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Feedback effect on the observable properties of $z > 6$ AGN

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ABSTRACT

Active galactic nuclei (AGNs) feedback has a major impact on to the supermassive black hole (SMBH) growth, the properties of the host galaxies, and their cosmic evolution. We investigate the effects of different kinetic feedback prescriptions on the observable properties of AGN and their host galaxies at $z > 6$ in a suite of zoom-in cosmological simulations. We find that kinetic feedback decreases the column density of the interstellar medium (ISM) in the host galaxy by up to a factor of ≈ 10 , especially when the SMBHs reach high accretion rates ($\approx 10\text{--}30\text{ M}_{\odot}\text{ yr}^{-1}$). In particular, kinetic feedback is required to extend the ISM size to > 1 kpc and match the observed sizes of the gas reservoirs in $z > 6$ AGN host galaxies. Moreover, it produces unobscured lines of sight along which the AGN can be detected in the rest-frame UV band with magnitudes consistent with observed values of $z > 6$ AGN. The assumed geometry of the outflow plays an important role in shaping the observed properties of high-redshift AGN. We find that a biconical geometry is favoured over a spherical one to reproduce the observed properties, but it overestimates the number of multiple AGN systems detectable in X-ray observations. This result suggests that simplistic BH seeding recipes widely employed in cosmological simulations produce too many X-ray detectable multiple AGN at $z = 6\text{--}7$, thus soliciting the adoption of more physically motivated seeding prescriptions.

Key words: methods: numerical – galaxies: active – galaxies: high-redshift – galaxies: nuclei – early Universe.

1 INTRODUCTION

The existence of supermassive black holes (SMBHs) with masses larger than billion solar masses at $z \gtrsim 6$ (e.g. Bañados et al. 2018; Matsuoka et al. 2018a; Wang et al. 2021a), when the Universe was < 1 Gyr old, challenges our current understanding of SMBH and galaxy formation and evolution, and is thus one of the most pressing open issues in modern astrophysics (e.g. Woods et al. 2019). Their distance and faintness make observations of these objects difficult and strongly biased towards the most luminous and massive accreting SMBHs. A complementary approach is to use numerical simulations as tools to study the largely unknown phases of SMBH growth in the early Universe (e.g. Sijacki, Springel & Haehnelt 2009; Tanaka & Haiman 2009; Habouzit et al. 2016, 2019). However, observed properties of high-redshift accreting SMBHs, or active galactic nuclei (AGN), and predictions of numerical simulations have been compared only seldom (e.g. Ni et al. 2020; Habouzit et al. 2022; Di Mascia et al. 2021a; Zana et al. 2022).

An important ingredient entering in numerical simulations focused on the early growth of SMBHs is the effect of AGN feedback (e.g. Costa et al. 2014; B18; Habouzit et al. 2019; Costa, Pakmor & Springel 2020; Valentini, Gallerani & Ferrara 2021), as it is often considered to have a major role in shaping the evolution of AGN and galaxies along the whole cosmic history (e.g. Fiore et al. 2017). In particular, optically selected luminous quasi-stellar objects (QSOs) in the early Universe often present evidence for the launching of fast and massive multiphase outflows (e.g. Maiolino et al. 2012; Ciccone et al. 2015; Bischetti et al. 2019; Carniani et al. 2019; Schindler

et al. 2020; Izumi et al. 2021; but see also, e.g. Decarli et al. 2018; Novak et al. 2020; Meyer et al. 2022), which are expected to affect the observable properties of the QSOs themselves and their host galaxies, such as X-ray obscuration, UV extinction, and gas content (e.g. Brusa et al. 2015; Ni et al. 2020).

Outflows observed in QSOs are thought to originate from fast nuclear winds, which, in turn, may be accelerated by several physical mechanisms, including radiation pressure, due to UV photons produced in the accretion disc, on dust grains or on partially ionized gas mediated by UV transitions, and magnetic effects (e.g. Proga, Stone & Kallman 2000; Murray, Quataert & Thompson 2005; Fabian, Vasudevan & Gandhi 2008; Yuan et al. 2015; Ricci et al. 2017). The physical scales involved in these processes are those of the accretion disc (e.g. Giustini & Proga 2019). Since such scales cannot be resolved by large-scale cosmological simulations, different authors have modelled AGN feedback using several different recipes (e.g. Barai et al. 2018; Costa et al. 2020; Ni et al. 2020). Moreover, the effect of the outflow on the surrounding material can potentially depend on its geometry (e.g. Zubovas, Bourne & Nayakshin 2016). Since the exact acceleration physics, and thus launching direction, of nuclear winds is not well understood, numerical simulations typically assume either spherical (e.g. Feng et al. 2016) or bi-conical (e.g. Sala et al. 2021) outflow geometry as study cases.

Beside the properties of the individual galaxies hosting accreting SMBHs, numerical simulations provide also information on the environment of high-redshift luminous AGN. While these objects are expected to reside in the peaks of the dark matter halo distribution, which are generally characterized by large overdensities of galaxies (e.g. Costa et al. 2014; Wise et al. 2019), although with some scatter (e.g. Habouzit et al. 2019), observations struggle to provide

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us with a clear view of typical high-redshift QSO environment. In fact, $z > 6$ QSOs have been reported to reside in a variety of environments, including underdense, normal, and overdense regions (e.g. Ota et al. 2018; Mazzucchelli et al. 2019; Overzier 2022). The first spectroscopically confirmed galaxy overdensity around a $z > 6$ QSOs was presented recently by Mignoli et al. (2020), followed by a tentative confirmation of another structure by Overzier (2022).

A significant fraction (≈ 40 per cent) of $z \gtrsim 6$ QSOs has ALMA-detected dusty companion galaxies at distances of a few kpc (e.g. Willott, Bergeron & Omont 2017; Decarli et al. 2018; Neeleman et al. 2019; Venemans et al. 2020). These satellite galaxies might host heavily reddened and buried AGN (e.g. Di Mascia et al. 2021a), although currently there is no strong observational evidence for the presence of accreting SMBHs in their centres (e.g. Connor et al. 2019, 2020; Vito et al. 2019a, 2021). Such objects would be typically brighter than inactive galaxies, especially in the X-ray band. Therefore, their predicted number in numerical simulations can be tested against observational results to infer how well simulations approximate reality.

In this paper, we present a study of the effect of AGN kinetic feedback on the observable properties of $z > 6$ AGN in cosmological simulations. In particular, we analyse a set of numerical simulations presented by Barai et al. (2018, hereafter, B18) with different kinetic feedback prescriptions, focusing on the most massive SMBH at $z = 6$ and its surrounding environment. We extract multiwavelength observables such as column density and radial extent of the gas distributed in the host galaxies, UV and X-ray AGN fluxes, and number of satellite AGN detectable over small (i.e. a few kpc) distances from the central SMBH. We compare these properties with results from multiwavelength observations.

The paper is structured as follows. In Section 2, we describe the numerical setup of the simulations, the AGN selection, and the method used to measure the gas column density and distribution. In Section 3, we discuss the redshift evolution of the column densities for the considered AGN. In Section 4, we present the observable properties of the simulated AGN and their host galaxies, and we compare them with empirical findings. In Section 5, we investigate the presence of multiple AGN systems over scales of a few kpc, and we compare their detectability rates in the X-ray band with results from observations of high-redshift AGN. Finally, in Section 6, we discuss and interpret the results, and in Section 7 we provide a summary. All quoted distances are physical unless otherwise noticed. We adopt a flat Λ CDM cosmology with $H_0 = 67.7 \text{ km s}^{-1}$ and $\Omega_m = 0.307$ (Planck Collaboration et al. 2016).

2 METHOD

2.1 Numerical model

We consider the simulation runs *AGNcone* and *AGNsphere* by B18, which include kinetic feedback. We provide here a summary of the numerical setup and refer to the original works for an in-depth discussion.

B18 used a modified version of the smooth particle hydrodynamics (SPH) N-body code GADGET-3 (Springel 2005) to follow the evolution of a comoving volume of $(500 \text{ Mpc})^3$, starting from cosmological initial condition generated with MUSIC (Hahn & Abel 2011) at $z = 100$, and zooming-in on the most massive (i.e. $4 \times 10^{12} M_\odot$) dark matter (DM) halo, corresponding to a $\approx 3\sigma$ overdensity (e.g. Barkana & Loeb 2001), inside the box down to $z = 6$. Therefore, the final zoomed-in simulations focus by construction on a highly biased cubic region, with a volume of $(5.21 \text{ Mpc})^3$. The highest level

of the simulation has a mass resolution of $m_{\text{DM}} = 7.54 \times 10^6 M_\odot$ and $m_{\text{gas}} = 1.41 \times 10^6 M_\odot$ for DM and gas particles, respectively. The softening length for gravitational forces for these high-resolution DM and gas particles is $R_{\text{soft}} = 1 h^{-1} \text{ ckpc}$.

The code accounts for gas heating and cooling (including metal-line cooling) depending on the gas metal content, based on eleven element species (H, He, C, Ca, O, N, Ne, Mg, S, Si, Fe) that are tracked in the simulation (Tornatore et al. 2007). Star formation in the interstellar medium (ISM) is implemented following the multiphase effective subresolution model by Springel & Hernquist (2003), adopting a density threshold for star formation of $n_{\text{SF}} = 0.13 \text{ cm}^{-3}$. The simulations include stellar winds, supernovae feedback, and metal enrichment, and assume a Chabrier (2003) initial mass function in the mass range $0.1\text{--}100 M_\odot$ (Tornatore et al. 2007; Barai et al. 2013; Biffi et al. 2016).

When a DM halo that is not already hosting a black hole (BH) reaches a total mass of $M_h \geq 10^9 M_\odot$, a $M_{\text{BH}} = 10^5 M_\odot$ BH is seeded at its centre. BHs are treated as collisionless sink particles and are allowed to grow by accretion of the surrounding gas or by mergers with other BHs. Gas accretion on to BHs is modelled via the classical Bondi–Hoyle–Littleton accretion rate $\dot{M}_{\text{Bondi}}$ (Hoyle & Lyttleton 1939; Bondi & Hoyle 1944; Bondi 1952), capped at the Eddington rate $\dot{M}_{\text{Edd}}$:

$$\dot{M}_{\text{BH}} = \min(\dot{M}_{\text{Bondi}}, \dot{M}_{\text{Edd}}). \quad (1)$$

Accreting BH radiate away a fraction ϵ_r of the accreted rest-mass energy, with a bolometric luminosity

$$L_{\text{bol}} = \epsilon_r \dot{M}_{\text{BH}} c^2, \quad (2)$$

where c is the speed of light. B18 fixed the radiative efficiency to $\epsilon_r = 0.1$, a fiducial value for radiatively efficient, geometrically thin, optically thick accretion discs around a Schwarzschild BH (Shakura & Sunyaev 1973).

A fraction $\epsilon_f = 0.05$ of the total output energy is distributed to the surrounding gas in a kinetic form.¹ In *AGNcone*, the kinetic energy is distributed along two cones with a half-opening angle of 45° . The direction of the cone axis is chosen randomly for each BH at the seeding time, and is kept fixed throughout the simulation (B18), similarly to what is done in Zubovas et al. (2016). Instead, the AGN feedback in *AGNsphere* pushes away the gas particles along random directions, thus mimicking a spherical geometry.

2.2 AGN selection

We analyse the simulation snapshots in steps of $\Delta z = 0.2$ from $z = 10$ to 8 and $\Delta z = 0.1$ from $z = 8$ to 6. In particular, we follow the most massive SMBH at $z = 6$ in each simulation set, and consider a box with side size of 60 kpc centred on it. We refer to all of the SMBHs in the box accreting at $\dot{M}_{\text{BH}} > 0.02 M_\odot \text{ yr}^{-1}$ (i.e. $L_{\text{bol}} \approx 10^{44} \text{ erg s}^{-1}$) as AGN. Fig. 1 presents the BH mass evolution of AGN in the two simulations. Each AGN is labelled with the initial letter of the run (C for *AGNcone*, S for *AGNsphere*).

AGNcone forms two very massive ($> 10^9 M_\odot$) BHs at $z < 7$, while only less massive BHs are formed in the *AGNsphere* run. This behaviour is linked to the implementation of the feedback: *AGNcone* allows the gas to accrete continuously along the equatorial directions, while the lack of a preferential direction along which the outflow is launched in *AGNsphere* does not allow for a steady and

¹We refer to B18 for details about the choice of the value for ϵ_f and the numerical implementation of the kinetic feedback.

