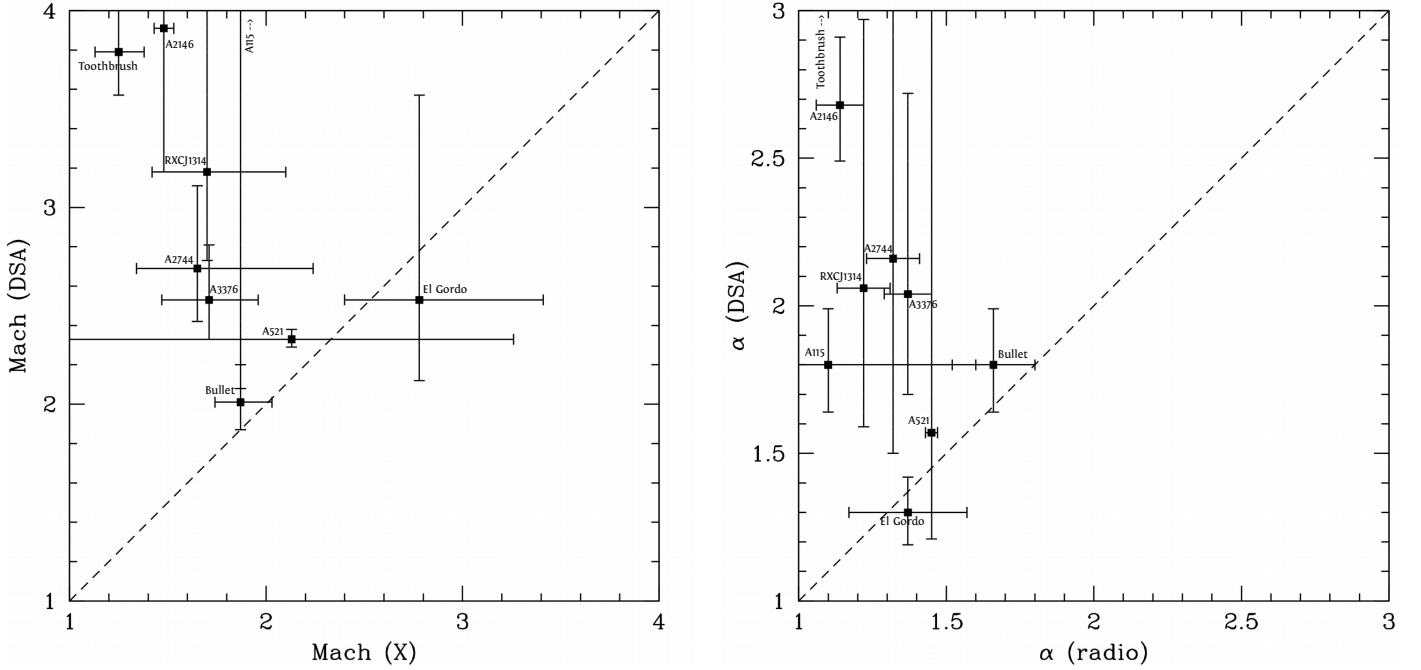


Fig. 3. Observed Mach numbers and spectral indexes vs. expected values from DSA theory. The values used to produce the plots are those listed in Table 3. The dashed lines indicate the linear correlation as a reference.

of previous studies that showed that the Mach numbers derived from radio observations under the assumption of DSA are generally biased high in comparison to those coming from X-ray data (e.g., Akamatsu et al. 2017; Urdampilleta et al. 2018). The inconsistency between radio and X-ray spectra is a long-standing problem of radio relics (e.g., Brunetti & Jones 2014). However, such tension might be understood by looking at numerical simulations which show that the inconsistency between radio- and X-ray-derived Mach numbers might emerge from projection effects of multiple shock surfaces (Skillman et al. 2013; Hong et al. 2015). Furthermore, modifications to the basic DSA theory (e.g., considering Alfvénic drift or including superdiffusion regimes at the shocks, Kang & Ryu 2018; Zimbardo & Perri 2018) change the expected value of the spectral index of the accelerated particles from DSA predictions.

4.2. Efficiencies

We calculate the acceleration efficiency of electrons η_e and the relevant parameters that are necessary to reproduce the bolomet-

ric synchrotron luminosity in the sample of relics presented in Section 3 as a function of the magnetic field B . From Eq. 1, the electron acceleration efficiency is

$$\eta_e \simeq \frac{2 \int_{\nu_0} L(\nu) d\nu}{A \rho_u V_{sh}^3} \left(1 - \frac{1}{C^2}\right)^{-1} \frac{1}{\Psi} \frac{B_{cmb}^2 + B^2}{B^2}, \quad (11)$$

where Ψ is given in Eq. 2 and depends on p_{min} . Since the Mach numbers observed in the X-rays and those predicted by DSA from the relic spectrum may be different (Fig. 3), we explore three approaches.

