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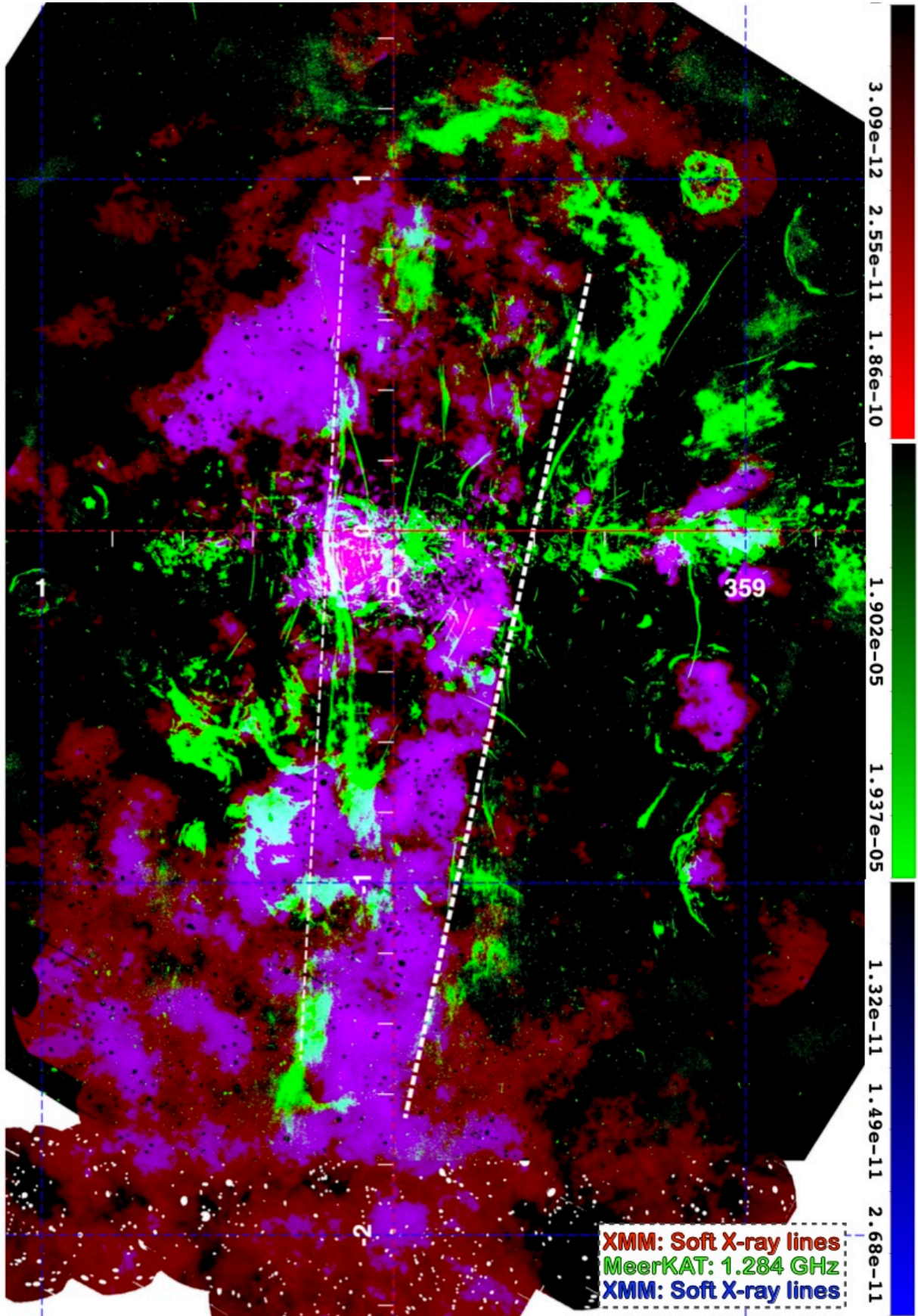


Fig. 3. Continuum-subtracted soft X-ray line image in red and blue (with different intensity cuts, to give a better indication of the extent of the chimneys). In green, the *MeerKAT* map with intensity cut at 2×10^{-5} Jy (see Fig. 1 of Heywood et al. 2019, and note the caveats regarding photometric accuracy in the Methods section).

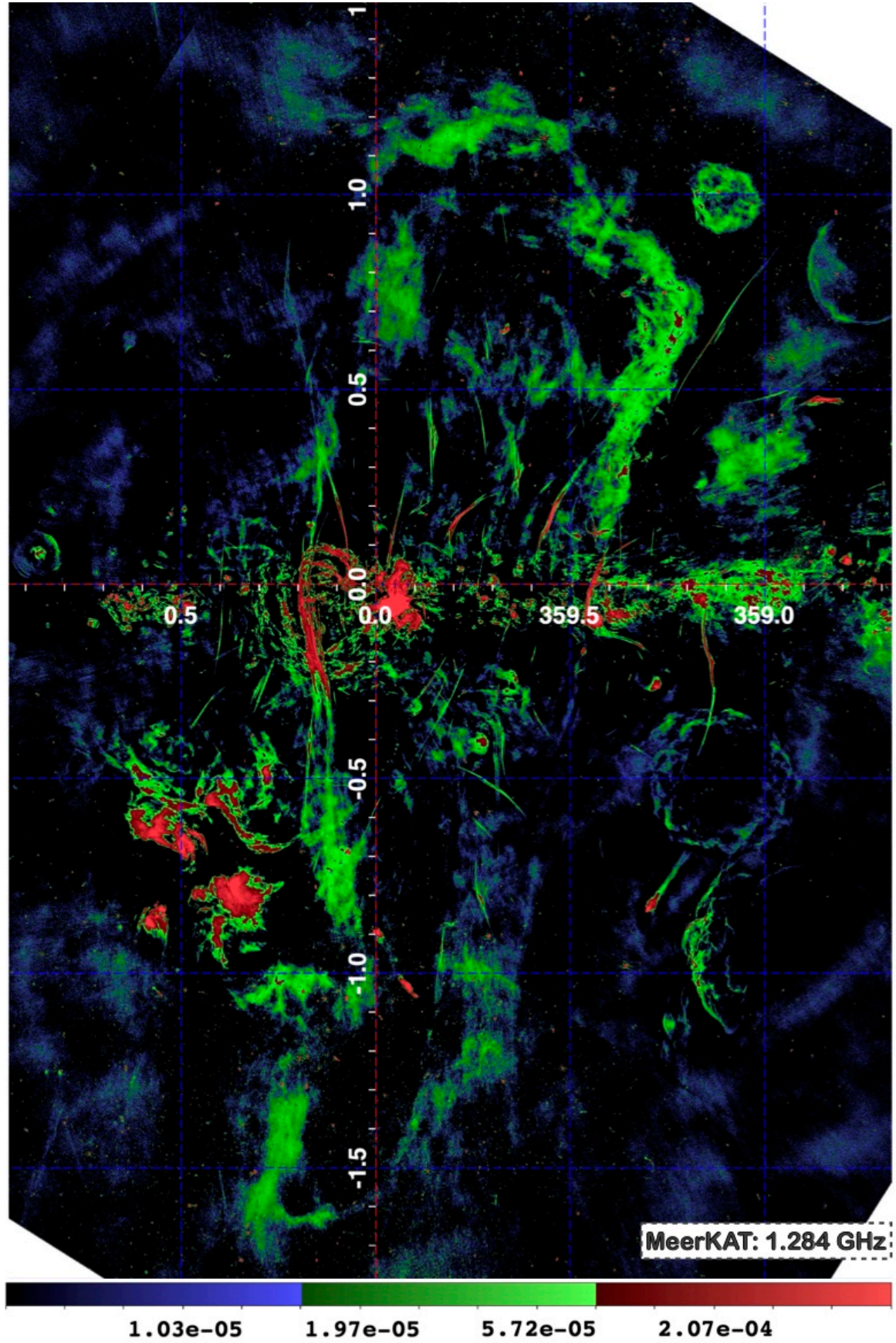


Fig. 4. *MeerKAT* map displayed with its entire dynamical range.