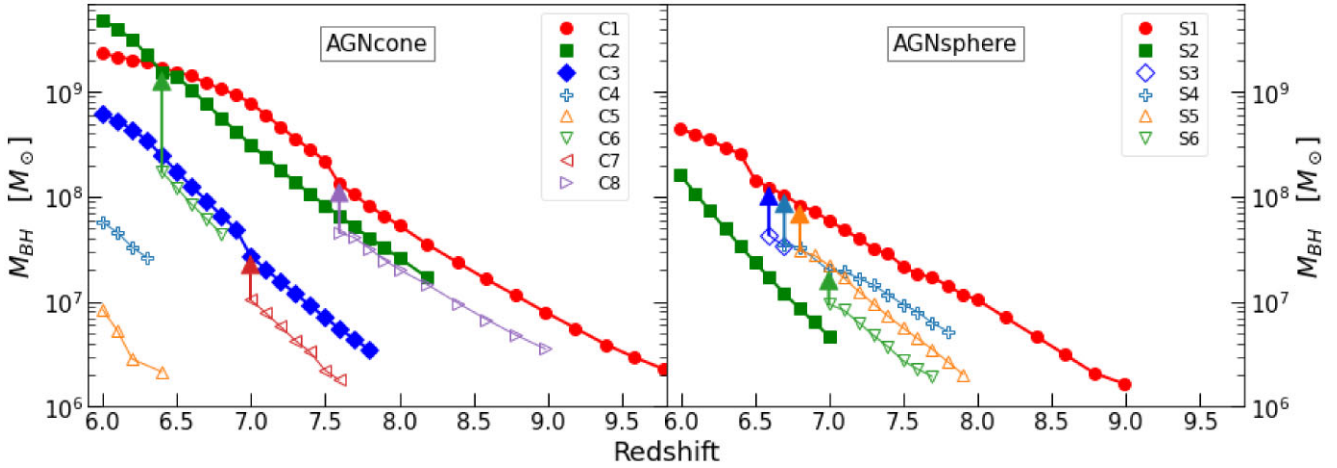


Figure 1. BH masses as a function of redshift for the *AGNcone* (left-hand panel) and *AGNsphere* (right-hand panel) runs. Only SMBHs accreting at $\dot{M} > 0.02 M_{\odot} \text{ yr}^{-1}$ are considered. The arrows mark the mergers between BHs. AGN considered in Sections 3 and 4 (i.e. those that reach $z = 6$ with $M_{\text{BH}} > 10^8 M_{\odot}$) are plotted as filled symbols.

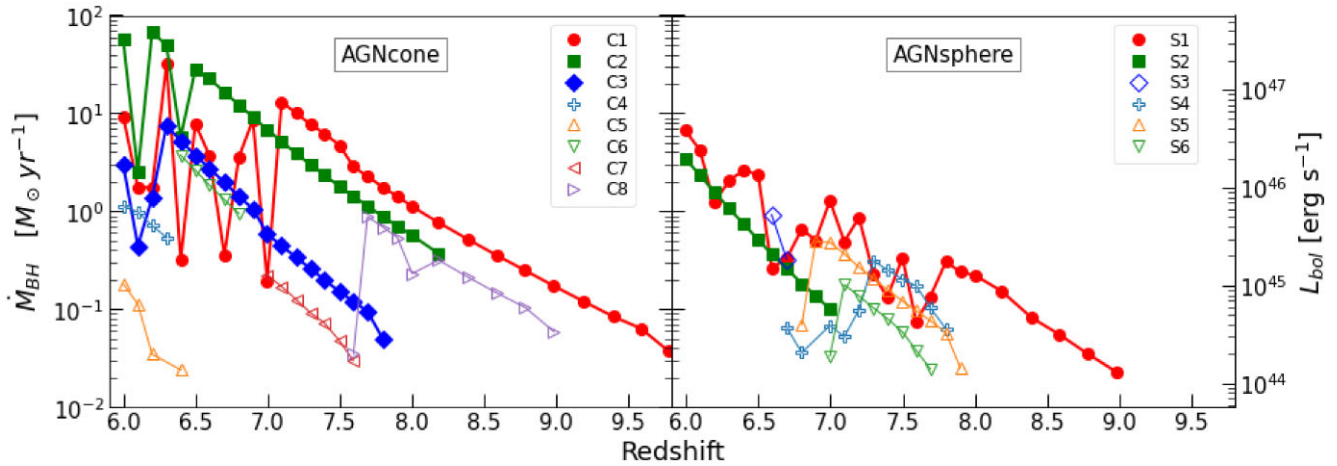


Figure 2. Mass accretion rate as a function of redshift for the *AGNcone* (left-hand panel) and *AGNsphere* (right-hand panel). The corresponding bolometric luminosity (equation 2) is reported in the right axis.

efficient accretion on to the SMBH. This effect can be appreciated in Fig. 2: The accretion rate of *AGNcone* is generally higher than that of *AGNsphere*, at least up to $\dot{M} \approx 1\text{--}30 M_{\odot} \text{ yr}^{-1}$. At higher accretion rates, which are reached by the most accreting BHs at $z < 7$, AGN feedback prevents further increase of the accretion rate.

Hereafter, we focus our analysis on the AGN that reach $z = 6$ with $M_{\text{BH}} > 10^8 M_{\odot}$ and $L_{\text{bol}} > 10^{46} \text{ erg s}^{-1}$; (see filled symbols in Figs 1 and 2), which we refer to as ‘bright AGN’ (i.e. C1, C2, and C3 in *AGNcone*; S1 and S2 in *AGNsphere*). These BH mass and luminosity values are typical of known $z > 6$ QSOs (e.g. Yang et al. 2021), allowing us to compare the physical properties of simulated and observed AGN in a consistent way. We note that, since the simulations focus on a single cosmic region at high redshift, the derived expectations on the AGN observable properties might be affected by cosmic variance.

2.3 Gas column density and radial distribution

Here, we describe the method that we use to derive the distribution of hydrogen, helium, and metal column densities in the ISM for galaxies hosting AGN in the considered simulations. We make use

of the hydrogen column density in the remaining of the paper to derive the observational properties predicted by the two considered simulations.

We estimate the distribution of the column densities for the bright AGN in the simulations by launching 1000 randomly selected line of sights (LOSs) toward each AGN from a distance $d = 30 \text{ kpc}$. Each LOS is considered as the axis of a cylinder with basis radius of R_{soft} . We note that the resolutions of the simulations do not allow us to probe structures on smaller scales, as, for instance, the existence of a dusty torus on pc scales. Then, each cylinder is divided along its length into bins of $l_{\text{bin}} = 0.25 \text{ kpc}$ width, for a total of $\frac{d}{l_{\text{bin}}} = 120$ radial bins. We compute the density of each chemical element in a bin of the cylinder from the mass carried by each particle included in that bin. With this approach, we also obtain the radial distribution of the gas density. Finally, we integrate along the cylinder to compute the total column density of hydrogen (N_{H}) and of the other elements. The resulting total N_{H} is not sensitive to reasonably different values of l_{bin} (i.e. from 0.25 to 1 kpc). Therefore, we used $l_{\text{bin}} = 0.25$ as this value allows us to sample well the radial distribution of the gas (see Section 4.2). Fig. 3 (upper panel) presents an example of the derived column-density map centred on the QSO C1 in *AGNcone*.

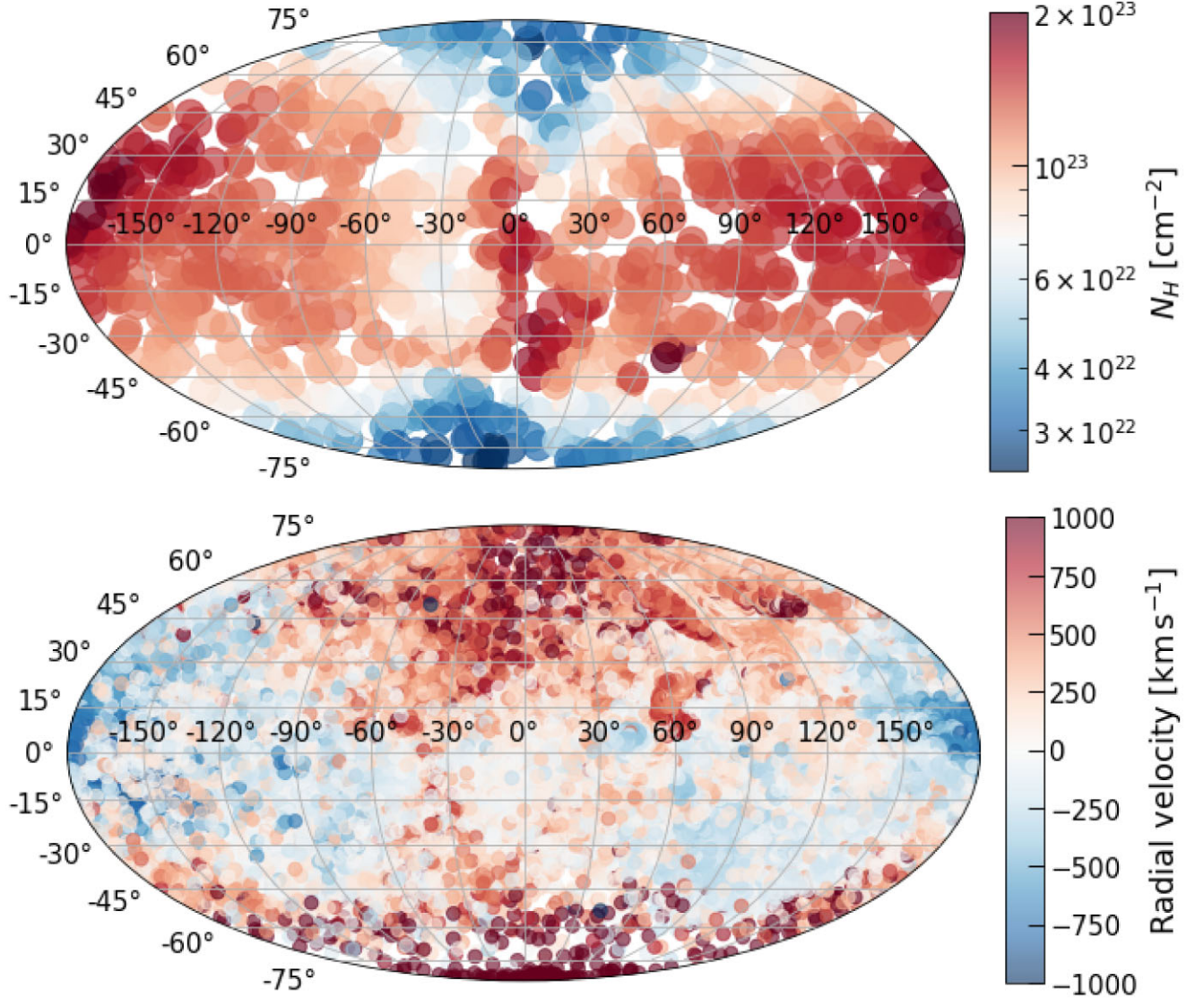


Figure 3. Upper panel: Mollweide projection of the column density along 1000 random lines of sights centred on the QSO C1 at $z = 7.1$. Lower panel: Mollweide projection of the radial velocities of all particles within 10 kpc from C1 at $z = 7.1$. The different sampling of the maps is intended to show the homogeneity of the 1000 LOSs used to compute N_H in the upper panel, and the velocity of the individual gas particles in the lower panel. The map is aligned with the outflow cone direction. Regions where the particles have high positive velocities correspond to the two cones along which the kinetic energy is distributed by the AGN feedback in the *AGNcone* simulation. Such cones are characterized by the lowest values of column densities.

Each circle represents one of the 1000 random LOSs, which sample homogeneously the entire solid angle as seen from C1.

To assess the effect of feedback on the column density (Section 3), we also consider an additional simulation run presented in B18, that is identical in terms of initial conditions and physical prescriptions to the *AGNcone* and *AGNsphere*, except that BHs are not seeded. The only type of feedback in this run, which we refer to as *noAGN*, is due to supernovae explosions (see B18 for detailed discussion).

We associate each AGN in a simulation to the corresponding galaxy in the *noAGN* runs following a method similar to that described in Zana et al. (2022): First, we identify the DM halo hosting the AGN as the one having its centre of mass closest to the position of the SMBH. Then, we identify the corresponding halo in the *noAGN* run by cross-matching the DM particle IDs in the two runs, and selecting the halo in *noAGN* which shares the largest fraction of particles with the initial AGN halo, further imposing

that the mass difference must be within a factor of 10–50. per cent². Finally, we repeat the procedure described above on the selected halo in *noAGN*, and derive the column density distribution in absence of AGN feedback. At $z > 8$, the redshifts at which the *noAGN* snapshots are taken are significantly different from those of the runs including AGN, making the DM-halo match procedure highly uncertain. Thus, we limit the identification of the AGN-hosting galaxies counterparts in the *noAGN* run to $z < 8$.