In the first two approaches, we assume standard DSA and use Eqs. 9 and 10 to connect the Mach number and the radio spectrum. The two approaches differ for the values of $\mathcal{M}$ that are assumed. In the first one (Section 4.2.1), we use $\mathcal{M}_X$ to calculate the spectrum of the accelerated particles. This approach is therefore based on X-ray observations. This scenario is at least relevant for those relics where the radio- and X-ray-derived Mach numbers are consistent (Fig. 3). In the second approach (Section 4.2.2), we assume that the average Mach number of

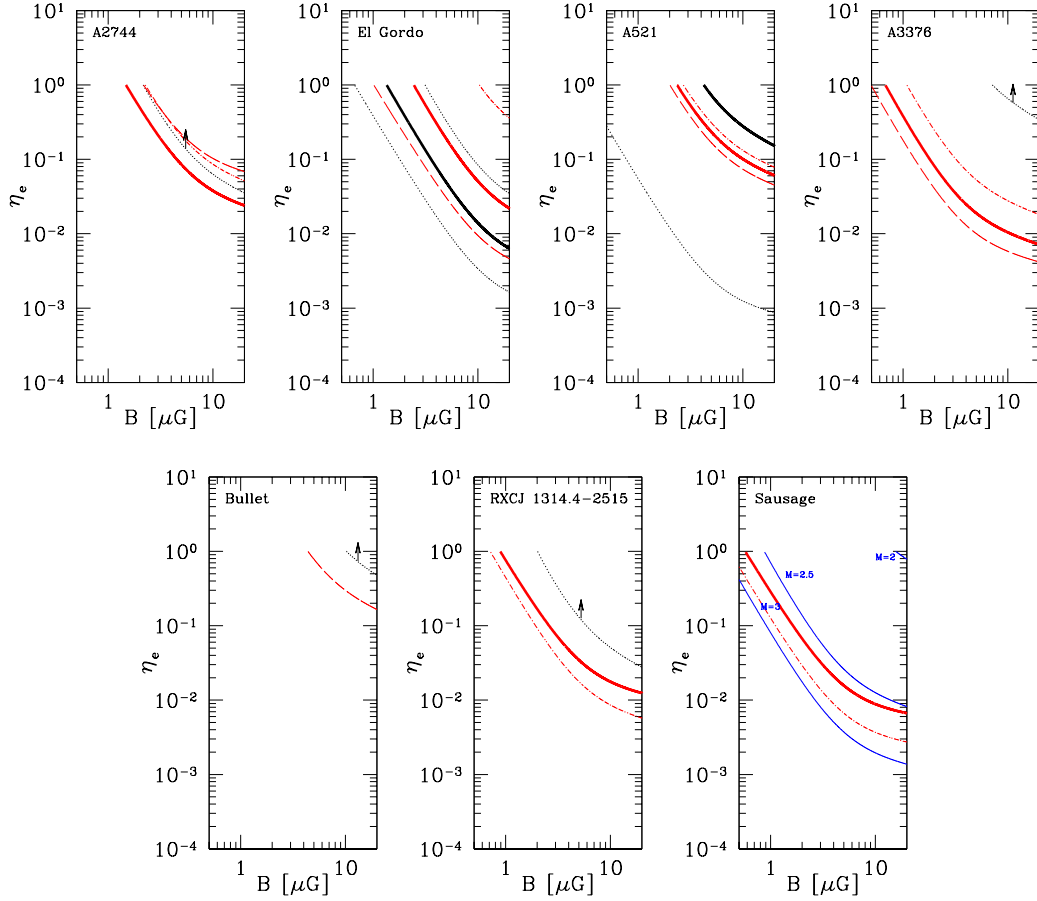


Fig. 4. Electron acceleration efficiency for the radio relics of the sample vs. magnetic field in the downstream region. Calculations were performed using the Mach numbers listed in Table 3 measured in the X-rays (*black*) and derived from the integrated spectral index in the case of DSA (*red*). Lines denote the best fit Mach numbers (*solid*), the X-ray upper and lower bounds (*dotted*), and the upper (*long dashed*) and lower (*dot-dashed*) bounds of $\mathcal{M}_{\text{DSA}}$. For the Sausage relic, lines in *blue* represent assumed Mach numbers of $\mathcal{M} = 2, 2.5$, and 3 (from top to bottom).

the shocks is given by $\mathcal{M}_X$ (and we therefore use $\mathcal{M}_X$ to get the energy flux at the shock; Eq. 3) and that the acceleration of electrons that generate radio relics is dominated by the sites in the shock surfaces where the Mach numbers $\geq \mathcal{M}_{\text{DSA}}$. This approach is motivated by numerical simulations (Hong et al. 2015; Roh et al. 2019) and might explain both radio and X-ray observations. In the third approach (Section 4.2.3), we go beyond the standard DSA model. We assume that a shock with Mach number $\mathcal{M}_X$ accelerates particles with a spectrum of $\delta = 2\alpha_{\text{radio}}$, where α_{radio} is the observed radio spectrum. This scenario allows us to match both radio and X-ray observations and is motivated by recent theoretical proposals (e.g., Zimbardo & Perri 2018). In all these approaches, we use the X-ray Mach numbers $\mathcal{M}_X$ estimated from the surface brightness analysis to calculate the energy flux at the shock surface (Eq. 3) because they are better constrained than those obtained with the spectral analysis.

4.2.1. Mach numbers measured in the X-rays

In Fig. 4 we show the electron acceleration efficiency versus magnetic field for the relics in the sample using the Mach number measured in the X-rays (black lines). We note that in most cases we report only a lower limit on η_e in the (B, η_e) plane because the observed radio luminosity cannot be matched even if the shock were to accelerate the entire distribution of thermal electrons (namely $p_{\text{min}} < p_{\text{th}}$ in Eq. 8, cf. Fig. 1). For A115, Toothbrush cluster, and A2146, no solution is found in Eq. 8 be-

cause the Mach numbers measured in the X-rays are too low and for this reason they are not reported in Fig. 4.