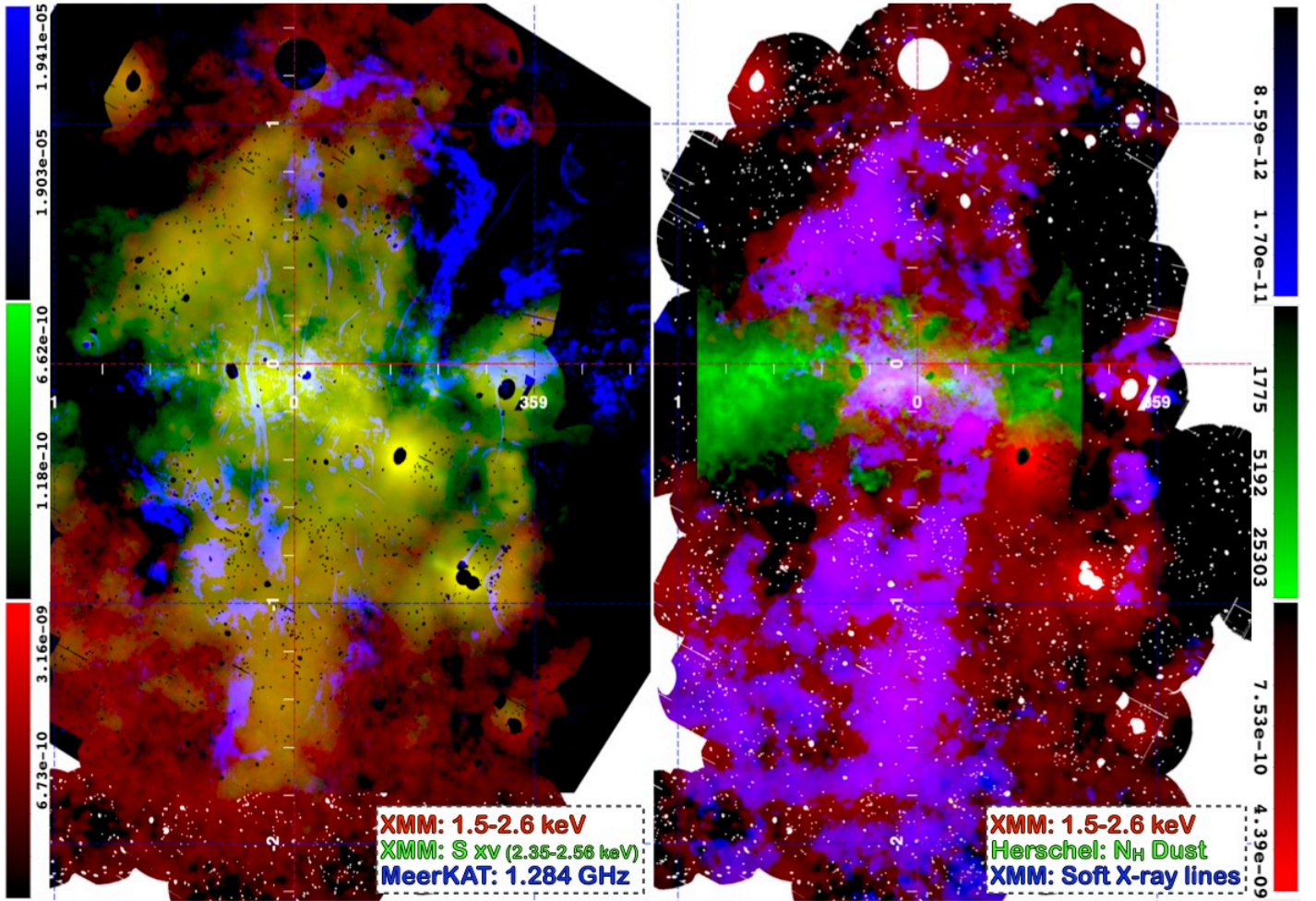


Fig. 5. Multiwavelength view of the chimneys. (*Left panel*) Red colors show the soft X-ray emission map (1.5–2.6 keV). Green colors show the S xv line emission (2.35–2.56 keV) and blue colors show the *MeerKAT* map. The S xv line emission shows clear gradients in rough coincidence with the edges of the *MeerKAT* bubbles. Additionally, it is less affected by interstellar absorption than the soft X-ray emission map (indeed the latter shows clear depressions in the emissivity of the chimneys at $b \sim 0.2^\circ$ and $b \sim -0.4 - 0.6^\circ$). (*Right panel*) RGB image of the GC. Red: Soft X-ray emission map (1.5–2.6 keV). Green: Column density map of neutral material as derived from dust emission of the central molecular zone, derived from *Herschel* observations (Molinari et al. 2011). The depressions of soft X-ray emission at $b \sim 0.2^\circ$ and $b \sim -0.4 - 0.6^\circ$ correspond to a high column density of neutral material, probably as a result of enhanced obscuration. We point out that the protrusion at $l = 0.2, b = 0.25^\circ$ (§3.6) might be associated with a fountain originating from the Sgr B molecular complex, one of the most massive, highly obscured and highest star formation rate locations of the Milky Way. Indeed, although the protrusion appears to be disconnected from the plane in the soft X-ray map (it extends from $b \sim 0.2^\circ$ to $b \sim 0.6^\circ$), this might be a byproduct of strong obscuration in the Sgr B region. Blue: Continuum-subtracted soft X-ray line emission. The continuum subtraction removes the contribution from the dust-scattering haloes around bright point sources, which otherwise dominate the emission west of the southern chimney.

are clearly observed to emerge almost vertically from the plane. The blue dashed ellipses indicate the positions of well-known, bright, foreground star-forming regions.

The bottom left panel of Fig. 6 displays an enlargement of the *WISE* map toward the northern chimney, where in red and blue are shown the 12.08 and 22.2 μm emission, respectively. The green colors show the ratio of infrared color defined as the 22.2 μm map divided by the 12.8 μm one. The bottom right panel shows in red and blue the *WISE* maps (as in the left panel) and in green the radio emission as observed by *MeerKAT*. The white dashed lines display the location of the spurs delineated on the basis of the *WISE* maps.

3. Discussion

We start the discussion by emphasizing two points.

First point: The overall degree of symmetry of the radio and X-ray emission (see Fig. 3) around the GC suggests that the chimneys are a single coherent feature located at the GC. Additionally, Fig. 1 demonstrates that the X-ray, radio and infrared emission are deeply connected. Indeed, Fig. 1 shows that they form coherent features, extending for hundreds of parsecs, that can be followed both in the X-rays, radio and the infrared bands. This demonstrates that the main players traced in each band, which are the hot plasma in X-rays, the warm dust in infrared and the shocks in radio, are interacting and deeply affecting each other. This strengthens the idea that they are all

byproducts of a single energetic phenomenon.

Second point: We note from the radio image (Fig. 3) that the surface density of well-defined nonthermal filaments is highest within ~ 0.5 degrees of the Galactic plane and drops rather abruptly with latitude beyond that. In some cases, the narrow non-thermal filaments become increasingly diffuse at higher latitudes, especially above and below the Radio Arc. In addition, at northern latitudes, the magnetic field lines delineated by the nonthermal filaments appear to diverge with increasing latitude. We conclude from these observations that the magnetic flux decreases with increasing Galactic latitude (see also Morris 2006b; 2015). Such a decrease would affect how the GC outflow is manifested as a function of latitude. At low latitudes, the shock occurring where the outflow impacts the surrounding ISM would encounter a relatively stronger field, and it would encounter it at a steep angle, so that the shock is likely to be a C-type shock in which the shock energy is distributed broadly over a relatively large region (c.f., Draine 2011), and without a velocity jump large enough to ionize the gas passing through the shock³. At high latitudes, however, with a weaker field and a more oblique shock, in the presence of a predominantly vertical field, the velocity jump could be sufficient to ionize the gas in the shock. The diffuse radio emission that we associate with the shock induced by the GC outflow as it impacts the surrounding medium would therefore appear most prominent at the higher Galactic latitudes, as is observed.

The following subsections will discuss the various components of the chimneys. Section 3.1 considers the multi-wavelength emission from the Arc super-bubble, which might be instructive for a deeper understanding of the chimneys. In §3.2 we debate the proposed foreground location of part of the northern chimney. In §3.3 we examine the northwestern edge of the chimneys, reporting evidence for a multiphase (i.e., hot and cold-molecular) and multi-epoch outflow that produces strong shocks at the edges of the chimneys. In Section 3.4 we talk about the northern "cap." In §3.5 we discuss the eastern edge of the chimneys, highlighting the differences compared with the western edge, stressing the importance of the GC magnetic field and introducing the concept of a "magnetic wall." Section 3.6 discusses possible origins of the X-ray protrusion. Section 3.7 considers the southern chimney. Section 3.8 shows that the continuity of hot plasma emission all the way to the higher latitudes, where the Fermi bubbles begin, provides evidence that the chimneys are the multi-epoch outflows that feed the Fermi bubbles with energetic particles. Finally, in §3.9 we ask whether the non-thermal radio filaments might be associated with the GC outflow.