3 COLUMN DENSITY EVOLUTION

Fig. 4 presents the evolution of the column density for bright AGN in the *AGNcone* and *AGNsphere* simulations. Considering the *AGNcone*

²The exact threshold is adjusted at each time-step in order to find at least one halo counterpart.

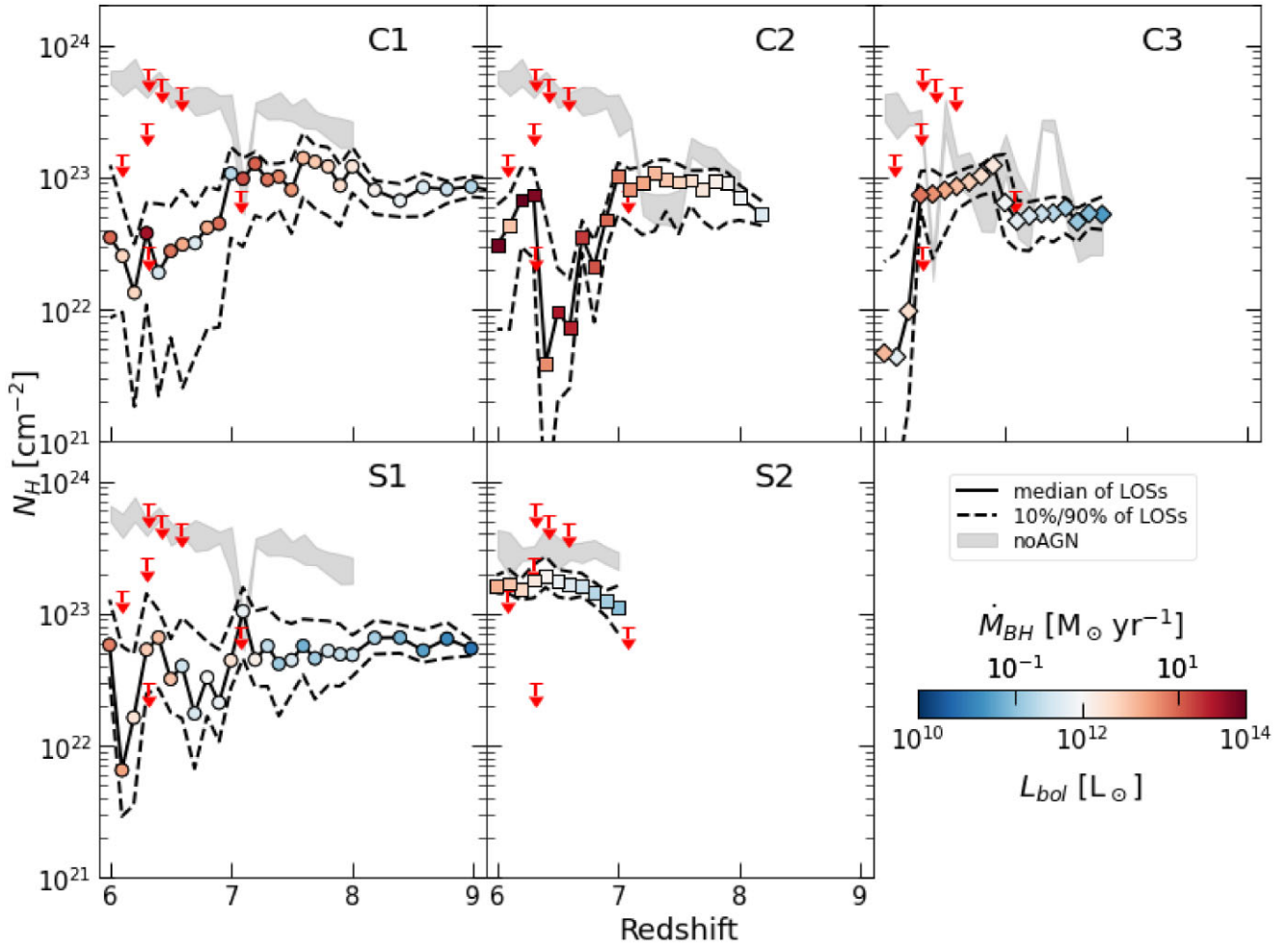


Figure 4. Evolution of column density for bright AGN in the *AGNcone* (C1, C2, C3) and *AGNsphere* (S1, S2) simulations. We show the median value (solid line, colour coded according to the AGN bolometric luminosity and accretion rate), and the 10 per cent and 90 per cent percentiles (dashed lines) computed by launching 1000 lines of sight. The grey stripes enclose the 10 to 90 per cent percentiles of the column densities of matched galaxies in the same simulation sets where, however, BHs have not been seeded (i.e. the *noAGN* case). To compare with observational results (Section 4.1), the red arrows mark the 3σ upper limits derived for X-ray detected QSOs with >10 counts from Nanni et al. (2018) and Connor et al. (2019).

simulation, the AGN column densities are similar to, or slightly lower than, those derived for the corresponding galaxies in the *noAGN* run until the AGN accretion rate reaches $\dot{M} \approx 10\text{--}30 M_\odot \text{yr}^{-1}$. This happens at $z \approx 7$ for C1 and C2, and $z \approx 6.3$ for C3 (see Fig. 2). At later times, the AGN column density drops significantly by up to ≈ 1 dex and the accretion rate starts to oscillate. The 10 and 90 per cent percentiles span up to one order of magnitude, especially at $z = 6\text{--}7$, when the accretion rates reach the maximum values, producing the most powerful conical outflows.

Instead, the column densities of the corresponding galaxies in *noAGN* (grey stripes in Fig. 4) keep on increasing relatively smoothly. This finding confirms the AGN N_H drop and the presence of unobscured LOSs to the effect of the conical kinetic feedback. At low accretion rates, the produced outflow cannot stop the infall of material, but once the accretion rate reaches high enough values, the energy carried by the outflow impacts a significant part of the gas in the halo, hindering further infalling, especially along the conical outflow directions. As a result, the N_H decreases, as well as the AGN accretion rate, until more material is allowed to accrete, producing a new burst of powerful feedback. Such a cyclic activity explains the decreasing median N_H , the wider N_H distribution, and

the oscillating $\dot{M}$ behaviour at later cosmic times. This result is in qualitative agreement with the self-regulation scenario discussed by, e.g. Sijacki et al. (2009), Dubois et al. (2013), Costa et al. (2014), Feng et al. (2014), Richardson et al. (2016), Trebitsch, Volonteri & Dubois (2019), according to which the AGN feedback controls the growth of the black hole and limits the duration of high-accretion episodes by emptying the host galaxy gas reservoir, provided that the accretion rate is sufficiently high.

However, we note that the physical interpretation of our results is complicated by the effect that one AGN may have on other AGN-hosting galaxies passing through its feedback cone. In fact, C1, C2, and C3 in *AGNcone* at $z < 7$ are always closer than 30 kpc, and reach minimum distances as small as 4 kpc. At these distances, powerful outflows launched from one AGN may affect nearby galaxies (e.g. Zana et al. 2022).

As an example of the feedback effect on the column density, in Fig. 3, we compare the N_H map centred on C1 with the radial velocity map of all particle within 10 kpc from C1. The maps correspond to $z = 7.1$, when C1 reaches a local maximum in accretion rate before the strong AGN feedback starts to impact significantly the N_H (Fig. 4) and $\dot{M}$ starts to oscillate. Comparing the column density

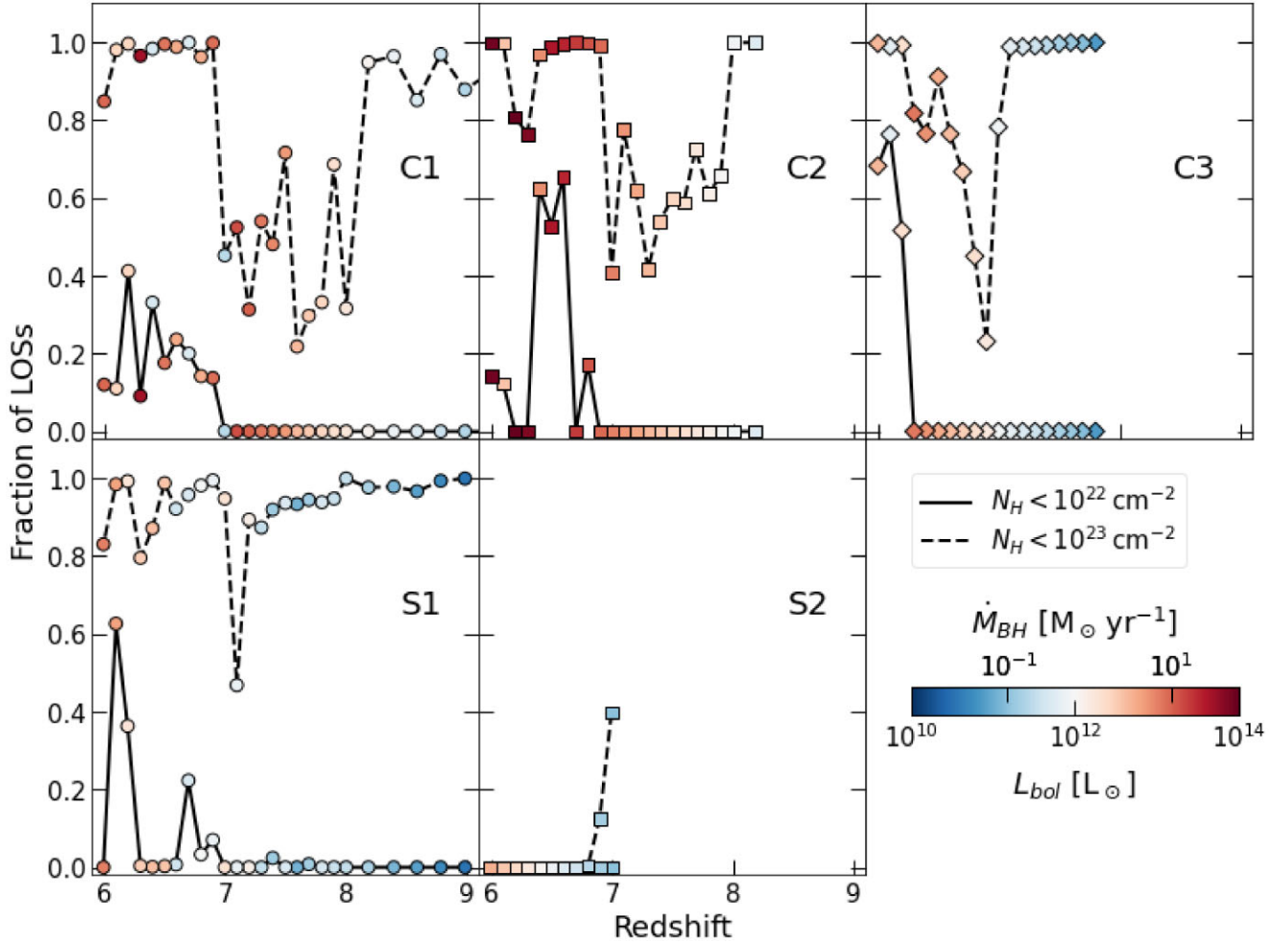


Figure 5. Fraction of lines of sight obscured by column densities $< 10^{22} \text{ cm}^{-2}$ (solid lines) and $< 10^{23} \text{ cm}^{-2}$ (dashed lines) as a function of redshift for the bright AGN in the *AGNcone* (C1, C2, C3) and *AGNsphere* (S1, S2) simulations. The symbols are colour coded according to the AGN bolometric luminosity and accretion rate.

map (upper panel) with the map of the radial velocity of individual particles (lower panel), we notice that the two conical outflows, identified as regions with positive radial velocities, correspond to LOSs with low column densities. Such LOSs are those along which high-redshift AGN are more easily to be detected in the rest-frame UV band, as we investigate in details in Section 4.3. Fig. 5 presents the fraction of LOSs along which $N_H < 10^{22} \text{ cm}^{-2}$ (solid lines) and $N_H < 10^{23} \text{ cm}^{-2}$ (dashed lines) for each bright AGN. Hereafter, we use the widely used threshold $N_H = 10^{22} \text{ cm}^{-2}$ to separate obscured and unobscured AGN.³ For instance, Merloni et al. (2014) found that such a value returns the best agreement between samples of obscured AGN as defined in optical (e.g. narrow emission-line AGN) and X-ray bands. From Fig. 5, we infer that only at $z \lesssim 7$ a fraction of the LOSs would appear as unobscured. In particular, at $z \lesssim 7$ C1 presents unobscured LOSs over 10–40 per cent of the solid angle, while this fraction is much more variable with redshift (i.e. 0–80 per cent) for C2 and C3.