The acceleration efficiency of CRp for weak shocks is likely $< 1\%$ (e.g., Ryu et al. 2019), in order for γ -ray emission from clusters to be consistent with the *Fermi*-LAT upper limits (e.g., Vazza et al. 2016; Ha et al. 2019). Although the ratio of CRe/CRp acceleration is not measured for weak ICM shocks, we can reasonably assume that, as in the case of SNRs, the great majority of the energy flux at these shocks goes into CRp. For this reason, in the following we consider a conservative value of $\eta_e < 0.1$, which is generally associated to (protons in) strong SNR shocks ($\mathcal{M} \sim 10^3$), and a more realistic value of $\eta_e < 0.01$. In Fig. 5 we compute the acceleration efficiency that is requested to match the observed radio luminosity of the relics as a function of the shock Mach number measured in the X-rays at fixed downstream magnetic field of $B = 5 \mu\text{G}$; smaller magnetic fields will increase the requested value of η_e (see Eq. 11). For the Sausage relic, where no surface brightness jump is observed in the X-rays, we assume $\mathcal{M}_X = 2.5$ and $\mathcal{M}_X = 3$. For all the relics except A521 and El Gordo, it is possible to compute (i.e., $p_{\text{min}} > p_{\text{th}}$ in Eq. 8) the efficiency using only the upper bounds on the shock Mach number, and for this reason lower limits on the efficiency are reported in Fig. 5.

Looking at the distribution of acceleration efficiencies in Figs. 4 and 5, the model of DSA of thermal electrons is ruled out for the great majority of our relics. Even if we adopt an optimistic threshold of $\eta_e < 0.1$, the only two relics whose bolo-

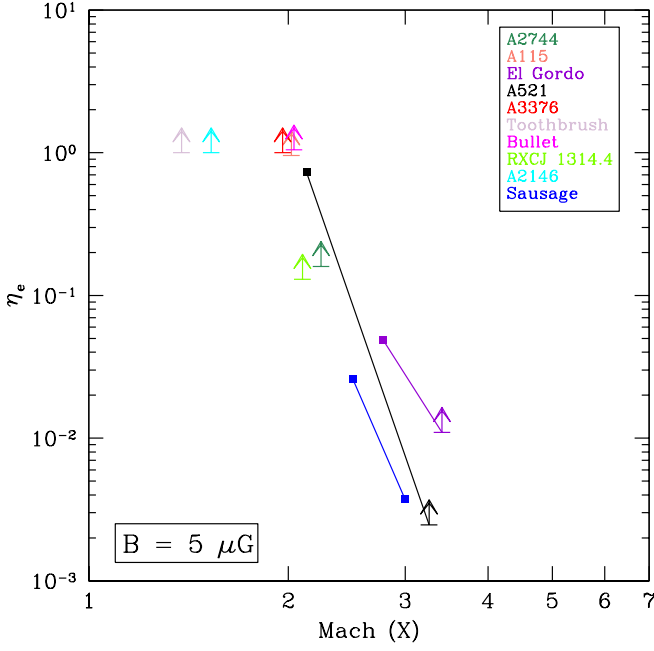


Fig. 5. Electron acceleration efficiency vs. M_X obtained for $B = 5 \mu\text{G}$. Most of the relics are reported as lower limits because it was only possible to compute the acceleration efficiency for the upper bounds of M_X . For the Sausage relic, we assumed $M_X = 2.5$ and 3 as no shock is observed in surface brightness.

metric radio luminosities can be reproduced by DSA from the thermal pool are those in El Gordo and A521, which are also those where M_X and M_{DSA} are consistent (Fig. 3); we note that this is only possible for A521 because of the large uncertainty of M_X , leading to an upper bound value which is significantly larger than M_{DSA} . In general, as already mentioned at the end of Section 2.1, Fig. 5 clearly shows the importance of the Mach number in the acceleration process: only $M \gtrsim 2.5$ have an energy flux and produce an accelerated spectrum that is sufficient to explain the luminosity of radio relics with the DSA of thermal ICM electrons.

The origin of radio relics can also be investigated by looking at possible correlations between the properties of shocks and the associated radio luminosity. Colafrancesco et al. (2017) studied a sample of radio relics and investigated the connection between relic radio power and shock Mach number. They interpreted the lack of correlation between these quantities as evidence against an origin of relics via DSA of thermal electrons. However, in general, the shock Mach number is not indicative of the energetics of the shock and M is only a fair measure of shock speed if the upstream temperature is the same for all the relics in a sample, and this is not the case. For this reason, we consider the kinetic energy flux across the surface of the shock, that is $1/2 A \rho_u V_{sh}^3$, which can be measured by X-ray observations (that provide upstream density, temperature, and Mach number) and radio observations (that provide the shock/relic surface area). We compare the radio luminosity and kinetic energy flux at the shocks for our radio relics but we do not find any clear correlation (Fig. 6). For similar kinetic energy flux at the shocks, we notice that the radio luminosity of relics differs by more than two orders of magnitude. This further suggests that DSA of thermal electrons is not the general mechanism acting for radio relics.