3.1. The Arc super-bubble as a template to understand the chimneys

Figure 7 shows a radio (red), mid-infrared (green) and X-ray (blue) overlay centered on the Arc super-bubble. The super-bubble appears brightly in all of these bands. Figure 7 fully supports the paradigm that the Arc super-bubble is filled with hot plasma occupying an X-ray bright "cavity," surrounded by warm, shock-heated dust (Egan et al. 1998; Levine et al. 1999; Rodriguez-Fernandez et al. 2001; Price et al. 2001; Sofue 2003;

Simpson et al. 2007; Ponti et al. 2015). Indeed, bright mid-infrared emission is observed along the entire limb-brightened perimeter of the Arc super-bubble. The radio emission runs along the circular ridge of mid-IR emission and encloses the X-ray emission from the interior of the super-bubble. In addition, patches of bright radio emission can be observed also projected toward the interior of the Arc super-bubble (see the dotted ellipses highlighting this emission in Fig. 7), corroborating the idea that the radio emission is primarily tracing ionization fronts and also shocks at the interface of the super-bubble with the ISM. A detailed description of such an array of concentric radio shells (observed in the historical VLA 20cm data; Yusef-Zadeh & Morris 1987a,b) can be found in Sofue (2003). That work attributed the creation of the shells to recent ($\sim 10^6$ yr) starbursts, the most likely candidate being the supernovae and stellar winds from the Quintuplet cluster (Egan et al. 1998; Sofue 2003; Ponti et al. 2015).

Therefore, the Arc super-bubble represents a textbook example of a powerful outflow within the GC environment that is currently still contained within the disk of the Milky Way (Fig. 7). Although the more powerful chimneys apparently succeeded in breaking through the Galactic density gradient and in overcoming the Galactic potential, thereby opening a channel to the halo, we expect that their multiwavelength emission maintains some degree of resemblance to the Arc super-bubble. Indeed, as with the Arc super-bubble, within the chimneys we also observe intense radio emission tracing shocks occurring primarily at their edges and X-ray emission located primarily inside the chimneys as well as warm dust emission at various locations along the borders of the chimneys.

3.2. On the recent suggestion that the northwestern edge of the chimney is a foreground feature

The observation of the silhouette of the northwestern and northern portion of perimeter of the northern chimney against low-frequency radio emission (LaRosa et al. 2005; Brogan et al. 2003; Hurley-Walker et al. 2019; Tsuboi et al. 2020) clearly implies that the northern chimney is located in front of most of the diffuse radio continuum, which is produced primarily within the inner few hundred parsecs. It is currently highly debated whether the northwestern chimney is located along the Galactic disk at a few kiloparsecs from the Sun (Nagoshi et al. 2019; Tsuboi et al. 2020; Wang 2020) or whether it is instead placed just in front of the GC. The overall morphological symmetry of the radio and X-ray chimneys strongly suggests that at least some portion of the northwestern edge is located at the GC. Additionally, the association of the northwestern edge of the chimney with the AFGL 5376 feature (see §3.3), which is characterized by a very high positive velocity, indicates that both AFGL 5376 and the northwestern edge of the chimney are located at the GC. We note that, on the sky plane, the chimneys have a small overall tilt of about 7° with respect to the vertical to the Galactic plane, perhaps as a result of local pressure gradients, cloud placements, and initial injection directions. In any case, this inclination of the chimneys raises the possibility that the northern chimney is also inclined toward us to some extent. If the tilt is such that the northern chimney is inclined toward us by a few tens of degrees, then the majority of the northern chimney would be located in front of the bulk of the GC radio continuum, which could account for its appearing in absorption in low-frequency radio maps, while its southern counterpart, if the chimneys are indeed colinear, would cast no shadow on those images, as observed (more detailed discussion on this topic is presented in

³ Indeed the shock compression of the vertical field by a C-type shock at the location of the Radio Arc might play a role in generating the bundle of nonthermal filaments constituting the Arc.