The most massive BH in the *AGNsphere* simulation, S1, follows a somewhat similar N_H evolution to that of the AGN in *AGNcone*: a

roughly constant median N_H value up to $z \approx 7$ followed by a slightly decreasing and wider N_H distribution (Fig. 4), and the appearance of unobscured LOSs (Fig. 5) at later cosmic times. However, some differences exist: first, the AGN N_H is always significant lower than that of the corresponding galaxy in the *noAGNrun* (grey stripe), even at $z > 7$. Secondly, the column density drop at $z < 7$ is not as strong as in the *AGNcone* case. Finally, at $z > 7$, the accretion rate of S1 is not as smooth as in the *AGNcone* case, and keeps on increasing even at $z < 7$. These differences may be due to the prescribed geometry of the kinetic feedback in the *AGNsphere* case, in which gas particles are accelerated in a random direction during every accretion event. Therefore, in contrast with the *AGNcone* case, there is no preferential direction (i.e. the equatorial plane of the conical outflow) along which material can keep on accreting undisturbed for long periods of time at $z > 7$. In particular, the accretion rate of S1 never exceeds $\approx 10 M_\odot \text{ yr}^{-1}$, which is the approximate threshold after which the AGN kinetic feedback affects more evidently the N_H distribution and the accretion rate of AGN in the *AGNcone* run.

The column density evolution of S2, instead, does not appear to be strongly influenced by the AGN feedback. Although the median N_H is slightly lower than the values found in the *noAGNcase*, it remains constant with time, and does not drop even at $z < 6.5$, when S2 reaches similar accretion rate to S1. As a result, S2 would never

³However, we note that we consider the dust extinction as a more relevant quantity when we study the AGN rest-frame UV emission in Section 4.3.

appear as an unobscured AGN. We note that the typical column density of S2 is a factor of ≈ 3 higher than that of S1 at any redshift, and its accretion rate rises smoothly from $z = 7$ to 6. These properties suggest that higher accretion rates than the values reached by S2 are required in order to launch outflows powerful enough to sweep away the gas in the case of large column densities (e.g. Trebitsch et al. 2019), even when kinetic energy is distributed along random directions by the AGN feedback.

The median values of N_H we derive from the B18 simulations are consistent with typical values found by Lupi et al. (2022). However, the resolution of that work is $\times 85$ higher than our simulations, and allows the authors to sample compact regions of dense gas with $N_H \gtrsim 10^{24} \text{ cm}^{-2}$, especially at $z > 8$, when AGN feedback has not yet affected significantly the ISM distribution and density in the host galaxies. One of the main methodological differences with that work is that we compare the ISM densities in the same galaxies in which SMBHs are actively accreting or are not seeded at all. Thus, we probe directly the effect of AGN feedback on the ISM in the host galaxy.

4 COMPARISON WITH OBSERVATIONS

In this section, we compare the observable properties derived from the N_H distributions of the AGN predicted by the simulations (Section 3) with observational results. In particular, we focus on the comparison with constraints from X-ray observations (Section 4.1), the radial distribution of the gas reservoirs (Section 4.2), and the observed UV magnitudes (Section 4.3).

4.1 X-ray obscuration

X-ray observations are routinely used to constrain the column density of obscuring material along the LOSs of AGN. Low and moderate values of column densities ($N_H \lesssim 10^{22} \text{ cm}^{-2}$) can absorb soft X-ray photons (rest-frame energies $\lesssim 2 \text{ keV}$), whereas larger column densities are required to absorb a high fraction of more energetic photons. However, X-ray observations of high-redshift QSOs (e.g. Vito et al. 2019b; Wang et al. 2021b) sample rest-frame energies $E > 3 \text{ keV}$, and are thus sensitive only to high column densities ($N_H \gtrsim 3 \times 10^{23} \text{ cm}^{-2}$), at least at the sensitivities of currently available facilities. Moreover, all of the known $z > 6$ QSOs have been selected based on their unobscured rest-frame UV emission (i.e. they are optically classified as type 1 QSOs), and thus are not expected to be heavily obscured in the X-ray band. For these reasons, existing X-ray observations of bright $z > 6$ QSOs provide us with only loose upper limits of N_H . The downward-pointing red arrows in Fig. 4 are the observed upper limits on N_H derived for a sample of $z > 6$ QSOs by Nanni et al. 2017 and Connor et al. 2019, with typical luminosities $L_{\text{bol}} = 10^{46} - 10^{47} \text{ erg s}^{-1}$. The column densities derived for bright AGN in all of the considered simulations are lower than, or consistent with, such loose upper limits. Although the N_H values found for the *noAGN* case are typically higher, they are still consistent with some measured upper limits. Therefore, the constraints on N_H obtained from X-ray observations of high-redshift QSOs only marginally favour the presence of kinetic feedback.

We note that constraining AGN obscuration using X-ray observations requires an assumption on gas metallicity, as X-ray photons are mainly absorbed by metal atoms. Typically, solar metallicity is assumed, whereas the ISM metallicity of the host galaxies of the AGN in the B18 simulations is sub-solar (e.g. by factors of $\approx 2-3$ at $z = 6$; e.g. Zana et al. 2022). This consideration reinforces the overall consistency between the N_H values constrained from X-ray observations and found in the simulations, as significantly

larger column densities would be required in the case of sub-solar metallicities to produce X-ray obscuration in excess to that observed in real QSOs. In Section 5, we discuss the X-ray detectability of the QSOs in the simulations.

4.2 Gas radial distribution

We investigate the effect of kinetic feedback on the observable sizes of the gas reservoirs in high-redshift QSOs. From the radial distribution of N_H derived for each LOS in Section 2, we computed the radius from the centre of the galaxy which includes 90 per cent of the gas contributing to the total N_H . Then, for each galaxy, we computed the median value considering all of the 1000 LOSs, and define it as R_{90} . We use such a quantity to quantify the size of the gas reservoir in a galaxy.

Fig. 6 presents R_{90} as a function of redshift for every bright AGN in the *AGNcone* and *AGNsphere* simulations, as well as for the matched galaxies in the *noAGN* runs. All of the bright AGN in the *AGNcone* simulation (C1, C2, C3) have a similar evolution of R_{90} : their gas reservoir sizes are constant ($\approx 1 \text{ kpc}$) at $z \gtrsim 7$. At lower redshift, where N_H decreases due to strong effect of the kinetic feedback, which is proportional to $\dot{M}$ and L_{bol} (see the colour-code of the circles in Figs 4 and 6), R_{90} increases up to several kpc. This behaviour is expected considering that the AGN feedback applies a mechanical push to the surrounding gas particles. In fact, the size of the gas reservoir in the *noAGN* run, where the AGN feedback lacks (grey dashed lines in Fig. 6), remains constant or tends to even decrease at later cosmic times.

The evolution of R_{90} for S1 in the *AGNsphere* simulation is similar to that of the AGN in the *AGNcone* simulation. However, the increase of R_{90} is stronger and begins at earlier cosmic times. We recall that the accretion rate of S1 is typically lower than that of the AGN in *AGNcone* (see Fig. 2), and therefore the stronger evolution of R_{90} is not due to intrinsically stronger outflows launched by the AGN, but, as discussed in Section 3, to the different geometry of the outflow: being launched along random directions at every accretion event, the outflow is more likely to transmit the kinetic energy to the gas particles in the galaxy even at low or moderate accretion rates. Instead, S2 does not follow the same evolution as S1. On the contrary, R_{90} decreases to sub kpc values approaching $z = 6$. As discussed in Section 3, we ascribe this behaviour to the relatively low accretion rate, which does not produce feedback strong enough to efficiently affect the gas distribution in the host galaxy.

We compare our findings with the observed extent of the [C II] emission of 25 $z > 6$ QSOs presented by Venemans et al. (2020), assuming that the [C II] emission line is a good tracer of the spatial extent of the total gas reservoir (e.g. Zanella et al. 2018; Sommovigo et al. 2021). We used the major axis of the deconvolved [C II] emission size (table 3 of Venemans et al. 2020), which represents the FWHM of the emitting source, and converted it into the radius that includes 90 per cent of the [C II] light, assuming a Gaussian distribution.⁴ The resulting values are reported in Fig. 6 as black ticks at the redshift of each QSO.

The AGN in the *AGNcone* simulation have R_{90} consistent with the observed values, while the median gas size of S1 is larger at nearly every redshift. S2 have a size consistent with the most compact QSOs

⁴We note that the conclusions hold if we use an exponential profile (e.g. Fujimoto et al. 2020) and convert the FWHM values reported by Venemans et al. (2020) into exponential scale lengths. In this case, we obtain larger radii than in the Gaussian case by a factor of ≈ 1.75 .

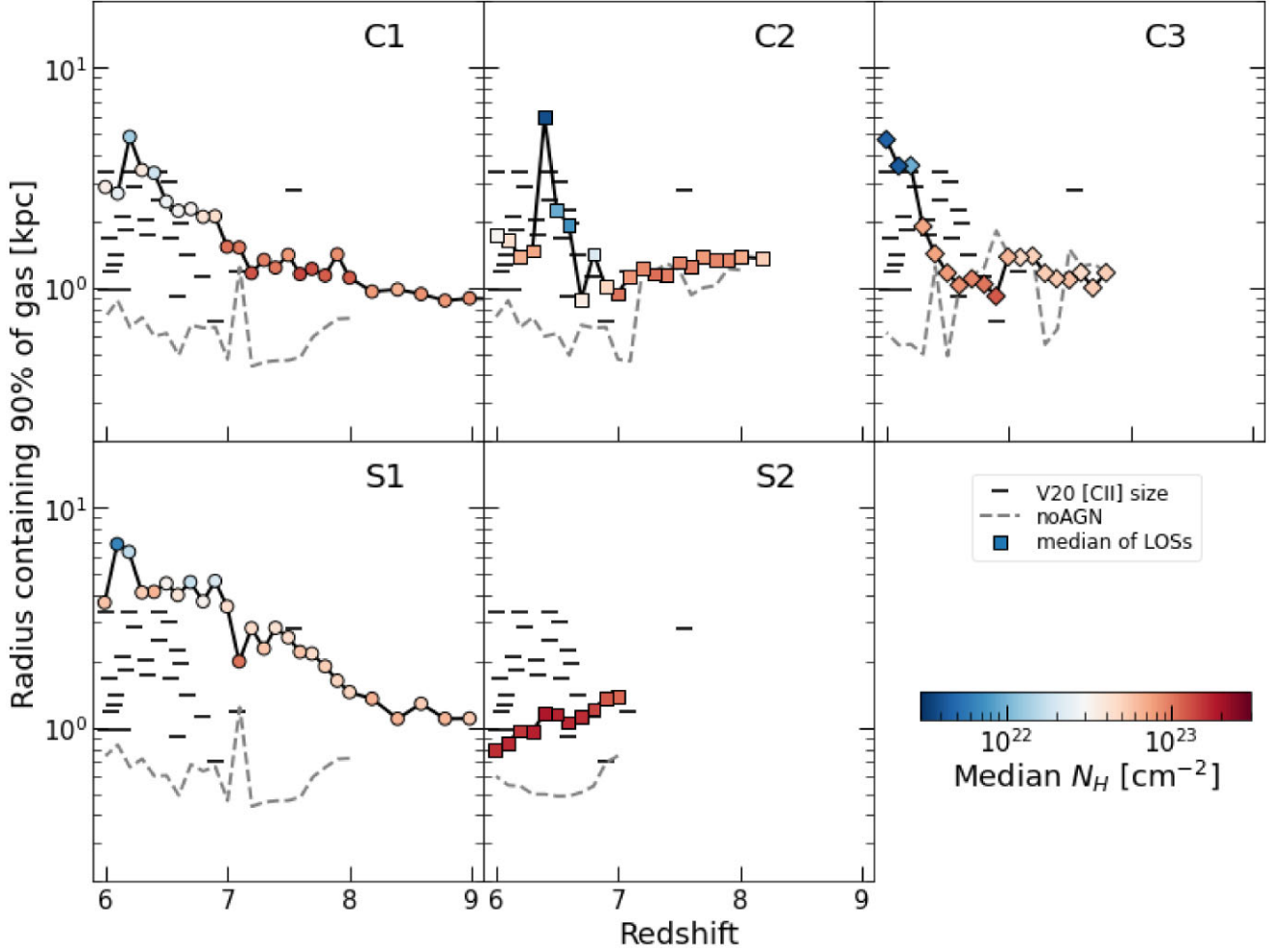


Figure 6. Median radius (R_{90} , computed over all of the lines of sight) containing the 90 per cent of the total gas as a function of redshift. The colour code indicates the median total N_H , averaged over all of the lines of sight. The dashed grey lines mark the same quantity computed for matched galaxies in the same simulation sets where, however, BHs have not been seeded (i.e. the *noAGN* case). The black ticks mark R_{90} for 25 $z > 6$ QSOs, as estimated from the [C II] emission beam-deconvolved sizes presented by Venemans et al. (2020).

in the Venemans et al. (2020) sample. However, this comparison is not fair: the ISM in S2 produces very large column densities at all redshifts and all LOSs (Fig. 4), and thus large expected values of dust extinction. All of the QSOs studied in Venemans et al. (2020) are instead rest-frame UV selected objects: We lack observational information about the extent of the gas reservoirs of buried high-redshift QSOs, as is S2. In all cases, the median gas size of the *noAGN* control galaxies are smaller than the observed values for QSOs, suggesting that kinetic feedback is required to produce the gas extents observed in real QSOs.