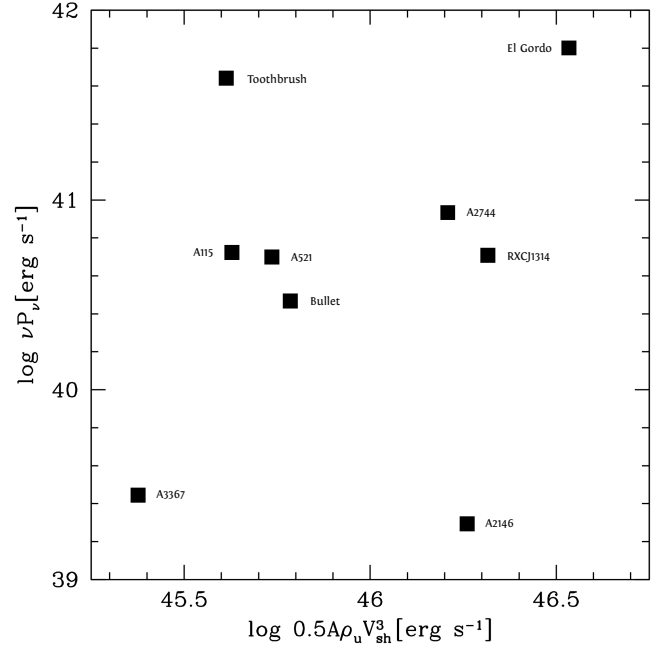


Fig. 6. Kinetic energy flux through shock surface vs. relic radio power. The values reported in Table 2 were used to compute the two quantities.

4.2.2. Matching X-ray and DSA Mach numbers with simulations

Numerical simulations show that Mach numbers implied by DSA from the radio spectra of relics in simulations are slightly larger than Mach numbers measured in the X-rays (e.g., Hong et al. 2015) because in simulations the Mach number is not constant across the shock surface and consequently the synchrotron emission from accelerated electrons is mainly contributed by regions where the Mach number is larger. For this reason, in this section we compute the electron acceleration efficiency for our relics by assuming the spectrum of the accelerated electrons and the Mach number implied by DSA from radio observations (Eq. 10). We use Eq. 11, but in this case the surface A is the area covered by large Mach numbers, $\geq M_{DSA}$.

The fraction f_{DSA} of the shock area that is covered by Mach numbers $\geq M_{DSA}$ is estimated with the simulations described in Roh et al. (2019). These simulations followed turbulence, magnetic fields, and shocks in the ICM using a model cluster in a controlled box to achieve a high resolution. With the data of the highest resolution (the size of grid zone $\Delta x = 3.9$ kpc) simulation, a number of mock radio relics are identified. We extract about 43000 grid zones, which together cover five shock surfaces of radio relics, and calculate the 3D Mach number of these cells, $M_{3D,sim}$ (see Roh et al. 2019). In Fig. 7 we show the distribution of $M_{3D,sim}$ normalized to the average Mach number of the shock surface to which the cells belong, $M_{X,sim}$. Following Hong et al. (2015), $M_{X,sim}$ is calculated from the temperature jump projected along the line of sight that is measured in simulations across the shock surfaces. In line with Hong et al. (2015), we find that shocked cells with a Mach number that is approximately equal to or greater than two times the X-ray Mach number of the shock are extremely rare. We determined f_{DSA} for our relics from the fraction of shocked cells with $M_{3D,sim}/M_{X,sim} \geq M_{DSA}/M_X$; the ratio M_{DSA}/M_X is derived from the values of M_{DSA} (with its error bounds) and the central value of M_X reported in Table 3. The resulting f_{DSA} are reported in Table 4.

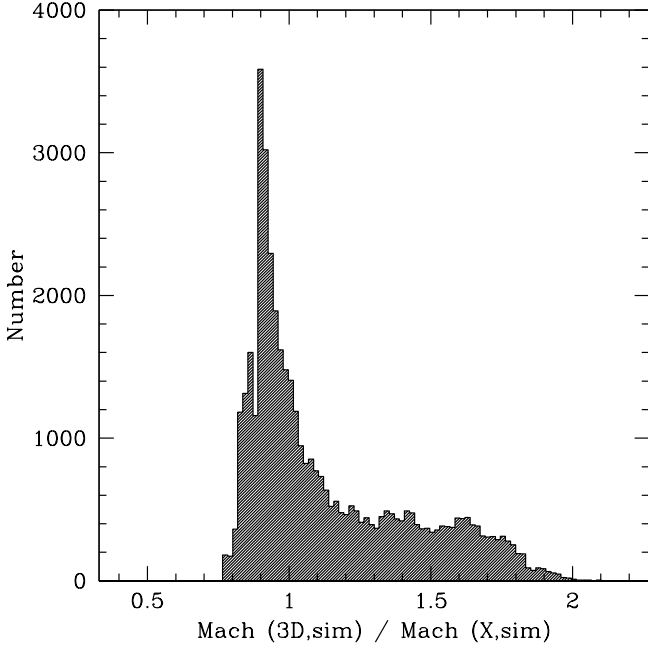


Fig. 7. Distribution of the ratio between the 3D Mach number $M_{3D,sim}$ and the average X-ray Mach number $M_{X,sim}$ for five radio relics extracted from our simulations.

Table 4. Values of f_{DSA} taking into accounts the error bounds of M_{DSA} and the best fit of M_X for the relics in our sample. For the Sausage Cluster we assumed $M_X = 2.5$.

Cluster name	f_{DSA}^-	f_{DSA}	f_{DSA}^+
A2744	0.1783	0.0985	0.0099
El Gordo	1.000	0.7788	0.2901
A521	0.4392	0.4195	0.4016
A3376	0.2405	0.1699	0.0882
Bullet Cluster	0.4202	0.4392	0.3450
RXC J1314.4-2515	0.1086	0.0117	–
Sausage Cluster	0.1535	0.0420	–

We use Eq. 11 to calculate the acceleration efficiency that is necessary to match the observed synchrotron luminosity assuming a surface $A' = A \times f_{DSA}$ and M_{DSA} (cf. Tab. 3). The electron acceleration efficiency versus magnetic field for the relics in the sample computed using the Mach numbers implied by DSA from the relic integrated spectral indexes are shown in Fig. 4 (red lines). As before, in Fig. 8 we summarize the results by showing the acceleration efficiency that is required to match the observed radio luminosity of the relics as a function of the shock Mach number derived under DSA assumption at a fixed downstream magnetic field of $B = 5 \mu\text{G}$.