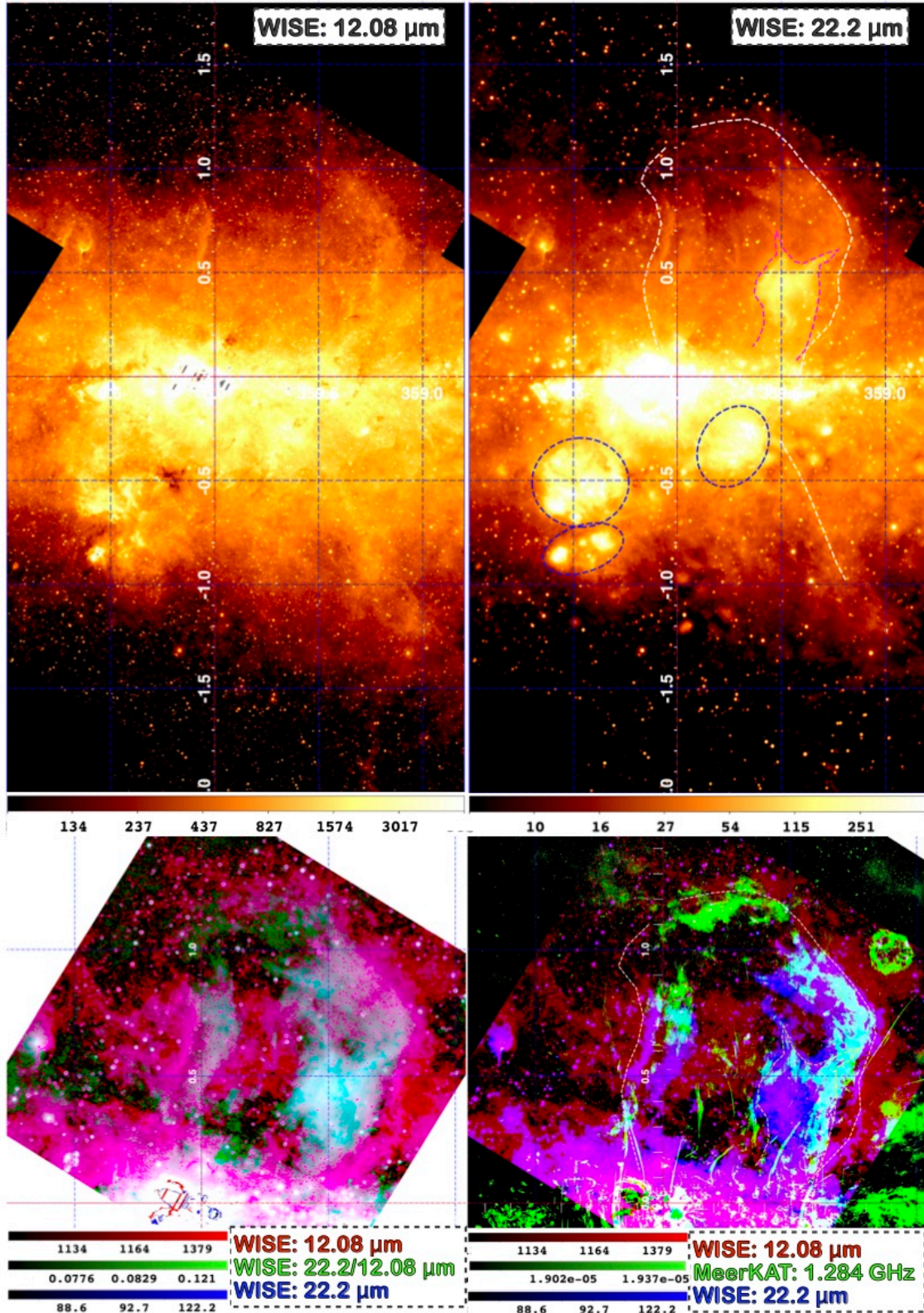


Fig. 6. Mid-IR emission from the central degrees of the Milky Way. (*Top left*) Emission at $12.08\ \mu\text{m}$ as observed by WISE. The highest concentration of warm dust is distributed along the Galactic plane. Spurs of material are clearly observed to emerge vertically from the plane. (*Top right*) Emission at $22.2\ \mu\text{m}$ as observed by WISE. The most prominent spurs of matter are highlighted by white dashed lines. At northern latitudes both the western and eastern spurs are clearly associated with the northern chimney and delineate a lobe clearly associated with the observed radio emission. The magenta dashed line delineates the silhouette of AFGL 5376. The detection of similar spurs in the southern hemisphere is complicated by the presence of well-known, bright foreground star-forming regions (highlighted by the blue ellipses). (*Bottom left*) Zoom-in toward the northern chimney. Red, green, and blue show the $12.08\ \mu\text{m}$, the ratio $22.2\ \mu\text{m}/12.08\ \mu\text{m}$, and $22.2\ \mu\text{m}$ emission, respectively. Warm dust surrounds the entire chimney. Also, the $22.2\ \mu\text{m}$ emission peaks inside the location of the western shock and of the $12.08\ \mu\text{m}$ emission spur. AFGL 5376 lies inside the edge of the warm gas and traces a $\sim 90\ \text{pc}$ long molecular shock. (*Bottom right*) Same as previous but substituting the ratio map with the radio emission observed by MeerKAT. The radio emission is very well correlated with the warm dust emission. In particular, the radio emission traces the perimeter of AFGL 5376 more tightly than the warm dust emission.

§A). Therefore, hereinafter, we will assume that the northwestern chimney is part of the chimneys that is rooted at the GC. However, we warn the reader that projection effects might be important and that a portion of the features here assumed to be located at the GC might be foreground features unrelated with the Galactic outflow.

3.3. Northwestern edge

3.3.1. Entrained molecular outflow producing a ~ 0.1 kpc long shock onto the ISM

Figure 3 shows that the western edge of the chimneys is very well defined, both in X-rays (see thick dashed line in Fig. 3 Ponti et al. 2019) and radio. Indeed, the radio emission runs parallel to the X-ray edge along its full extent of ~ 350 pc, that is, all the way from the northern part of the northern chimney to the southerly tip of the southern chimney. However, the peak of the radio emission on the western side of the chimneys is displaced toward more negative longitudes by $\sim 0.2^\circ$ (~ 30 pc) than the western edge of the X-ray emission. Additionally, Figure 6 shows that warm dust surrounds the entire northern chimney. Indeed, both the $12.08 \mu\text{m}$ and the $22.2 \mu\text{m}$ maps show a spur of material (IR1, see Fig. 2) emerging from the Galactic plane (at the location of the Sgr C molecular complex) and reaching a cap of material observed at high Galactic latitudes around $b \sim 1.1 - 1.4^\circ$ (Fig. 3, 2, 6 and 1). Again we note that, while the X-ray, radio and warm dust distributions run parallel to each other, they are significantly displaced in longitude (by $\sim 0.1 - 0.3^\circ$). This is discussed further below.

3.3.2. AFGL 5376: The brightest part of a molecular shock

Almost three decades ago, Uchida et al. (1994) studied the molecular line emission around this region and discovered two high-velocity components of molecular material, defining a vertical rift coinciding with the vertical ridge of the mid-IR source, AFGL 5376. The large velocity separation of the two molecular components — 65 km s^{-1} — led Uchida et al. (1994) to suggest that the strong IR emission from AFGL 5376 results from the energy deposited in a 90-pc long shock where the two molecular components meet (Fig. 8). Furthermore, the large positive velocities of both components are best explained by expansion motions away from the GC, so we presume that they are participating in the GC outflow⁴. The authors hypothesize the presence of two shocks, represented by the thick-solid blue lines in Fig. 2. The western blue arc would represent the leading edge of a strong ionizing shock with a westward velocity component, delineating the location where the outflow encounters the ISM (see also Fig. 8). A shock with $v \geq 65 \text{ km s}^{-1}$ could dissociate CO and H₂ molecules and ionize hydrogen, inducing free-free emission. The eastern blue line, which coincides with the IR ridge, was suggested to represent a dissociation front, associated with the reverse shock, where the fast outflowing material encounters the

⁴ The line of sight velocity of AFGL 5376 is consistent with the one expected if such a cloud were located along the innermost non-self-intersecting X1 orbit induced by the Galactic bar of stars (see e.g., Binney et al. 1991). If so, the velocity of AFGL 5376 would not be due to expansion, but rather to the streaming motions of gas along the bar, characteristic of the X1 orbits. However, such a scenario leaves two things completely unexplained: 1) the fact that AFGL 5376 is rather far out of the Galactic plane, and 2) the fact that AFGL 5376 is the site of such a strong internal shock. For these reasons, we believe that AFGL 5376 participates in the outflow from the GC region.

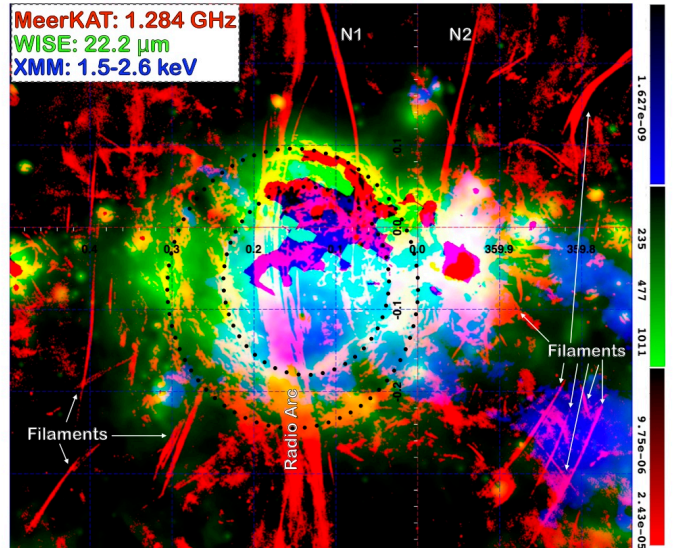


Fig. 7. Multiwavelength view of the Arc super-bubble. Red colors show the *MeerKAT* map. Green colors show *WISE* $22.2 \mu\text{m}$ emission and the blue colors show the soft X-ray emission map ($1.5\text{--}2.6 \text{ keV}$) of the region around the Arc super-bubble. The radio and mid-infrared emission are edge-brightened and trace the edges of the Arc super-bubble, while the X-ray emission appears to fill the central volume. Despite being edge-brightened, several bright stripes and arcs of radio emission, presumably tracing shocks, are projected toward the interior of the Arc super-bubble, as they are for the chimneys (some are highlighted by the two black dotted ellipses). The Radio Arc as well as many narrow, non-thermal filaments (some of which are highlighted by the white arrows) are clearly visible (see Fig. 2 for a definition of the Radio Arc and most prominent nonthermal filaments). Saturation and bright source removal in the *WISE* and *XMM-Newton* maps produce some visible artifacts.

more slowly moving post-shock material preceding it (Uchida et al. 1994).

3.3.3. Corroborating the shock-interpretation with fresh data

The data presented here corroborate this interpretation. Indeed, the *WISE* versus *MeerKAT* overlay shows intense radio emission downstream of the strong shock, on the western side of AFGL 5376 (Figs. 6 and 1). Additionally, the current data reveal that the entire perimeter of the AFGL 5376 feature is bright in radio emission, supporting the idea of a structured interaction between the outflow and the ISM, possibly through a secondary (or "reverse") shock between AFGL 5376 and the surrounding ISM (see Fig. 8). Such a reverse shock is consistent with being a dissociative shock. Indeed, the lack of radio emission running down the infrared ridge at the center of AFGL 5376, disfavor it being an ionizing shock. Additionally, the radio emission on the eastern perimeter of AFGL 5376 is at the interface between the X-ray emitting plasma and the molecular cloud, so that ionization could be induced by particle impact (e.g., conduction).

Figure 1 reveals that almost no X-ray emission is observed between the strong and reverse shock, in agreement with the idea that the shocked molecular outflow, visible as warm dust, might form a structured shock with the ISM, with the low-density, shock-heated plasma radiating in the X-ray band located in the internal post-shock region, inside of longitude $\sim 359.6^\circ$ (Fig. 1, 8). Finally, the *WISE* maps (Fig. 6) clearly show the presence of warm dust running along the edge of the strong shock proposed by Uchida et al. (1994). We point out that it remains un-

clear whether AFGL 5376 impacts the ambient ISM "head on" or whether there is considerable shear parallel to the shock.