4.3 UV magnitudes

In Section 4.1, we discussed how the available X-ray observations of $z > 6$ QSOs are not sensitive to the column density values that we derived for bright AGN in the simulations. Instead, the rest-frame UV emission of high-redshift AGN is expected to be severely affected by dust extinction even for low values of N_H . In this section, we compare the expected rest-frame UV magnitudes of bright AGN in the simulations with the observed values of known $z > 6$ QSOs.

We assumed that the intrinsic (i.e. unextincted) rest-frame UV spectra of the AGN-hosting galaxies in the simulations are dominated

by the AGN (i.e. we do not include stellar emission) and are well represented by the Vanden Berk et al. (2001) composite spectrum of type 1 QSOs, rescaled to their bolometric luminosity via the bolometric correction of Venemans et al. (2016) and Decarli et al. (2018)

$$\log \left(\frac{L_{\text{bol}}}{\text{erg s}^{-1}} \right) = 4.553 + 0.911 \times \log \left(\frac{\lambda L_{\lambda}(1450\text{\AA})}{\text{erg s}^{-1}} \right). \quad (3)$$

We assumed a simple uniform slab of dust located in front of each AGN and an SMC extinction curve, and computed the measured rest-frame UV flux as

$$F_{\lambda}^{\text{obs}} = F_{\lambda}^{\text{intr}} e^{-\tau_{\lambda}}, \quad (4)$$

where $\tau_{\lambda} = k_{\lambda} \Sigma_m f_{\text{dust}}$, k_{λ} is the extinction cross-section at wavelength λ , Σ_m is the mass column density of metals, which we computed in Section 2.3, and the fraction of metal mass locked into dust is assumed to be $f_{\text{dust}} = 0.15$ as in Di Mascia et al. (2021b). Finally, we computed the apparent magnitude at the wavelength corresponding to rest-frame 1450 Å, that is m_{1450} .

For all considered AGN, the metal mass is computed from the column densities of metals derived in Section 2.3 for 1000 LOSs at every simulation snapshot. Thus, we obtain a distribution of 1000

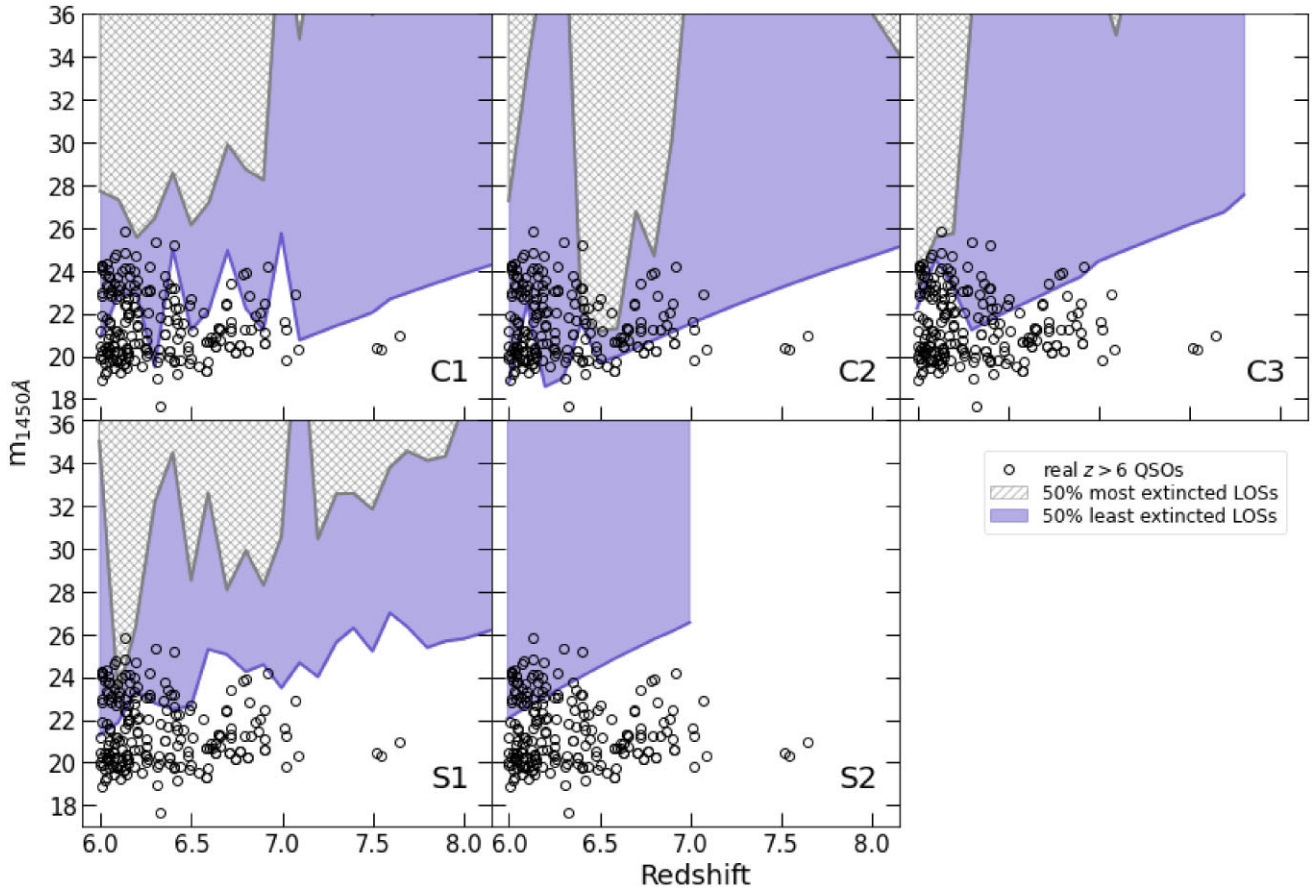


Figure 7. Apparent magnitude at rest-frame $\lambda = 1450 \text{ \AA}$ as a function of redshift for bright AGN in the *AGNcone* (C1, C2, C3) and *AGNsphere* (S1, S2) simulations. The purple regions encompass the 50 per cent least extinguished LOSs, while the grey hatched regions represent the 50 per cent most extinguished LOSs. The grey circles are $z > 6$ QSOs collected from Bañados et al. (2016, 2018), Chehade et al. (2018), Matsuoka et al. (2018a, 2018b, 2019a, 2019b), Mazzucchelli et al. (2017), Reed et al. (2017), Tang et al. (2017), Wang et al. (2017, 2018a, 2018b, 2019, 2021a), and Yang et al. (2020a, 2020b).

values of m_{1450} at every redshift. In Fig. 7, we show the magnitudes obtained for the 50 per cent least (purple regions) and most (grey hatched regions) extinguished LOSs. To allow for a comparison with observations, we add the magnitudes of a sample of $z > 6$ QSOs collected from Bañados et al. (2016, 2018), Chehade et al. (2018), Matsuoka et al. (2018a), Matsuoka et al. (2018b, 2019a, 2019b), Mazzucchelli et al. (2017), Reed et al. (2017), Tang et al. (2017), Wang et al. (2017, 2018a, 2018b, 2019, 2021a), Yang et al. (2020a, 2020b), with typical magnitudes of $19 \lesssim m_{1450} \lesssim 24$.

Among the considered simulations, *AGNcone* produces the UV brightest AGN, which are consistent with the magnitudes of known QSOs at $z \lesssim 7$. As discussed in Section 3, such redshift range corresponds to the period when the AGN strong kinetic feedback strongly affects the gas column density in the host galaxy, strongly suggesting that known, optically selected $z > 6$ AGN are indeed observed preferentially along directions where AGN feedback has cleared the LOS of most of the gas and dust. This prediction is hard to be tested observationally. Not only estimating the outflow direction is a difficult task, but the incidence of outflow in high-redshift AGN itself is still a matter of debate (e.g. Maiolino et al. 2012; Ciccone et al. 2015; Bischetti et al. 2019; Novak et al. 2020; Izumi et al. 2021; Meyer et al. 2022). Moreover, $z > 6$ QSOs might have been detected along LOSs which have been previously cleared of most of the gas and dust by past outflows. In this respect, a caveat

arises from the numerical implementation of the ISM properties in the *B18* simulations, which, as described in Section 2.1, follow the prescription of Springel & Hernquist (2003). This model does not capture the ISM porosity and therefore is not able to resolve clumpy structures on pc scales. Resolving such structures might decrease the effective opacity of the medium and possibly produce more unobscured lines of sight, even in the absence of AGN feedback.

In Fig. 7, only < 50 per cent of the LOSs of an individual AGN have extinction values small enough to reproduce the observed magnitudes. We computed the probability that multiple AGN appear as UV bright (i.e. $m_{1450} \lesssim 24$) sources along the same LOS, and found that it is negligible. This result is consistent with observations, according to which, to date, no such a system of multiple UV-bright AGN has been discovered at high redshift.

The most luminous AGN in the *AGNsphere* run, S1, reaches magnitudes as bright as the observed values only at $z \approx 6.5$, while it fails at reproducing the magnitudes of $z > 6.5$ QSOs. This is due to the lower accretion rate, and thus lower intrinsic luminosity, of S1 than the accretion rates of bright AGN in *AGNcone*. The large column density of S2 results in dramatic extinction levels along all of the LOSs, such that only along a small fraction of the LOSs S2 has apparent magnitude consistent with those of observed high-redshift QSOs, despite its intrinsic luminosity being similar to that of S1 at $z < 6.5$ (Fig. 2).

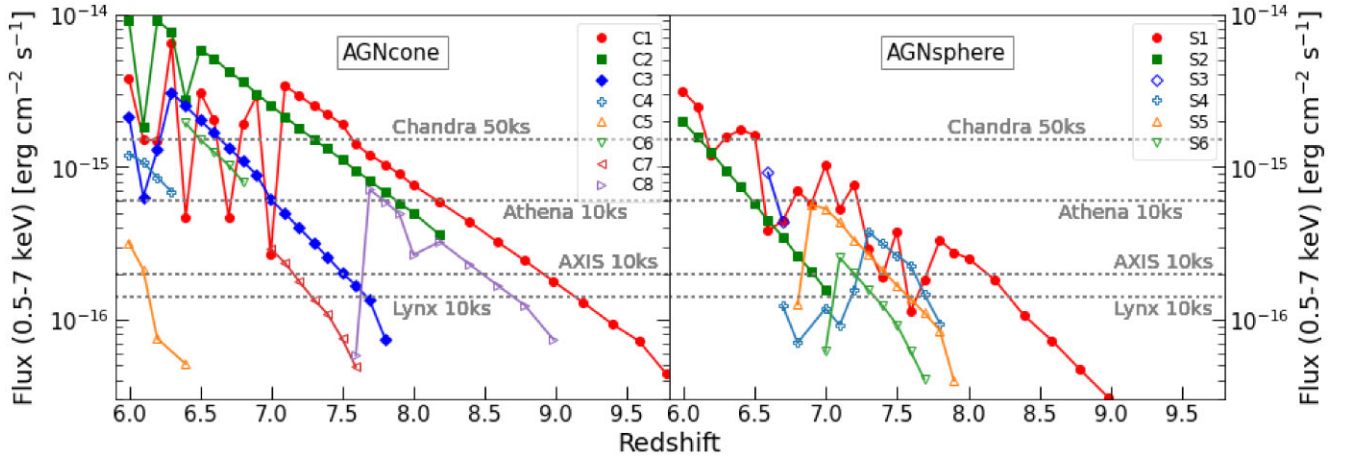


Figure 8. Expected X-ray flux in the 0.5–7-keV band as a function of redshift for the *AGNcone* (left panel), and *AGNsphere* (right panel) runs. For each AGN at each redshift, we assumed the median N_H computed for 1000 lines of sight. The horizontal dotted lines mark the flux limit computed for *Chandra* (50 ks observation), *Athena* (10 ks), *AXIS* (10 ks), and *Lynx* (10 ks).