Even if we consider an optimistic acceleration efficiency of $\eta_e < 0.1$, we find that only four out of ten relics (namely: A2744, El Gordo, A3376, and RXCJ 1314.4-2515) can be explained via DSA assuming an optimistic value of $B > 5 \mu\text{G}$. On the other hand, an efficiency $\eta_e < 0.01$ is obtained only for A3376 and only by considering $M_{DSA} = 3.18$ (i.e., its upper bound on M_{DSA}) and $B > 5 \mu\text{G}$. Therefore, also in this case, we conclude that even if this approach matches the synchrotron spectral index and the average Mach number measured in the X-rays, it energetically fails to reproduce the synchrotron luminosity observed in the great majority of our radio relics.

We note that the correction factor and the efficiency obtained

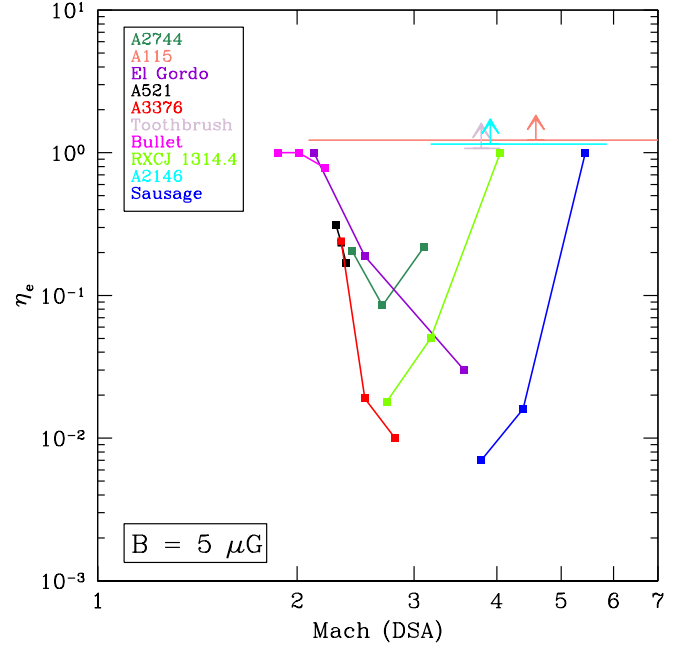


Fig. 8. Electron acceleration efficiency vs. M_{DSA} obtained for $B = 5 \mu\text{G}$. For the Sausage relic, f_{DSA} was calculated by assuming $M_X = 2.5$. Upper limits at $\eta_e > 1$ are slightly offset for display purposes.

with this procedure are anchored to the Mach number measured in the X-rays for the relics (Table 3). Although there are possible systematic errors involved in measuring Mach numbers in the X-rays (e.g., Markevitch & Vikhlinin 2007), we believe that these Mach numbers for our relics are relatively reliable for two main reasons: (i) the Mach numbers from density and temperature jumps (that have very different dependencies on M) are generally consistent for our relics, and (ii) the presence of double relic systems in our sample suggests that at least in a large fraction of our relics the shocks propagate close to the plane of the sky, thus minimizing projection effects.

4.2.3. Beyond standard DSA

Another *ad hoc* possibility is shock acceleration from the thermal pool where the spectrum differs from the classical expectations based on the DSA theory. To explore this possibility, we evaluate the acceleration efficiencies assuming the values of Mach number and radio spectrum measured from X-ray and radio observations, respectively. Results are shown in Fig. 9; the acceleration efficiencies are similar to those computed assuming M_{DSA} and $f_{DSA} = 1$. We find that such a modified DSA model can explain the properties of the radio relics in our sample and their connection with the underlying shocks observed with $B = 5 \mu\text{G}$ and $\eta_e < 0.1$ except for the case of El Gordo, A521, and the Bullet Cluster; this latter is the most critical case due to the very steep synchrotron spectrum of the relic. The point here is to understand why clusters shocks, contrary to shocks in SNRs, produce a spectrum that is different from that implied by DSA. One possibility that is worth mentioning is that the transport of electrons across the shock is not diffusive but is superdiffusive, as recently proposed by (Zimbardo & Perri 2018) for example. However, at this stage this hypothesis is rather speculative.