The extraordinary result revealed by the superposition of the *XMM-Newton*, *MeerKAT* and *WISE* maps is the discovery that such an association of X-ray, radio and mid-infrared emission is not confined to the small region next to AFGL 5376, but it continues to latitudes of $b \sim 1.2^\circ$ and all the way to the cap (Fig. 1). This suggests the presence of a shock, of heated dust and swept up material accumulating at this edge of the chimneys.

The *MeerKAT* map shows that the radio emission traces the presumed location of the strong shock defining the northwestern chimney (Figs. 3 and 1). The ratio of the $22.2 \mu\text{m}$ over the $12.08 \mu\text{m}$ emission shows that the $22.2 \mu\text{m}$ emission along most of the northwestern edge, traces the radio emission at the location of the shock, as expected for shock heated dust (Fig. 6 and 1). Furthermore, the $22.2 \mu\text{m}$ emission extends for many (up to tens of) parsecs behind the shock itself, indicating a strong interaction between the GC outflow and the mechanism heating the dust.

3.3.4. Indications of stochastic heating of pre-shock dust grains

Figure 6 shows that the peak of the $12.8 \mu\text{m}$ emission along most of the northwestern edge occurs either at the presumed location (or even outside) of the strong shock defining the northwestern edge of the chimney (Fig. 6). At first sight, this appears surprising because the $12.08 \mu\text{m}$ emission typically traces warmer dust than the $22.2 \mu\text{m}$ one, leading us to the paradox of having warmer dust ahead of the shock than behind it. This might be understood if the strong, ionizing shock acts as an efficient source of the photons that stochastically heat the small dust grains, producing enhanced $12.08 \mu\text{m}$ emission, even ahead of the shock front, despite the low temperature of the larger dust grains in the pre-shock region. After the shock, all the dust gets significantly heated, not just the small grains, therefore enhancing the $22.2 \mu\text{m}$ emission in addition to the emission at $12.08 \mu\text{m}$. We point out that the mean free path of the diffuse Lyman continuum is likely too short ($\sim 0.01 \text{ pc}$ for a density of 10 cm^{-3} ; Draine 2011) to contribute to such an effect⁵. Therefore, we speculate that if the Ly α emission is produced in a region with high velocities, then the Lyman α line might be sufficiently Doppler shifted to be out of the local line profile, thus reducing the optical depth considerably. In support of such a scenario, the overlay of the *WISE* versus *MeerKAT* maps (bottom right panel of Fig. 6) shows, along the full extent of the northwestern spur (IR1) and cap, intense radio emission (tracing the shock front) peaking in the zone where the $12.8 \mu\text{m}$ start dominating over the $22.2 \mu\text{m}$ emission, with increasing distance from the center.

3.3.5. AFGL 5376: Evidence for a multiphase outflow

For a mass of AFGL 5376 of $M_{\text{AFGL 5376}} \sim 5 \times 10^4 M_\odot$, an estimated volume of $V_{\text{AFGL 5376}} \sim 18^2 \times 7 \text{ pc}^3$ and a relative velocity to the pre-shock gas of $v_{\text{AFGL 5376}} \sim 65 \text{ km s}^{-1}$, we obtain a density of $\sim 10^3 \text{ cm}^{-3}$ and a ram pressure of $p_{\text{ram}} \sim 6 \times 10^{-8} \text{ dy cm}^{-2}$, which corresponds to $\sim 40 \text{ keV cm}^{-3}$. Such pressure is about two orders of magnitude larger than the thermal pressure of the observed hot plasma within the chimneys, which has been estimated to be $p_{\text{hot}} \sim 0.1 - 0.2 \text{ keV cm}^{-3}$ (Ponti et al. 2019). Assuming that the thermal pressure of the hot plasma is accelerating AFGL 5376, we estimate that it would require $> 10^7$ years to reach a speed of 65 km s^{-1} or more. Even if the hot plasma

⁵ Even shorter mean free paths are associated with Lyman α photons.

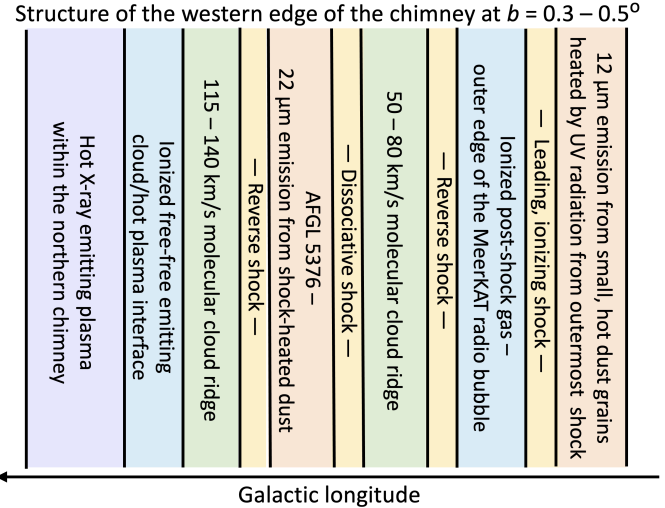


Fig. 8. Schematic representation of the structure of the northwestern edge of the chimney.

has a subsonic bulk motion as high as $v \sim 500 \text{ km s}^{-1}$, the ram pressure then results to be on the order of $\sim 0.5 \text{ keV cm}^{-3}$, therefore still requiring a long time to accelerate AFGL 5376. Therefore, this appears to disfavour models in which AFGL 5376 has been accelerated by the hot plasma that we currently observe within the chimneys, suggesting that an alternative agent might be needed to account for the outflow speed of the molecular cloud.

3.3.6. Origin of the AFGL 5376 feature

We note that the origin of AFGL 5376 might be connected with that of the recently discovered high-velocity HI clouds located at the base of the Galactic bulge (McClure-Griffith et al. 2013; Di Teodoro et al. 2018; 2020; Lockman et al. 2020). Indeed, recent large-scale HI surveys of the Galactic bulge led to the discovery of a large population (~ 200) of anomalous high-velocity HI clouds extending up to $\sim 3 \text{ kpc}$ north and south of the Galactic plane, in a bi-conical configuration, with a neutral cloud lifetime of $\sim 4 - 10 \text{ Myr}$ and with kinematics consistent with an outflow from the GC. The neutral cloud component of the outflow is characterized to have an opening angle of $> 140^\circ$, a maximum outflow velocity of $v_{\text{out}} \sim 330 \text{ km s}^{-1}$, a mass outflow rate of $\sim 0.1 M_\odot \text{ yr}^{-1}$, and a kinetic luminosity of the outflow of $L_k > 3 \times 10^{40} \text{ erg s}^{-1}$ over the past 10 Myr (McClure-Griffith et al. 2013; Di Teodoro et al. 2018; 2020; Lockman et al. 2020). For densities of such HI clouds in the range $n_H \sim 1 - 20 \text{ cm}^{-3}$ and local standard of rest velocities of the clouds in the range $v_{\text{LSR}} \sim 100 - 360 \text{ km s}^{-1}$, the driving wind should be imposing a ram pressure of $\sim 0.1 - 30 \text{ keV cm}^{-3}$ to accelerate such clouds. The upper bound on the required pressure range is one-to-two orders of magnitude larger than the thermal pressure of the hot plasma currently observed within the chimneys and similar to the ram pressure needed to accelerate AFGL 5376. We suggest that the process that accelerated the HI clouds distributed throughout a biconical volume at the base of the *Fermi* bubbles has also accelerated some massive clouds in directions closer to the Galactic plane, and is thus responsible for the dynamics of AFGL 5376.