5 MULTIPLE HIGH-REDSHIFT AGN ON 1-10 ARCSEC SCALES

Typical separations between AGN in the *B18* simulations are $\approx 5 - 50$ kpc, corresponding to only a few arcseconds in projection. To date, no multiple AGN system has been discovered observationally at $z > 6$ (e.g. Greiner et al. 2021), with the highest redshift AGN pair being recently discovered at $z = 5.7$ (Yue et al. 2021). This result could be due to dust extinction preventing the detection of other possible accreting SMBHs close to high-redshift QSOs, as we found in our simulations (Section 4.3). Alternatively, QSOs observed at $z \gtrsim 6$ intrinsically have no AGN satellite. The latter hypothesis implies that the simulations overpredict the number of bright AGN, due to, for example, the specific numerical setup and seeding prescription. In addition, as discussed in Section 2.1, the simulations focus on an overdense region, which maximizes the probability of forming multiple SMBHs, and thus bright AGN, in a small volume.

To investigate better the relation between the predicted and observed number of systems of multiple AGN at high redshift, in Section 5.1, we produce mock X-ray observations with the *Chandra* X-ray observatory⁵ based on the *AGNcone* and *AGNsphere* simulations. Then, in Section 5.2, we compute the probability of detecting multiple AGN on small angular separations, and compare the findings with observational results. Finally, in Section 5.3, we investigate the potential of future X-ray facilities in detecting possible multiple faint AGN over small scales around bright high-redshift QSOs.

5.1 Mock X-ray observations

As discussed in Section 4.1, the column densities that we derived in Section 3 for simulated $z > 6$ AGN have a negligible effect on the X-ray emission at the observed-frame energies probed by X-ray telescopes, allowing us to factor out the effect of varying N_H along different LOSs. However, we have to take into account another effect related to the specific choice of the LOS: the emission of different AGN might be blended along some LOSs due to projection effects,

and appear as a single X-ray source. This effect might be important as the projected angular separations of the AGN in the considered simulations are comparable with the angular resolution of *Chandra* (i.e. ≈ 0.5 arcsec), which is the existing X-ray observatory with the sharpest view.

We produce mock observations using the SOXS v. 3.0 software,⁶ using *Chandra* response matrices and ancillary files suitable for Cycle 20. SOXS accounts for three background components: a uniform Galactic component, a cosmic background due to point-like sources, and an instrumental component. For each simulation, we produce two sets of mock images, assuming an exposure time of 30 ks or 50 ks, which are typical lengths of real *Chandra* observations of $z > 6$ QSOs (e.g. Vito et al. 2019a; Wang et al. 2021b). For each set, we considered 100 random LOSs, along which all AGN have been projected on the sky plane according to their tri-dimensional positions in the simulations. This allows us to statistically take into account (1) the possible blending of multiple sources due to projection effects, and (2) the Poisson fluctuations of the number of detected X-ray photons at a given intrinsic flux.

We convert the bolometric luminosities of AGN in the simulations into X-ray luminosities in the rest-frame 2–10-keV energy band using the Duras et al. (2020) relation. Then, we compute the fluxes in the 0.5–7-keV band (i.e. one of the standard energy bands used to analyse *Chandra* observations) for every AGN, and use them as input values to simulate the images. We adopt intrinsic power-law emission with photon index $\Gamma = 2$. This is a typical value for AGN up to $z \approx 6.5$ (e.g. Nanni et al. 2017; Vito et al. 2019b), although Vito et al. (2019b) and Wang et al. (2021b) find hints for a steepening at higher redshifts. We also include absorption due to the measured value of column density along the considered LOS, although, as discussed above, the produced obscuration is negligible for our high-redshift objects, and a Galactic absorption component with $N_H = 5 \times 10^{20} \text{ cm}^{-2}$. These computations have been performed with XSPEC v.12.11 (Arnaud 1996; model *phabs* $\times$ *zphabs* $\times$ *powerlaw*).⁷ Fig. 8 presents the expected X-ray flux of every AGN in the simulations as a function of redshift.

⁶<https://hea-www.cfa.harvard.edu/soxs/>

⁷<https://heasarc.gsfc.nasa.gov/xanadu/xspec/>

⁵<https://cxc.harvard.edu/>

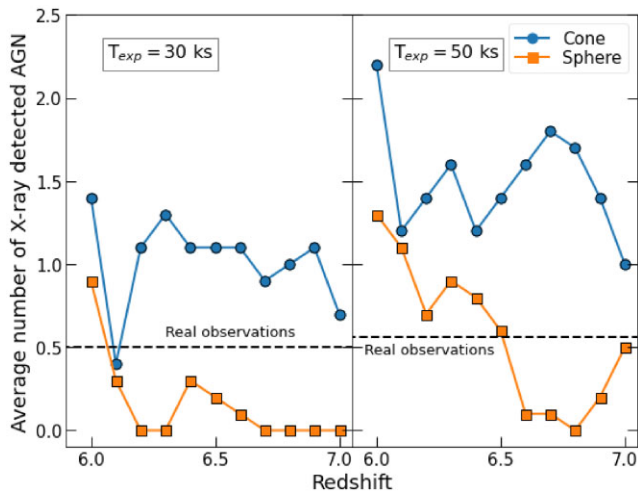


Figure 9. Average number of detected X-ray sources, averaged over 100 LOSs, detectable in the two simulations within <30 kpc from the central AGN with 30 ks (left-hand panel) and 50 ks (right-hand panel) *Chandra* observations. The black dashed line mark the average number of detected sources in real observations of $z > 6$ QSOs.

5.2 X-ray detection of multiple AGN

We ran a blind source detection procedure on the *Chandra* mock observations in the 0.5–7-keV band using the *wavdetect* tool in CIAO v.4.12⁸ (Fruscione et al. 2006), with a significance threshold of 10^{-5} , over an area corresponding to <30 kpc from the central QSO, to be consistent with the volume considered throughout this work (see Section 2). We repeated this procedure for all snapshots in the $z = 6–7$ range, which includes most of the $z > 6$ QSOs observed with *Chandra*, thus allowing for a fair comparison with real observations.

Fig. 9 presents the number of AGN detected in the mock *Chandra* observations with 30 and 50 ks exposures, averaged over the 100 LOSs, for each simulation. *AGNcone* predicts an average of ≈ 1 detectable AGN already with relatively short exposures (30 ks) and multiple detected X-ray sources using slightly longer observations (50 ks) over all of the considered redshift range. Instead, according to the *AGNsphere* run, 30 ks (50 ks) *Chandra* observations of $z \gtrsim 6.2$ ($z \gtrsim 6.5$) should typically return no detected source, but the probability to detect one or more AGN increases quickly approaching $z = 6$.

In order to compare these results with real data, we collected all of the available *Chandra* observations of $z = 6–7$ QSOs with exposure times of 20–40 and 40–80 ks (Tab. 1). The median exposure time of the 20–40 ks (40–80 ks) observations is 38 ks (54 ks) and the median redshift of the targeted QSOs is $z = 6.4$ ($z = 6.5$). These values are well matched to our sets of 30 and 50 ks mock images, respectively. We repeated the detection procedure described above on the real *Chandra* observations, considering only an area of $R < 30$ kpc from the targeted QSO, to allow for a fair comparison with the mock image results. We stress that the blind detection procedure prevents any bias related to rest-frame UV pre-selection of possible X-ray sources.

The last column of Table 1 reports the number of detected sources in the real observations,⁹ which are almost equally split between no

detected source and one detected source (i.e. the targeted QSO): The average numbers of detected X-ray sources in one observation are 0.50 and 0.56 for the 20–30- and 40–80-ks samples, respectively. Similar values are obtained by splitting each sample according to its median redshift. Comparing these results with the expected numbers of detected sources in simulations (Fig. 9), we find that *AGNcone* overestimates the number of detectable AGN at all redshifts, assuming both 30 and 50 ks exposure times. Instead, *AGNsphere* underestimates such number assuming 30 ks observations, while shows a strong dependence on redshift for longer exposures: At $z > 6.5$ and $z < 6.5$, it underestimates and overestimates, respectively, the average number of detected X-ray sources.

Due to the small sample sizes of real QSO observations and the narrow range covered by the number of detectable X-ray sources, it is difficult to provide a quantitatively robust comparison with the predictions from simulations. None the less, we attempt to do it by comparing the normalized histograms of detected sources in the mock and real observations over the entire $z = 6–7$ range (Fig. 10). This is justified by the relatively flat redshift distribution of the QSOs targeted by real observations (Tab. 1). For each set of mock images, we computed the two-sample Anderson–Darling test.¹⁰ The null hypothesis is that the mock and real observations are drawn from the same parent population, for what the number of detected X-ray sources is concern. We found that the null hypothesis can be rejected with high significance (i.e. Anderson–Darling test significance level $\lesssim 0.001$) for almost all combinations of simulations and exposure times. Fig. 10 confirms that *AGNcone* and *AGNsphere* overestimate and underestimate, respectively, the number of detectable X-ray sources. Mock simulations of *AGNsphere* with $t_{\text{exp}} = 50$ ks is the only set for which the null hypothesis cannot be rejected, although this simulation is not consistent with real observations for $t_{\text{exp}} = 30$ ks. It is worth noting that few $z > 6$ QSOs have been pointed with long *Chandra* exposures (100–500 ks; e.g. Nanni et al. 2018, Connor et al. 2020, Vito et al. 2021). Some of these observations were performed to check the presence of faint and possibly obscured AGN around $z > 6$ QSOs, for which companion galaxies have been detected with ALMA and *HST*. However, to date, no solid detection of such satellite AGN has been obtained (Vito et al. 2019a, 2021; Connor et al. 2019, 2020).

5.3 Predictions for future X-ray facilities

The high sensitivities of future X-ray facilities will allow us to push the search for AGN satellites of luminous optically selected QSOs at $z > 6$ down to intrinsic luminosities significantly lower than those probed with *Chandra*. In Fig. 8, we report as dotted grey lines the approximate expected sensitivity limits of future missions such as *Athena*/WFI (Nandra et al. 2013), *AXIS* (Mushotzky et al. 2019; Marchesi et al. 2020), and *Lynx*/HDXI (Gaskin et al. 2019), each one computed assuming 10 ks exposure time, and compare them with the sensitivity of a 50-ks *Chandra* observation. We computed these values by simulating X-ray observations of an X-ray source, assuming a simple power-law spectrum with photon index $\Gamma = 2$ and varying flux. In particular, for each instrument, we loaded

⁸<https://cxc.harvard.edu/ciao4.12/>

⁹We note that for almost all of the QSOs considered here, the results of the blind detection procedure agree with what reported in the literature, but for J084229.43+121850.5. Vito et al. (2019b) reported a detection of X-ray

emission from this QSO, while here we report it as undetected. This apparent discrepancy is due to the different detection procedure (i.e. blind detection versus rest-frame UV pre-selection of the target position) and significance threshold.

¹⁰We used the *anderson_ksamp* method of the SciPy package (Virtanen et al. 2020).

Table 1. Comparison sample of *Chandra* observations of $z = 6-7$ QSOs (see Section 5.2).

ID (1)	z (2)	Reference (3)	Observation ID (4)	t_{exp} (ks) (5)	N_{det} (6)
20–40-ks sample					
J002429.77+391319.0	6.621	W21	20416	20	0
J005006.67+344521.6	6.253	V19	20393	34	1
J022601.87+030259.4	6.541	V19	20390	26	1
J084229.43+121850.5	6.076	V19	20392	29	0
J104819.09-010940.2	6.676	W21	20415	35	0
J150941.78-174926.8	6.122	V19	20391	27	1
J152637.84-205000.7 ^a	6.586	C20	22165	33	0
J163033.90+401209.7	6.065	V19	5618	27	1
40–80-ks sample					
J010953.13-304726.3	6.791	V19	20398, 22214	66	0
J030516.92-315055.9	6.614	V19	20394	50	0
J103027.11+052455.1 ^a	6.308	N17	19926	50	1
J111033.98-132945.6 ^a	6.515	V19	20397	54	0
J114816.65+525150.4	6.419	G17	17127	78	1
J164121.73+375520.2	6.047	V19	20396, 21961	54	1
J203210.0-211402.3 ^a	6.24	C19	20470	45	1
J223255.14+293032.3	6.666	V19	20395	54	1
J234833.34-305410.0	6.902	W21	20414	42	0

^aThese QSOs have been observed with multiple ObsIDs, resulting in longer total exposure times than those reported here. We only consider the reported ObsIDs to allow for a fair comparison with our 30 ks and 50 ks mock observations. Notes. Column (1) – ID of targeted QSO; Column (2) – redshift of targeted QSO; Column (3) – reference for published X-ray data. C19: Connor et al. (2019). C20: Connor et al. (2020). G17: Gallerani et al. (2017). N17: Nanni et al. (2017). V19: Vito et al. (2019b). W21: Wang et al. (2021b). Column (4) – *Chandra* observation ID considered in this work; Column (5) – Exposure time; and Column (6) – number of detected X-ray sources according to the procedure described in Section 5.2.