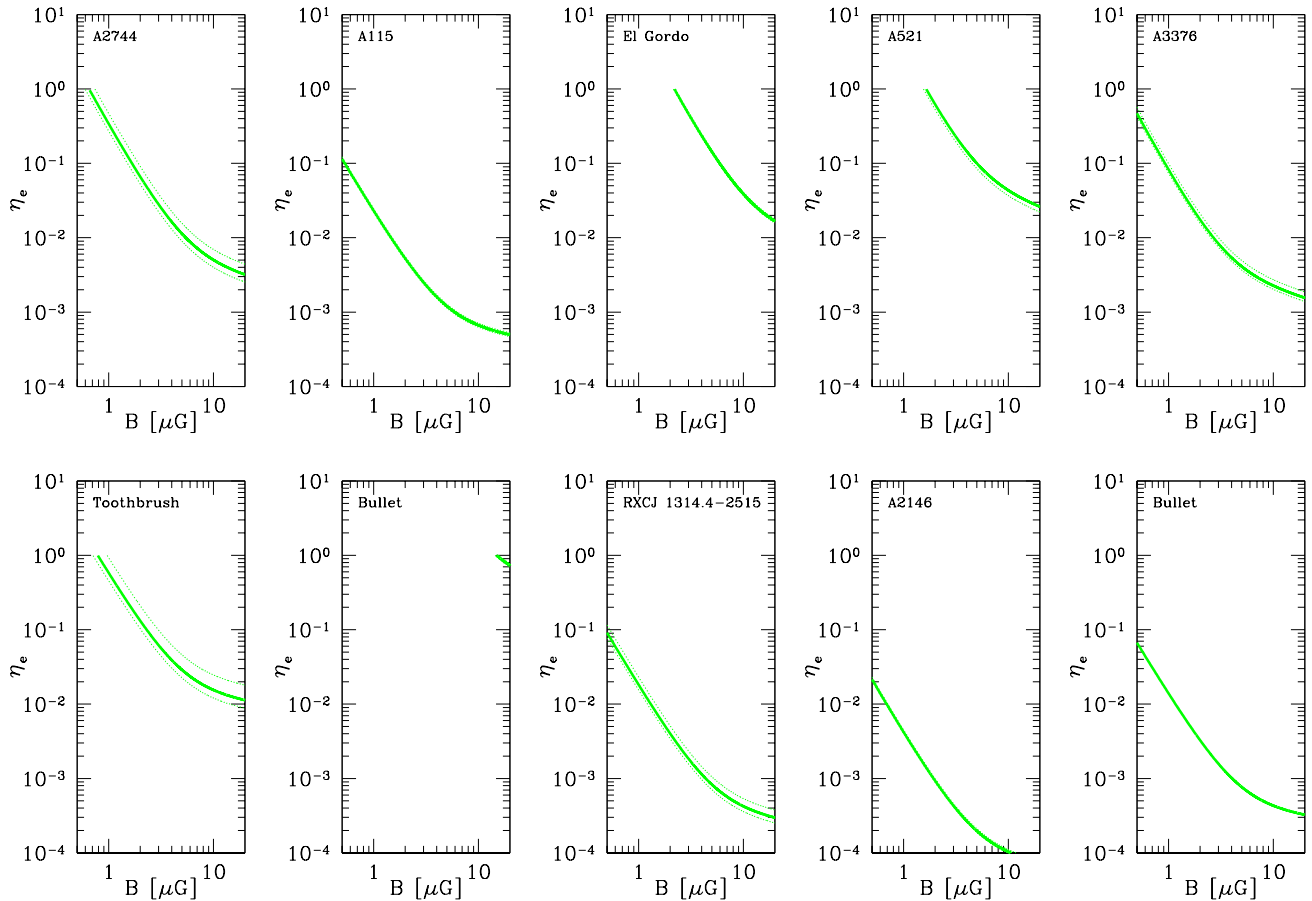


Fig. 9. Same as Fig. 4 where now the green lines denote the case where the electron acceleration efficiency was computed adopting the Mach number measured in the X-rays (with its error bounds) and the central value of integrated spectral index measured in radio (cf. Table 3).

5. Conclusions

We selected ten well-studied radio relics with detected underlying shocks to study the electron acceleration efficiency that is required to explain the origin of the diffuse radio emission. Our aim was to test whether or not DSA from the thermal pool is a general mechanism capable of explaining radio relics. Under this condition, the accelerated electrons are injected with a power-law momentum distribution from a minimum momentum p_{min} . Theoretically, this momentum is difficult to determine, and therefore we adopted the reverse approach to calculate p_{min} and efficiency that would be necessary to explain the observed radio assuming DSA and shock parameters.

In general, the slope of the momentum distribution of the emitting electrons in our relics does not match DSA prediction if we assume the Mach number of the underlying shocks; this confirms previous findings (e.g., Akamatsu et al. 2017; Urdampilleta et al. 2018). As a consequence we adopted two approaches to test DSA of thermal particles. Firstly, we assumed the Mach number measured in the X-rays: this model is based on the X-ray measurements but it only matches the radio spectrum of three relics in our sample. We find that this Mach number is generally too low – typically $M_X \lesssim 2 - 2.5$ – to reproduce the observed synchrotron luminosity of our radio relics, pinpointing a fundamental energy argument against this scenario. Under the optimistic assumption of a large B , only A521 and El Gordo can be explained by this mechanism assuming $\eta_e < 0.1$. We also show that the kinetic energy flux of the shocks constrained by the X-ray observations of our relics does not correlate with the

observed radio luminosity, further disfavoring a direct connection between DSA of thermal electrons and radio relics.

Secondly, we assumed the value of the shock Mach number measured in X-rays but also assumed that the Mach number across the shock surface is not uniform and that the acceleration of electrons is dominated by the regions where M is higher. More specifically, we assumed that in these cases the acceleration is associated with Mach numbers that are implied by DSA from the spectrum of the radio relics. These Mach numbers for our relics are significantly larger than those measured in the X-rays and we used numerical simulations to correct for the fraction of shock surface that is occupied by these Mach numbers, f_{DSA} . This fraction is small, implying that the acceleration efficiency that is required to match the observed radio luminosity is still relatively large in these regions. Under the optimistic conditions of large B , only A3376 can be explained under the assumption of $\eta_e < 0.01$, whereas the other three relics (A2744, El Gordo, RXC J1314.4-2515) can be explained if $\eta_e < 0.1$.

Our main conclusion is therefore that, even if models of DSA can be reconciled with radio and X-ray measurements of Mach numbers, in general they fail to reproduce the observed synchrotron luminosity of radio relics due to the small amount of energy dissipated into particle acceleration at cluster shocks.