3.3.7. Further evidence for a multiphase outflow

Other high-velocity molecular features are observed toward the chimneys, further supporting the multiphase nature of the outflow (Hsieh et al. 2015; 2016). The polar arc represents one such example. The polar arc is an extra-planar molecular cloud, located ~ 30 pc above Sgr A*, with a high radial velocity ($v > +100$ km s $^{-1}$) and a positive velocity gradient perpendicular to the Galactic plane, suggesting an accelerating and expanding motion off the plane (Bally et al. 1988; Hsieh et al. 2015; 2016). This cloud appears connected with an extra-planar hourglass-shaped feature having an extent of ~ 13 pc perpendicular to the plane and centered on the central parsec of the Milky Way, and a dynamical timescale of $\sim 3 \times 10^5$ yr (Hsieh et al. 2016). The hourglass feature runs along the edges of Sgr A's bipolar lobes containing hot plasma (see the remarkable agreement represented in Fig. 2 by the hourglass feature in blue and the edges of Sgr A's bipolar lobes in orange; Morris et al. 2003; Ponti et al. 2015). Indeed, the observed molecular components might represent the entrained molecular gas within the hot plasma of Sgr A's bipolar lobes (Hsieh et al. 2016; Ponti et al. 2015; 2019). The morphology and kinematics of these molecular clouds within tens of parsecs of the GC are reminiscent of kpc-scale molecular outflows in nearby starburst galaxies (Garcia-Burillo et al. 2001; Walter et al. 2002; Bollato et al. 2013). Indeed, they are consistent with an origin in the Galactic plane from which the clouds have been lifted. As a matter of fact, they have been suggested to have orbital paths altered either by an explosion $\sim 10^5$ yr ago or by outflows from the central parsec of the Milky Way (Hsieh et al. 2015; 2016).

A recent extensive survey of H $_3^+$ absorption toward bright stars within the central molecular zone demonstrated the presence of an outflow of warm ($T \sim 200$ K) diffuse ($n \sim 50$ cm $^{-3}$) gas (Oka et al. 2020). The diffuse, warm gas is observed to expand with radial velocities of ~ 150 km s $^{-1}$ and to be as extended as ~ 150 pc from Sgr A* (Oka et al. 2020). The energy, momentum and timescale required to create such an outflow have been estimated to be $\sim 5 \times 10^{53}$ erg, 5×10^8 M $_{\odot}$ km s $^{-1}$ and $\sim (0.5 - 1) \times 10^6$ yr (Oka et al. 2020). This revives the idea of either an expanding ring (Kaifu et al. 1972; Scoville 1972) or an expanding bi-polar vertical cylinder with total length as great as ~ 170 pc (Sofue 2017).

3.3.8. Evidence for a multi-epoch outflow

Despite the large uncertainties (e.g., line of sight distance, 3-d velocity, launching point, etc.), it appears that the molecular clouds associated with AFGL 5376 were first launched onto their relatively high-latitude trajectories a few 10^6 yr ago. We note that this timescale is comparable to the one associated with the outflow of warm diffuse gas, while it is about 10 times longer than that of the hourglass-shaped feature. Additionally, it is $\sim 10 - 100$ times longer than the sound crossing time of the chimneys ($t_s \sim 3 \times 10^5$ yr) and of Sgr A's bipolar lobes ($\sim 3 \times 10^4$ yr) and more than ten times longer than the relatively recently formed recombining plasma, with estimated age less than $\sim 10^5$ yr (Nakashima et al. 2013). We conclude that these greatly different timescales indicate that the chimneys were shaped by events occurring at widely different times, and therefore they are the product of multi-epoch events.

3.4. The cap

The superposition of the radio, mid-infrared and X-ray images reveals the presence of a cap above the northern chimney (see Figs. 3, 2, 6 and 1). This is clear evidence that the GC outflow to the north has swept up material, or is being at least partly impeded by moderately dense material there.

We speculate that the different latitudinal extent and longitudinal width of the northern and southern chimneys can be attributed to a greater initial ISM mass in the northern volume, leading to more mass swept up to create the "northern cap" and to a broader longitudinal extent because the enhanced confinement to the north causes a greater pressure that leads to a greater expansion in the longitudinal direction. This is also consistent with the fact that the AFGL 5376 cloud and the polar arc cloud are located to the north, but no comparable clouds have been seen near or in the volume occupied by the southern chimney.

We note that the X-ray maps show a clear gradient of decreasing hot plasma emission (clearer in the S xv map; Fig. 5) at the location of the cap. However, intense hot plasma emission is also observed beyond the cap (Fig. 3, 5, 1). In particular, the surface brightness, temperature, density and pressure of the hot plasma are consistent with a rather smooth transition across the cap, suggesting that the outflow associated with the last event of a quasi-continuous series of intermittent energy releases is only partially impeded by the material in the cap.

3.5. Northeastern edge

3.5.1. Differences between the eastern and western edges

Because of the absence of a structured shock similar to that seen at the northwestern boundary, the eastern edge of the chimneys is less well defined (Fig. 3). The *MeerKAT* map indicates an association of the eastern edge of the chimneys with the Radio Arc, however it also reveals that the morphology of the eastern edge, within $\sim 0.5^\circ$ of the plane, is dominated by bundles of filaments, therefore very different from the "fuzzy" appearance of the northwestern edge and cap. We also note that the filamentary appearance of the eastern edge becomes fuzzier at higher latitudes. This could be understood in terms of a magnetic field that dominates the pressure at the eastern edge of the chimneys, but starts to diverge beyond $|b| \sim 0.5^\circ$, and therefore become less confining of the synchrotron emitting particles that occupy the filaments.

Additionally, the radiative mechanism (synchrotron vs. free-free emission) and degree of polarization are two of the major differences between the eastern and western edges of the northern radio bubble, which implies that the northeastern edge is undergoing very different physical processes compared with the western edge (Fig. 3, 2, 6, 1; see also Reich et al. 1987; Haynes et al. 1992).

Finally, the morphology of the hot plasma emission along the northeastern edge represents another major difference compared with the northwestern edge. Very intense X-ray emission is observed all the way from the interior of the chimney to and beyond the location of the Radio Arc. The X-ray emission at that location is primarily thermal (with temperatures of $\sim 0.7 - 1$ keV; Ponti et al. 2019), therefore we exclude a major contribution due to synchrotron emission associated with the nonthermal filaments, although X-ray counterparts to the filaments do make a minor contribution (Wang et al. 2002, 2020 preprint, Zhang et al. 2014; Ponti et al. 2015; Mori et al. 2015).

3.5.2. Hints of a dominant vertical magnetic field

We note that on the eastern edge of the chimney, the radio map offers many manifestations of a pressure-dominant magnetic field, including: i) exceptionally long nonthermal radio filaments such as the N1, N2 and XMM J0.173-0.413 filaments, running north and south of the plane; ii) the Radio Arc; and iii) the highly polarized radio plumes running both north and south of the Galactic plane. All of these features underline the presence of a strong (likely dominating the ambient pressure) magnetic field with a strength that some have estimated to be as high as a few mG (Seiradakis et al. 1985; Morris & Yusef-Zadeh 1985; Tsuboi et al. 1986; Yusef-Zadeh & Morris 1988; Lang et al. 1999; Blanton 2008; Magilli et al. 2019).

Also, the mid-infrared spur (IR2, see Figs. 6 and 2) can be best understood in the framework of a high magnetization. We note that, close to the Galactic plane, IR2 runs parallel to the N1 filament up to $b \sim 0.4 - 0.5^\circ$, which is radiating synchrotron emission due to relativistic electrons in a highly ordered vertical magnetic field. Between $b \sim 0.4$ and $b \sim 0.7^\circ$, the mid-infrared spur is defined by the Double Helix Nebula (DHN; top left panel of Fig. 6 Morris et al. 2006; Tsuboi et al. 2010). The radio continuum emission in the general direction of the DHN is observed to be highly polarized, indicating a highly ordered magnetic field with synchrotron-emitting relativistic electrons (Tsuboi et al. 2010). Molecular line surveys of this region have revealed two molecular components at 0 and -35 km s^{-1} associated with the DHN, with no clear evidence of shocks and with turbulent line broadening of $\sim 3 - 5 \text{ km s}^{-1}$. The mass of the two molecular components is estimated at $\sim 3.3 \times 10^4$ and $\sim 0.8 \times 10^4 M_\odot$, respectively (Enokiya et al. 2014; Tori et al. 2014). These authors, following Morris et al. (2006), propose that the warm dust has been forced into such a double helix configuration by a strong magnetic field.