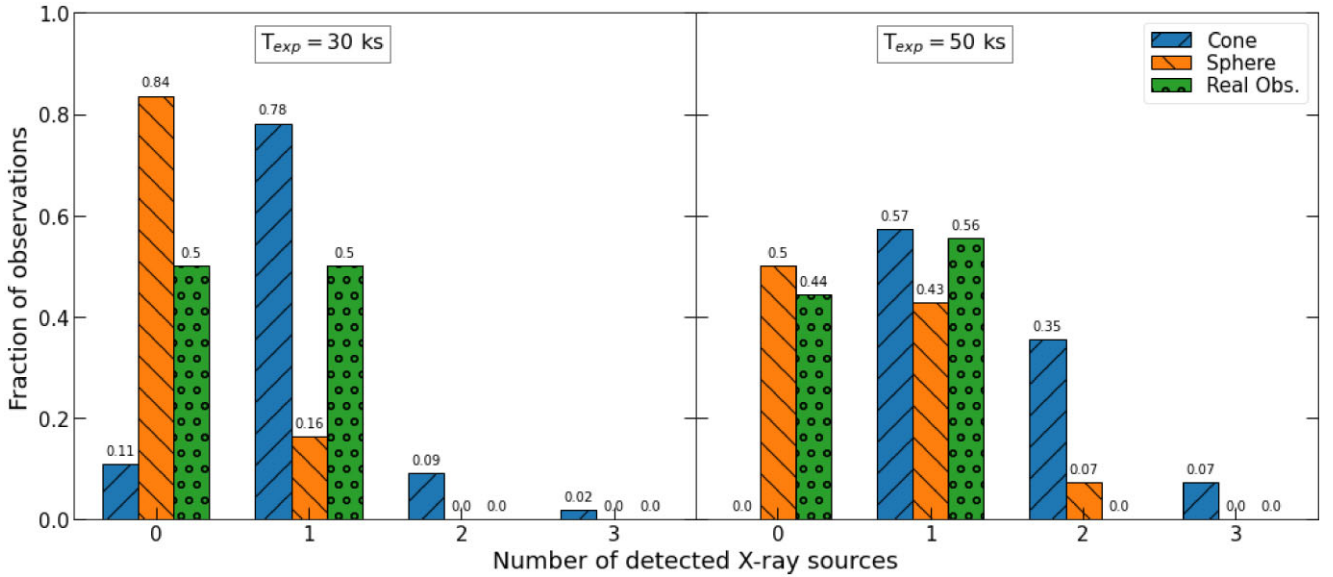


Figure 10. Normalized histograms of the number of detected X-ray sources in the mock and real *Chandra* observations of $z = 6-7$ AGN, for $t_{\text{exp}} = 30$ ks (left-hand panel) and 50 ks (right-hand panel).

response matrices and background files¹¹ in XSPEC, and computed the expected source and background count rates in a region including ≈ 90 per cent of the expected point spread function (PSF); i.e.

¹¹We use real response matrices and background files for *Chandra*, and the preliminary files included in SOXS for *Lynx*, *AXIS*, and *Athena*.

$R = 1$ arcsec for *Chandra*, *AXIS*, and *Lynx*, and $R = 5$ arcsec for *Athena*. Then, we computed the flux that returns a binomial no-source detection probability (i.e. P_B ; Weisskopf et al. 2007) such that $(1 - P_B) = 0.997$, corresponding to 3σ in the Gaussian approximation.

Fig. 8 shows that all of the considered next-generation X-ray mission will provide us with a huge improvement in the capability

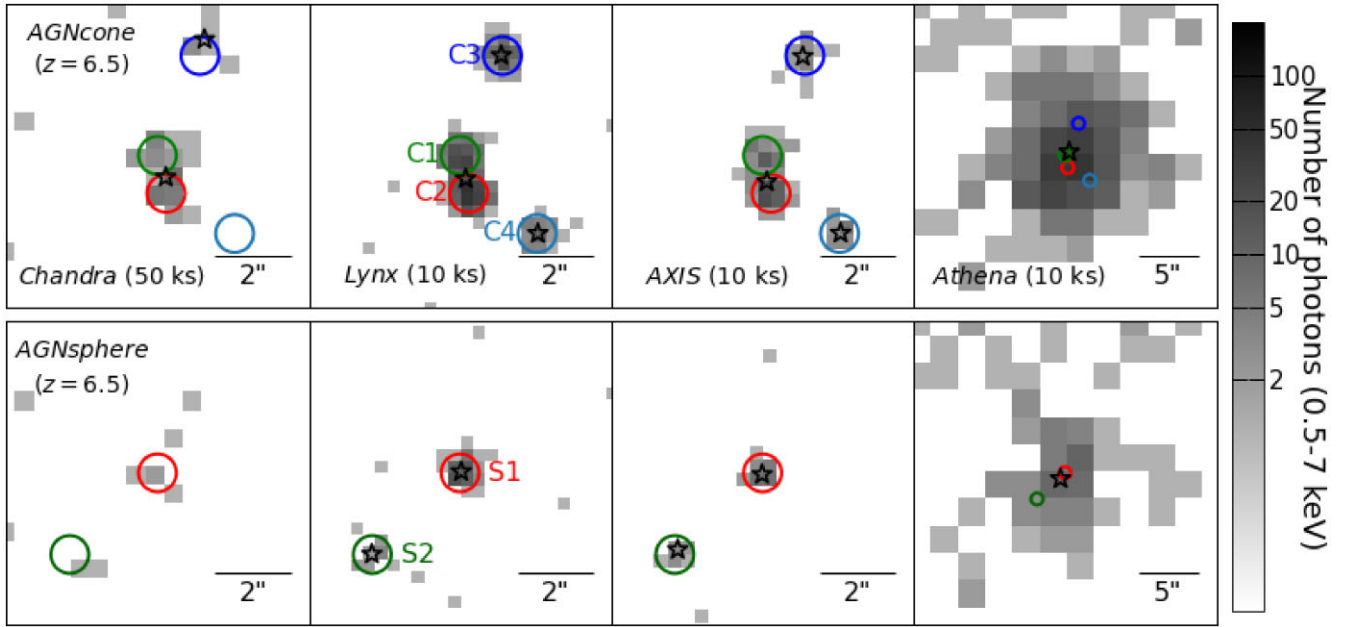


Figure 11. Simulated X-ray observations in the 0.5–7-keV band of the most-massive AGN at $z = 6.5$ and the surrounding satellite AGN in the *AGNcone* (upper row) and *AGNsphere* (lower row) simulations. From the leftmost to the rightmost columns, we simulated observations with *Chandra*/ACIS-S (50 ks), *Lynx*/HDXI (10 ks), *AXIS* (10 ks), and *Athena*/WFI (10 ks). For presentation purpose, the angular scale of the *Athena* image is different from the other cases, due to the larger PSF. The circles mark the location of the simulated AGN for a representative line of sight, and are colour coded as in Fig. 1. The black stars mark the position of X-ray detected sources obtained with a blind detection procedure.

of detecting faint AGN at $z > 6$, including satellite AGN around bright QSOs at $z > 6$, in a fraction of the time of a typical *Chandra* observation. Fig. 11 presents simulated X-ray observations with *Chandra* (50 ks), *Lynx* (10 ks), *AXIS* (10 ks), and *Athena* (10 ks) of a representative snapshot (i.e. $z = 6.5$) and LOS of the two simulation runs. The satellite AGN will appear as multiple X-ray sources on a few arcsec scales. This implies that, in addition to high sensitivity, excellent angular resolution, such as that provided by *AXIS* and *Lynx*, is required to detect them individually. To probe this issue, we performed a blind detection run with *wavdetect* on these images, and compared the detected sources (black stars in Fig. 11) with the input AGN (colored circles): The identification of close objects like C1 and C2 is difficult even with missions with ≈ 0.5 arcsec angular resolution. The problem is clearly more evident with *Athena*, due to its PSF of a few arcsec.

6 DISCUSSION

As discussed in Section 2.1, the outflow directions in the considered simulations are assumed not to be physically related to the host-galaxy properties and to be time-independent. In particular, the *AGNcone* simulation does not assume the outflow to be perpendicular to the plane of the host galaxy, as suggested by several observations of kpc-scale outflows or radio jets in the local Universe (e.g. García-Burillo et al. 2014; Cresci et al. 2015; Morganti et al. 2015; Venturi et al. 2021), where the outflow geometry can be studied in details, and by some numerical simulations (e.g. Hopkins et al. 2012).

Several physical mechanisms can concur in the acceleration of winds at sub-pc scales that eventually produce large-scale outflows, including magnetohydrodynamic effect (e.g. Sadowski et al. 2013), thermal driving (e.g. Proga 2007), radiation pressure acceleration, either applied on dust (e.g. Ishibashi & Fabian 2015) or mediated by UV transitions (e.g. Proga & Kallman 2004; Mizumoto et al. 2021),

which might produce outflows with different geometries. Moreover, the outflow geometry might be affected by interactions with the surrounding environment as the outflow expands (e.g. Nelson et al. 2019; Talbot, Sijacki & Bourne 2021), and might change with time. Cosmological simulations cannot describe in detail such a complex, and largely unknown, physics and evolution of outflows with relatively simple numerical recipes.

The goal of this paper is to investigate the effect of two particular large-scale outflow geometries (i.e. a spherical outflow and a bi-conical outflow parametrized as described in Section 2.1) on the observable properties of high-redshift AGN, regardless of the sub-grid physical mechanisms responsible for their acceleration. Extensive numerical simulations with identical initial conditions and physics except for the outflow parameters would be required to check whether and how the results are sensitive to different choices of the outflow parameters.

Kinetic feedback produced during the phases of fast accretion of SMBHs in the B18 simulations has a significant impact on the surrounding material and is required to match the predicted observable properties of bright AGN with observational results. One of the strongest pieces of evidence is represented by the study of the gas extent in the AGN host galaxies (Fig. 6): The gas reservoirs in the *noAGN* case (i.e. in absence of AGN feedback) are always more compact than those derived from ALMA observations of $z > 6$ QSOs (see also, e.g. van der Vlugt & Costa 2019). The effect of AGN feedback pushes the gas in the host galaxies to larger distances (i.e. up to a few kpc) from the centres, in agreement with observations (e.g. Ciccone et al. 2015; Bischetti et al. 2019; Venemans et al. 2020; Izumi et al. 2021). Although other mechanisms related to AGN feedback may produce such an observable, by, for instance, preventing gas infall from large scales (e.g. Trussler et al. 2020) or causing fluctuations in the gravitational potential, which may lead to a radial migration of the material (e.g. van der Vlugt & Costa

2019), B18 found that the mechanical removal of gas from the inner region of the host galaxies is the main process that affects their gas content in their simulations. We underline that also some $5 < z < 7$ star-forming ($1\text{--}70\text{ M}_\odot\text{ yr}^{-1}$) galaxies have been found to show both an extended [C II] halo (e.g. Fujimoto et al. 2020) and broad wings in the [C II] emission-line profile (e.g. Gallerani et al. 2018; Ginolfi et al. 2020), suggestive of outflows possibly powered by a yet undetected accreting MBH (e.g. Orofino, Ferrara & Gallerani 2021).

At $z < 7$, the feedback produces a general decrease of the N_H (Fig. 4), allowing for the appearance of unobscured (i.e. $N_H < 10^{22}\text{ cm}^{-2}$) LOSs (Fig. 5). Such directions are most probably those along which known $z > 6$ QSOs are preferentially observed, as the rest-frame UV selection of these objects requires low dust extinction. In fact, at $z \lesssim 6.5$, when the feedback effect is the strongest, bright AGN in the *AGNcone* simulation are able to reach the UV magnitudes observed for known $z > 6$ QSOs (Fig. 7). However, such LOSs represent only a fraction of the total LOSs of an AGN (see also Trebitsch et al. 2019; also, e.g. Ni et al. 2020; Lupi et al. 2022): more than half of the LOSs would appear too faint to be selected as high-redshift objects in current optical/near-IR surveys, suggesting that a large fraction of the high-redshift, intrinsically luminous QSO population is observationally missed due to strong UV extinction produced by the ISM only. The presence of a dusty torus on pc scales, which is not included in the simulations we have analysed, would further increase such a fraction.