For this reason we also investigated a “hybrid” case where the acceleration efficiency is evaluated assuming the Mach number and synchrotron spectrum derived from X-ray and radio observations. In this case the majority of radio relics in our sample can be explained in the case of high B and assuming $\eta_e < 0.1$, making this solution very appealing. Nonetheless, invoking a

scenario that produces synchrotron spectra that differ from DSA predictions is rather speculative at present and requires further theoretical follow-up studies.

The main conclusion of the present study is that other mechanisms, such as shock re-acceleration of supra-thermal seed electrons or a modification of standard DSA, are required to explain the formation of radio relics. The assumption of mildly relativistic electrons (i.e., higher p_{min}) would lower the required acceleration efficiencies and explain the luminosity of relics at weak shocks through re-acceleration mechanisms, as we discussed for A115 in an earlier work (Botteon et al. 2016a).

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Appendix A: Details of the X-ray analysis

We retrieved all the ObsIDs available for the clusters of our sample from the *Chandra*² and *XMM-Newton*³ archives. In the cases where the clusters have been observed with both instruments, we used *Chandra* because its higher angular resolution allows us to better characterize the sharp edges of the shocks and to excise the point sources more accurately (cf. Table 1).

We performed a standard data reduction using CIAO v4.10, *Chandra* CALDB v4.7.8 and the ESAS tools developed within the *XMM-Newton* SAS v16.1.0. After the excision of contaminating point sources, surface brightness profiles across the shocks were extracted from the 0.5 – 2.0 keV exposure-corrected images of the clusters and fitted with PROFIT v1.5 (Eckert et al. 2011). An underlying broken power law with a density jump was assumed to fit the data, which were rebinned to reach a minimum S/N of 7. PROFIT performs a modeling of the 3D density profile that is numerically projected along the line of sight under spherical assumption (following the Appendix in Owers et al. 2009). Deprojected density profiles were recovered from the emission measure of the plasma evaluated in the case of an absorbed APEC model (Smith et al. 2001) with metallicity assumed to be 0.3 $Z_{\odot}$ and total (i.e., atomic + molecular) hydrogen column density measured in the direction of the clusters fixed to the values of Willingale et al. (2013). The choice of the soft band 0.5 – 2.0 keV ensures that the bremsstrahlung emissivity is almost independent of the gas temperature for $kT \gtrsim 3$ keV (e.g., Ettori et al. 2013).

For the Sausage relic, no surface brightness discontinuity has been detected in the X-rays (Ogrean et al. 2013, 2014). Indeed, this is the only case where we used a single power-law model to fit the surface brightness profile. In this case, we used the density measured at the location of the relic from the single power-law model and assumed different Mach numbers in the analysis.

In Fig. A.1 we show the *Chandra* and *XMM-Newton* images in the 0.5 – 2.0 keV band of the clusters in the sample while in Fig. A.2 we report the X-ray surface brightness profiles extracted across the relics considered in the analysis.