Torii et al. (2014) inferred densities of the DHN of $n_e \sim 0.7 - 2 \times 10^3 \text{ cm}^{-3}$. Assuming such densities and that the matter within the DHN is shaped into such a configuration by a turbulent pressure that is also the source of the line broadening, we compute the turbulent pressure to be in the range of $\sim 0.1 - 0.5 \text{ keV cm}^{-3}$. Such a pressure would be in balance with the magnetic pressure for a magnetic field with a strength of 0.1 mG. Therefore, the vertical magnetic field with a strength of $\sim 1 \text{ mG}$ (whose presence is corroborated by the presence of the nearby Radio Arc, filaments and polarized plumes) would, indeed, be able to dominate the dynamics of the DHN.

3.5.3. A magnetic wall?

We point out that if, as suggested, on the eastern side of the chimneys the magnetic field has a strength as high as $\sim 1 \text{ mG}$ and a uniform vertical configuration, then it cannot be neglected when considering the dynamics of the hot plasma⁶. Indeed, the longitudinal expansion of any outflow with a pressure lower than $\sim 100 \text{ keV cm}^{-3}$ (corresponding to a magnetic field strength of 1-2 mG) would encounter the resistance of the magnetic field, while it would allow the flow to continue almost unperturbed

perpendicular to the Galactic plane, therefore with the net effect of collimating the outflow. We also note that, adding to the collimation by the magnetic field, the vertical density gradient in the Galactic plane would contribute to the collimation. In this scenario, an outflow would be currently flowing through the chimneys. Alternatively, the hot plasma emission would be the relic of a past major outburst at the GC, now in hydrostatic equilibrium within the gravitational potential of the Milky Way and still radiating because of the very long cooling time ($2 \times 10^7 \text{ yr}$; Ponti et al. 2019). In such a scenario, it seems plausible that the hot plasma was generated and then remained at the edge of the outflow.

Despite its pivotal importance, this scenario fails to explain what is sustaining the magnetic field. Several authors have connected the creation of the strong (in the mG range), vertical GC magnetic field to the accretion of plasma through the Galactic disk over the entire life of the Milky Way (Sofue et al. 1987; 2010; Sofue & Fujimoto 1987; Howard & Kulsrud 1997; Chandran et al. 2000). In such scenarios, the toroidal component of the field would be amplified by differential rotation, and toward the GC, radial compression would amplify the vertical field, creating a nearly vertical magnetic field because of inefficient ambipolar diffusion (Chandran et al. 2000). In particular, Sofue et al. (2010) showed that the winding of the primordial magnetic field can evolve into composite configurations, comprising bisymmetric spiral, axisymmetric spiral, plane-reversed spiral, and/or sing fields in the disk, and vertical fields in the center, similar to what is observed in nearby spiral galaxies. We speculate that, if the magnetic field is mass-loaded, then its more relevant role on the eastern side of the GC might be associated with the larger amount of mass on the Galactic plane at positive longitudes (Fig. 5 Bally et al. 1988; Tsuboi et al. 1999; Molinari et al. 2011; Jones et al. 2012). Additionally, the effect of the magnetic field would appear more evident than on the northwestern side due to the fact that the ISM appears less dense above the plane on that side.

3.6. Protrusion

The X-ray map shows bright X-ray emission at $(l, b) \sim (0.4, 0.4^\circ)$, which appears as a "protrusion" from the northern lobe (Fig. 3 and 1 Ponti et al. 2019). The origin of such emission is unclear.

The physical properties of the hot plasma (e.g., temperature, density, energetics) within the protrusion could be consistent with an interloper supernova remnant (SNR). To validate this hypothesis, we note that the *XMM-Newton* and *MeerKAT* maps show at least four SNRs at Galactic latitudes comparable to or higher ($|b| > 0.3^\circ$) than that of the protrusion, therefore confirming the high projected density of SNRs in this region. Additionally, we note that, along the line of sight toward the Radio Arc a crescent-shaped radio feature is observed at the western border of the protrusion, at $(l, b) = (0.07^\circ, 0.18^\circ)$, which is consistent with being a radio shell (Figs. 4, 3 and 1). Furthermore, radio emission is observed around most of the remaining portions of the protrusion, although it is weak (compare Figs. 3 and 4), and not definitively associated with the protrusion. Therefore, the lack of a clear radio counterpart in either *MeerKAT* or previous radio surveys leaves the association of the protrusion with an SNR unclear (Haynes et al. 1992; LaRosa et al. 2000; Heywood et al. 2019).

Figure 5 shows that the S xv emission (which is less affected by absorption than the soft X-ray emission map), extends the emission of the protrusion all the way to the Galac-

⁶ Such a scenario has been already considered by several authors. Sofue (2020a), for example, considers a magnetic cylinder as the origin of the chimneys and nonthermal filaments. He hypothesizes that magnetohydrodynamic compression waves ejected from the nucleus might be reflected and guided through the magnetic field, therefore appearing as the nonthermal filaments (when viewed tangentially) and produce feedback loops between accretion events and magnetic outflows Sofue (2020b).

tic plane at the location of the Sgr B1 and Sgr B2 molecular complexes. Indeed, the Sgr B molecular complex represents the region of the Milky Way with the highest specific star formation rate ($\text{SFR}_{B2} \sim 0.04 \text{ M}_{\odot} \text{ yr}^{-1}$; Armillotta et al. 2019). Therefore, based on these data, we do not exclude the possibility that the protrusion might represent either a super-bubble or the early phases of the formation of a Galactic fountain, which could contribute to energizing and excavating the chimneys. Forthcoming deep radio maps of the protrusion hold the key to understanding the origin of such a peculiar hot plasma feature.

3.7. Southern chimney

The study of the southern chimney is complicated by the presence of well-known bright foreground star forming regions down to latitudes of $\sim 0.5^{\circ}$ and $\sim 0.7 - 0.9^{\circ}$ (highlighted by the blue ellipses; Fig. 2). Indeed, ambiguity remains regarding how much of the observed radio, X-ray, and mid-infrared emission is associated with foreground features.

The western side of the southern chimney appears as an impressive extension of the northwestern edge (Figs. 1 and 3). Indeed a single straight line well represents the border between the hot plasma distribution and the radio emission. This morphological correspondence suggests a direct link between the two in the form of a structured shock such as its northern counterpart, although not as clearly defined. The lack of an evident mid-infrared spur of warm dust and of shocked molecular clouds might be the consequence of its much shorter cooling time as well as of smaller ISM densities present at these negative latitudes (c.f., discussion in §3.5.1).

The eastern side of the southern chimney is defined by the extension of the Radio Arc down to $b \sim -1^{\circ}$. Also here, the morphology of the Radio Arc changes from an array of thin filaments to a fuzz of emission, consistent with a dominant magnetic field with degrading flux density away from the plane, as a consequence of the diverging magnetic field with latitude (Fig. 4 and 3).

The entire southern chimney appears to be filled with hot plasma and surrounded by radio emission. This is even more evident at its southernmost extension (Figs. 3, 5 and 1). Again this is easily understood as the aftermath of an outflow carving its way out and accumulating ISM material at the edges. We note that hot X-ray emitting plasma leaks beyond the southernmost radio edge. As for the northern cap, this implies that material at the edge is likely only partially impeding the observed outflow. Additional support for the "escaping hot plasma" scenario comes from the detection of recombination lines in its spectrum at this location (Nakashima et al. 2013). Those authors estimate an expansion time-scale on the order of $\sim 8 \times 10^4 \text{ yr}$, with an upper limit of $\sim 1.1 \times 10^5 \text{ yr}$ (Nakashima et al. 2013). We note that recombining plasma is often associated with plasma expanding on a time-scale shorter than that required to reach ionization equilibrium. Therefore, the claimed presence of recombining plasma could be readily understood if some portion of the outflow were unimpeded in its expansion by the radio shell.

3.8. High-latitude hot plasma: Further evidence for a multi-epoch outflow

We observe that both the northern and southern chimneys merge at high latitudes with cooler, X-ray-emitting plasma, which is consistent with the X-ray emission observed in the *ROSAT* map

and attributed to the plasma around the edges of the *Fermi* bubbles (Ponti et al. 2019).

The physical conditions of the high-latitude hot plasma and of the chimneys support the idea that the chimneys represent the channel replenishing the *Fermi* bubbles with energy and particles and that the current detailed morphology of the chimneys reflects the most recent episodes of energy injections.