The outflow geometry likely plays an important role: in the case of a conic outflow, SMBH accretion proceeds at maximum efficiency through equatorial infalling of gas until $\dot{M} \approx 10\text{--}30\text{ M}_\odot\text{ yr}^{-1}$ (Fig. 2), producing BHs with masses of $> 10^9\text{ M}_\odot$ at $z = 6\text{--}7$ (Fig. 1). At these accretion rates, the feedback regulates further accretion and reduces the typical obscuring column density, in particular along the cone direction (Fig. 3). In the case of outflows launched along random directions, the feedback can affect the growth of the SMBH and the N_H distribution even at lower accretion rates, resulting in $< 10^9\text{ M}_\odot$ BHs at $z = 6$, provided that the gas in the host galaxy is not too dense, as in the case of S2. Thus, the ISM properties (i.e. N_H and radial size of the gas) of the brightest AGN in the *AGNsphere* run is in agreement with observations. However, hindering the formation of $> 10^9\text{ M}_\odot$ BHs, the spherical geometry of the feedback in *AGNsphere* prevents AGN from reaching intrinsic luminosities comparable to known $z > 6$ QSOs at most redshifts (Fig. 7).

Interestingly, even the most luminous AGN in *AGNcone* cannot explain the detection of UV-bright QSOs at $z \approx 7.5$ (Fig. 7), due to the combination of the relatively small BH masses, and hence low accretion rates, which, by construction, are capped at the Eddington rate, and typically high N_H at that early cosmic time in this simulation. The existence of bright QSOs at $z \approx 7.5$ (e.g. Bañados et al. 2018; Wang et al. 2021b) requires different physical conditions for the SMBH formation and mass growth from those adopted in the considered simulations.¹² Future numerical simulations may explore such conditions as viable ways to reconcile the expected and observed properties of $z > 7$ AGN. Non-mutually exclusive possibilities are as follows:

(i) Different BH seeding mechanisms, that is, bright and massive QSOs discovered at $z \approx 7.5$ may be grown from more massive BH seeds or have been seeded at earlier redshift than the SMBHs in the simulations.

(ii) Sustained periods of super-Eddington accretion at $z > 7.5$, whereas in the simulations the SMBH accretion rate is capped at the Eddington limit.

(iii) Mass accretion characterized by a lower radiative efficiency than the value used in the simulations (i.e. $\epsilon_r = 0.1$). In this case, the mass that is not converted into radiation contributes to the growth of SMBH, which can reach higher masses than those found in simulations at a given time. For instance, Davies, Hennawi & Eilers (2019) report observational evidence for possible low radiation efficiency ($\epsilon_r \approx 0.001$) in high-redshift QSOs.

(iv) High-redshift AGN typically reside in regions which are even more overdense than that investigated in the B18 simulations, thus favouring the formation of SMBHs at earlier epochs. However, this possibility would arguably make the discrepancy between the observed and expected number of multiple X-ray detected AGN on small scales even worse. In addition, observational studies return contradictory results on the typical large-scale environment of high-redshift AGN (e.g. Ota et al. 2018; Mazzucchelli et al. 2019; Mignoli et al. 2020; Overzier 2022).

The analysis that we have performed demonstrates that the comparison between several observable properties of AGN predicted by the B18 simulations and the observational results, including both the properties of the individual galaxies and the environment, can help us to validate the recipes and assumptions adopted in numerical simulations. In particular, we found that AGN in the considered simulations match the gas radial distributions and apparent UV magnitudes of high-redshift QSOs. In addition, the same set of simulations has been demonstrated to reproduce well a number of physical properties of $z > 6$ QSOs, such as dust properties (Di Mascia et al. 2021b), multiwavelength spectral energy distribution (Di Mascia et al. 2021a), and the number of UV-detected and [C II]-detected satellite galaxies (Zana et al. 2022).

However, we also found that the predicted number of X-ray detectable satellite AGN located over small scales around luminous high-redshift QSOs both in the *AGNcone* and *AGNsphere* simulations does not agree with the observational results. This observable is relatively easy to estimate from simulations as it depends primarily on the BH accretion rate only, once a suitable conversion to X-ray luminosity is assumed. Moreover, gas and dust absorption does not affect significantly the observed X-ray emission from high-redshift AGN, as opposed to UV emission, up to high column densities ($\log \frac{N_H}{\text{cm}^{-2}} \approx 23.5\text{--}24.0$; see Sections 4.1 and 4.3).

The mismatch between the number of multiple X-ray detected AGN on small scales between simulations and observations may be related to numerical issues and physical prescriptions. In particular, the simplistic BH seeding recipe implemented in the considered simulations (i.e. a $10^5\text{ M}_\odot$ BH is placed in the centre of a galaxy when this reaches a given mass threshold) naturally leads to the formation of a large number of SMBHs, that would appear as bright AGN at later cosmic times. Similar seeding recipes have been commonly adopted by most cosmological simulations (e.g. Costa et al. 2014; Di Matteo et al. 2017; B18; Smidt et al. 2018; Lupi et al. 2019; Valentini et al. 2021), and typically mimic the ‘heavy seed’ formation channel for SMBHs (e.g. Lodato & Natarajan 2006; Ferrara et al. 2014). However, theoretical models of ‘heavy seed’ formation require stringent physical conditions on, e.g. metallicity, physical state of the gas, ad radiation fields (e.g. Ferrara et al. 2014). Accounting for such conditions in cosmological simulations is particularly difficult, but would reduce the number of formed SMBHs, and thus the discrepancy with observational results.

¹²As mentioned in Section 2.2, cosmic variance may affect our conclusions, as the simulations focus on a single cosmic region at high redshift.

Another possibility is that observed QSOs at high redshift do not reside in regions as dense as those probed in the analysed simulations (but see, e.g. Zana et al. 2022). In this case, the formation of multiple SMBHs is expected to be hindered, helping us reconcile the expected number of X-ray sources with observational results. In addition, we would also expect to form less massive BHs, with direct consequences on the observational expectations discussed in this paper, as the BH mass is tightly linked with the maximum accretion rate, and thus AGN luminosity and feedback strength. Qualitatively, we would expect to derive fainter rest-frame UV and X-ray fluxes, weaker feedback, and, as a consequence (see Fig. 6), more compact gas reservoirs (i.e. similar to the *noAGN* case) than the values discussed in Sections 4.2, 4.3, and 5. Future X-ray facilities will provide us with the required sensitivity and angular resolution to investigate the presence of multiple faint AGN around bright high-redshift QSOs down to unprecedented flux limits (see Section 5.3).

7 SUMMARY AND CONCLUSIONS

We studied the observable properties of $z = 6-10$ bright AGN in a suite of zoom-in cosmological simulations by B18 characterized by the inclusion of AGN kinetic feedback with either bi-conical (namely, *AGNcone*) and spherical (*AGNsphere*) outflow geometry. We focused our investigation on the gas column density and size in the host galaxies, the AGN rest-frame UV magnitude and X-ray fluxes, and the detectability of systems of multiple AGN over a few kpc scale in the X-ray band. We compared these quantities with a control simulation in which SMBHs are not seeded (i.e. *noAGN*), and observational results of $z > 6$ AGN. We summarize our findings as follows.

(i) *AGNcone* produces three bright AGN that grow up to $5 \times 10^8 < M_{\text{BH}} < 5 \times 10^9 M_{\odot}$ at $z = 6$. These objects are characterized by a steady increase of their accretion rate up to $\approx 10-30 M_{\odot} \text{ yr}^{-1}$. Once such high values are reached (at $z \approx 6.5-7$), the strong AGN feedback prevents further increase of the accretion rate. This behaviour is linked with the bi-conical geometry of the outflow, which allows steady infalling of material along the equatorial directions, at least until the feedback grows strong enough to affect most of the gas in the galaxy halo. In *AGNsphere*, the spherical geometry of the outflow affects gas accretion already at low and moderate SMBH growth rates. For this reason, the two bright AGN produced in *AGNsphere* reach lower values of BH masses (i.e. $2 \times 10^8 < M_{\text{BH}} < 5 \times 10^8 M_{\odot}$) and accretion rates ($\dot{M} < 10 M_{\odot} \text{ yr}^{-1}$) than objects in *AGNcone*.

(ii) AGN host galaxies in *AGNsphere* have gas column densities of $N_H \approx 10^{23} \text{ cm}^{-2}$ from their formation up to $z = 6.5-7$, when N_H presents a remarkable drop due to the strong AGN feedback. In fact, the N_H in matched galaxies in *noAGN* continues to increase during the entire considered redshift range. The brightest AGN in *AGNsphere* presents a similar behaviour as those in *AGNcone*, although the N_H is typically slightly lower. We interpret this difference again as due to the assumed spherical symmetry of the outflow. Instead, the second bright AGN in *AGNsphere* do not reach accretion rate sufficiently high to significantly affect the gas in the host galaxy. Our findings are consistent with the upper limits on N_H recently reported for a set of $z > 6$ AGN observed in the X-rays.

(iii) Kinetic feedback is required to match the gas extent reported for high-redshift QSOs (i.e. up to a few kpc). In fact, galaxies in *noAGN* present typical gas sizes of < 1 kpc, while the extents of the gas reservoirs of AGN in *AGNcone* and *AGNsphere* increase up to the observed values of a few kpc at $z \lesssim 7$. The exception is the

second bright AGN in *AGNsphere*, due to its relatively low values of accretion rate.

(iv) All AGN in the simulations would appear as obscured (i.e. $N_H > 10^{22} \text{ cm}^{-2}$) along all lines of sight (LOSs) at $z > 7$. These objects would be missed by currently employed UV-based selection methods, which are heavily affected by dust extinction, and would require observations in different bands (e.g. X-ray or infrared) to be unveiled. At later cosmic times, a fraction of LOSs (up to ≈ 80 per cent, depending on the specific AGN and redshift) have $N_H < 10^{22} \text{ cm}^{-2}$. These are the preferential directions along which known, UV-selected $z > 6$ QSOs are observed.

(v) Under simple, but reasonable, assumptions on the gas-to-dust mass scaling and dust distribution, we estimate the apparent UV magnitudes (m_{1450}) of the AGN in the simulations along different LOSs. We found that AGN in *AGNcone* have m_{1450} consistent with those observed for real high-redshift QSOs (i.e. $m_{1450} < 25$) along $\lesssim 50$ per cent of the LOSs at $z < 7$. AGN in *AGNsphere*, instead, have fainter magnitudes, due to the lower intrinsic luminosities, and, for the second AGN, the high extinction levels along most of the LOSs. No AGN in the simulations can reproduce the observed UV magnitudes of the few $z \approx 7.5$ QSOs known to date, whose formation and accretion history are likely not well captured by the prescriptions assumed in the simulations.

(vi) The presence of multiple bright AGN over scales of a few kpc led us to investigate their detectability in X-ray observations with *Chandra*, and to compare the results with real observations of $z > 6$ QSOs. We found that the *AGNcone* run significantly overpredicts the number of X-ray detected multiple AGN at high redshift. Instead, *AGNsphere* produces AGN with lower rate of X-ray detection than typical values derived in relatively shallow (i.e. 30 ks) observations, while it is consistent with the results obtained with longer (i.e. 50 ks) observations.

These results demonstrate that the AGN in the considered simulations have physical properties consistent with those of real QSOs for what the column density and extent of the gas in the host galaxies and the UV magnitudes are concerned. A bi-conical geometry for the outflow is favoured over a spherical geometry, as it reproduces AGN with the high luminosities and SMBH masses observed for $z = 6-7$ QSOs. However, both simulations cannot explain the recent discovery of luminous QSOs at $z \approx 7.5$, which may have been formed at higher redshift than the assumed seeding time in our simulations, or may have undergone extensive periods of super-Eddington accretion.

Moreover, we showed that the number of multiple AGN detectable in X-ray band over few kpc scales is the observable property that the considered simulations struggled the most to reproduce. We propose that this issue can be due to the simplistic BH seeding methods generally implemented in cosmological simulations, that do not account for the complex physics related with the formation and rapid growth of massive BHs in the early Universe. Future X-ray observatories will provide us with the sensitivity required to investigate the possible presence of multiple faint AGN satellites around luminous QSOs at high redshift.

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DATA AVAILABILITY STATEMENT

The data underlying this article will be shared on reasonable request to the corresponding author.

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