² <http://cda.harvard.edu/chaser/>

³ <http://nxsa.esac.esa.int/>

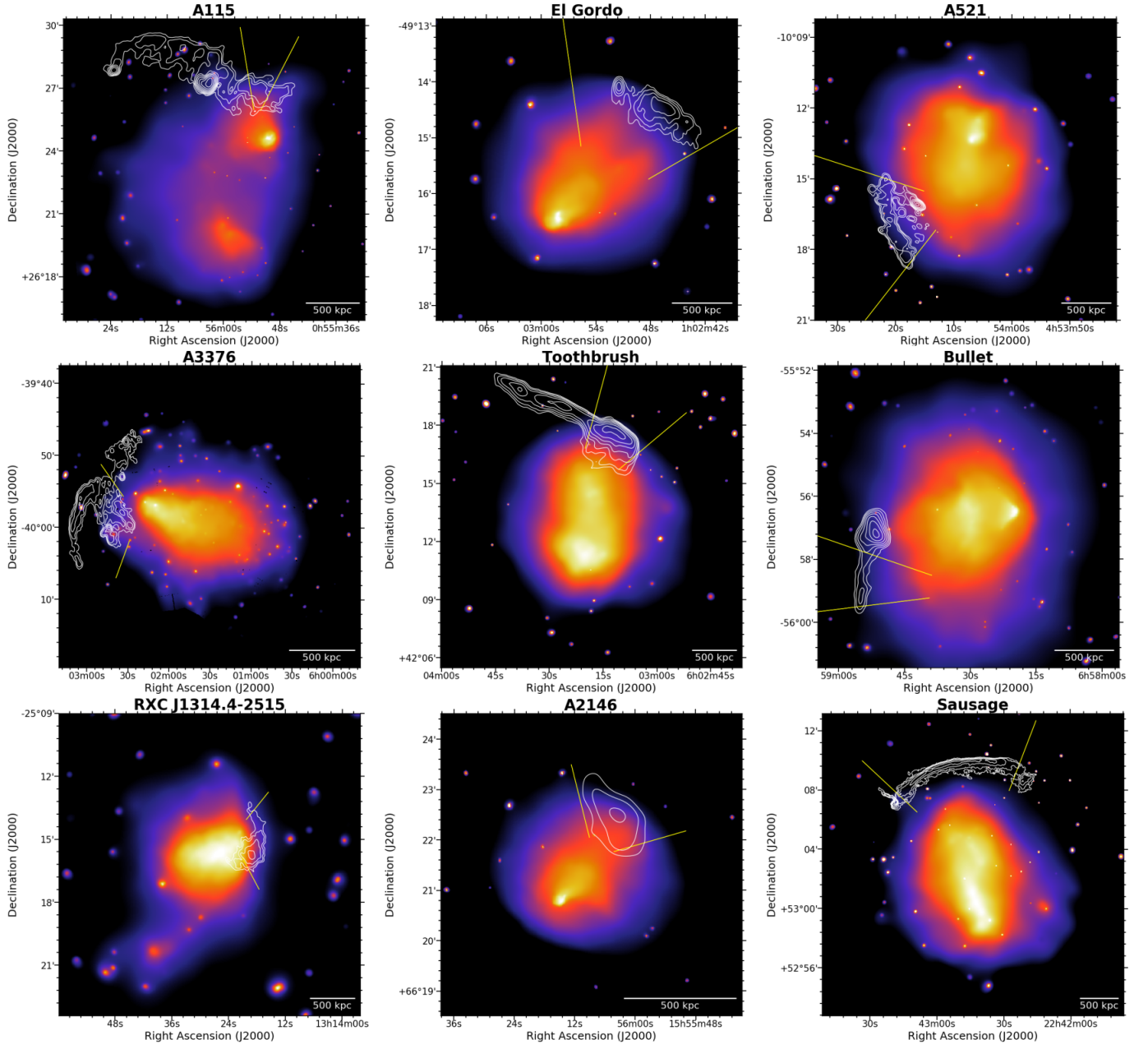


Fig. A.1. X-ray images with overlaid radio contours from the works listed in Table 1 and the sectors used for the analysis of the shock fronts.

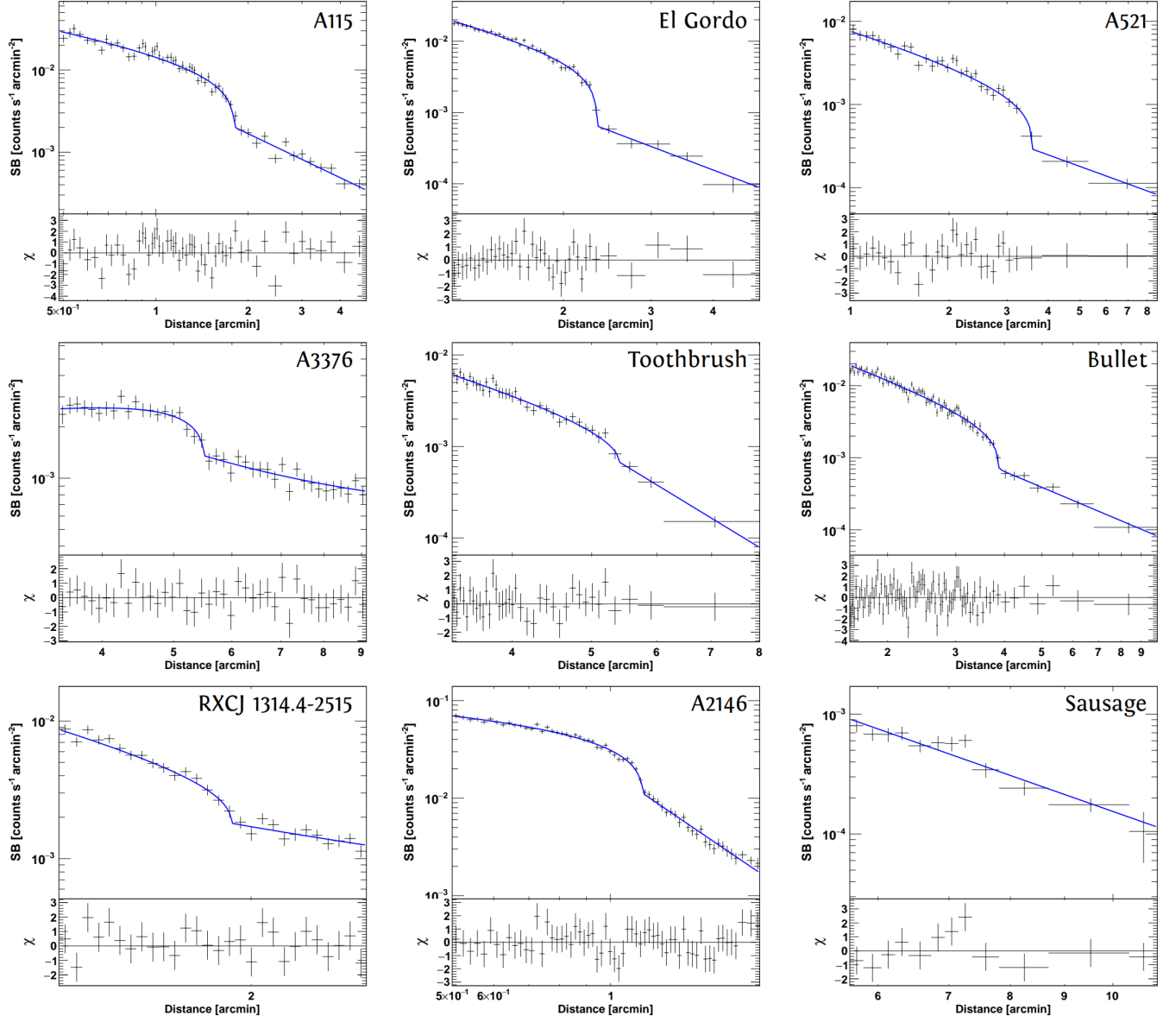


Fig. A.2. X-ray surface brightness profiles extracted across the relics in the sample. A broken power-law model was used to fit the data for all the relics but the Sausage.