3.9. Apparent association of the most prominent nonthermal radio filaments with the GC outflow

Since their discovery in the eighties, it has been suggested that the nonthermal radio filaments are tracers of an intense (dominating), pervasive and vertical magnetic field with an intensity of $\sim 1 \text{ mG}$ (Morris & Serabyn 1996). The nonthermal filaments would then appear anywhere there is a source of relativistic particles that illuminate the tube of field lines in which they are trapped (Morris & Serabyn 1996). The recent discovery of groups of nonthermal filaments spatially organized to resemble "harp," provides considerable credibility to such an interpretation (Thomas et al. 2020). Several other interesting processes have also been invoked to explain the origin of the filaments (Lesch & Reich 1992; Serabyn & Morris 1994; Rosner & Bodo 1996; Shore & LaRosa 1999; Bicknell & Li 2001; Yusef-Zadeh 2003; 2019; Bykov et al. 2017; Sofue 2020a).

Does the GC outflow play a role in creating the nonthermal filaments? We note that the overlay of the *WISE* and *MeerKAT* maps with the location of the strong shock suggested by Uchida et al. (1994) shows that the bright nonthermal filament C16 might originate at this location, although it could be coincidental (Figs. 4, 3 and 1). Indeed, in theory, strong shocks are expected to enhance the magnetic field strength and to accelerate particles, the two ingredients necessary to illuminate the filaments.

In agreement with this scenario, we note that a large fraction of the nonthermal filaments are indeed observed either within or just outside of the chimneys (Fig. 2). This raises the possibility that a fraction of these magnetic filaments might have been produced in a strong-field ($B \sim 0.1 - 1 \text{ mG}$) boundary zone surrounding the chimneys, perhaps by the compression of the GC magnetic field at that interface by the high gas pressure associated with the outflow. Alternatively, some filaments might be due to the (subsonic) stretching of a buoyant fluid element that will stretch the magnetic field, creating adjacent field lines of opposite polarity in its wake, which then gives rise to magnetic reconnection that accelerates particles to sufficient energies to shine via synchrotron emission in the radio band (see e.g., Churazov et al. 2013). Of course, if the nonthermal filaments are distributed within a cylindrical region surrounding the chimneys, then we should expect that some of them would be projected toward the interior of the chimneys as is observed.

4. Emerging picture

Despite the complexity of the region—where projection effects are certainly relevant, where several competing physical effects might be at play, and where obscuration limits the breadth of our observing windows—a picture appears to emerge from this multi-wavelength study. Indeed, the comparison of the X-ray, radio and infrared emission of the chimneys demonstrates the high level of interconnection of the emission in these various bands. Indeed, the same coherent features, extending for hundreds of parsecs, can be followed in all three bands. This demonstrates that the main players traced in each band, which are the hot plasma in X-rays, the warm dust in infrared and the shocks and nonthermal

filaments in the radio, are interacting and deeply affecting each other. This strengthens the idea that they are the byproduct of the same phenomenon that drives them all.

These multiwavelength data support the idea that the chimneys represent the channel connecting the quasi-continuous, but intermittent activity at the GC with the base of the *Fermi* bubbles. The prominent edges observed in the radio band are then the signposts of the most powerful and recent outflow from the central parsecs, which has created the radio shell partially filled with X-ray emitting plasma and swaths of warm dust at the boundary.

On the relation between the activity currently observed at the GC and the outflow of hot plasma which shaped the chimneys.

As shown by Ponti et al. (2019), the steep pressure gradient within Sgr A's bipolar lobes indicates that there is an outflow of hot plasma within the central tens of parsecs. However, this small-scale hot outflow is not the driver of the outflow reaching the top of the chimneys.

On the northwestern side of the chimneys, we observe that the pressure associated with the molecular outflow created a strong shock possessing a pressure that is about two orders of magnitude larger than the thermal pressure of the hot plasma (if volume filling) within the chimneys⁷. Again, on the northeastern side, the pressure appears to be dominated by the vertical magnetic field with a field strength of $\sim 0.1 - 1$ mG. These comparisons suggest that either the hot outflow is supersonic or the chimneys were inflated and sculpted by an entity that is different from the currently observed hot plasma.

In the latter case, the observed hot plasma might trace only a small fraction of the energy of the more powerful outflow which would be the driver of the observed phenomenology. Several candidates, including cosmic rays (Breitschwerdt et al. 1991; Yusef-Zadeh et al. 2019), Alfvén or MHD waves (Sofue 2020a,b), fast-and-cold outflow resulting after rapid adiabatic expansion (Chevalier & Clegg 1985; Heckman et al. 1990; Suchkov et al. 1994; Krumholz et al. 2017), very hot plasma, among others have been proposed. We stress that the presence of a fast-and-cold outflow is often invoked in starburst galaxies, where the initially hot plasma is rapidly adiabatically expanding, therefore transforming into a fast and cold flow (Heckman et al. 1990; Suchkov et al. 1994; Krumholz et al. 2017). On the other hand, if indeed an enhancement of very hot plasma ($kT \sim 7 - 10$ keV) is present within a degree of the Galactic plane at the GC, then this might be the driver of the outflow (Yamauchi et al. 1993; 2008; Koyama et al. 1986; but see Revnitsev et al. 2009).

The "relic" outflow scenario.

Alternatively, the outflow might have been significantly more powerful in the past, while its pressure has now dropped by orders of magnitude, although some pressure is still maintained by continued, lower-level activity. In such a scenario, hot plasma (with a significantly higher pressure than the one observed today) was previously flowing through the chimneys, then leaving the chimneys (on a sound crossing time of $t \sim 3 \times 10^5$ yr) to energize the *Fermi* bubbles. The currently observed hot plasma within the chimneys would still be radiating because of its long cooling time and it would likely be in hydrostatic equilibrium with the Galactic potential. In this scenario, the relic hot plasma within the chimneys would currently drive no strong shocks. Therefore, in agreement with the multiwavelength observations,

we would observe strong shocks at the boundaries of the chimneys only where they are driven by the slower moving phases of the GC outflow (i.e., AFGL 5376), which are still tracing the older, more powerful activity. Additionally, in this relic outflow scenario, the high magnetic field strength (with $\sim 0.1 - 1$ mG) might have been close to equipartition with the much more powerful past outflow. This picture appears in line with an active galactic nucleus-like type of activity, where the power of the outflow can change by orders of magnitudes on a relatively short time-scale. Clearly, a lower level of ongoing nuclear activity would still be required to produce features such as Sgr A's bipolar lobes, etc. In all of these scenarios, the formation of the northern cap and southern partially "enclosing" radio emission is likely associated with the bursty nature of the source powering the chimneys.

Interchanging dominant terms in the pressure balance and the role of the magnetic field.

The vertical magnetic field, present in the diffuse component of the GC, is likely to play an important role in collimating the outflow from the plane. At the same time, the compression and stretching of the fluid, induced by the outflow at the edges of the chimneys, will enhance the magnetic field, likely facilitating the generation of nonthermal filaments. The toroidal distribution of molecular material at the GC (Fig. 5) must also have played a role in collimating the outflow. Indeed, it has been suggested that the $\sim 5 \times 10^7 M_\odot$ of molecular mass within the central ~ 200 pc of the plane (Dahmen et al. 1998; Molinari et al. 2011) might be able to redirect out of the plane even a powerful active galactic nucleus outflow (Zubovas et al. 2011).

The total pressure driving the outflow is likely composed of several contributors. The data suggest that the magnetic pressure might be the dominant term in several places at the GC. Indeed, we ascribe the morphological difference between the eastern and western sides of the chimneys to a difference in the balance between the dominance in the pressure of matter and of magnetic field.

Testing these scenarios.

Future radio and X-ray observations of the GC will soon allow an assessment of the global framework proposed here and will provide insights to some of the many questions left open by this work. The most immediate question relates to the true location of the northwestern edge of the chimney. If this edge were located just in front of the GC as we have assumed, it would be consistent with the scenario proposed here. Were it instead located in the foreground Galactic disk, then it would have serious implications for the coherence of the north-south chimneys, and radio and infrared emission from the vicinity of the AFGL 5376 feature would then need to be sorted out from that arising in the foreground.

Additionally, high-resolution radio polarization and spectral index observations should eventually yield a much clearer distinction between thermal and nonthermal emission regions, thereby pointing to where the magnetic field plays an important role. Also, a more detailed X-ray and radio investigation of the nature of the Protrusion feature is needed to determine whether it is an offshoot of the northern chimney, perhaps created by some directed energy flow from Sgr A*, or an unrelated but superimposed plasma volume created by supernovae in the central molecular zone. Finally, many fundamental things will be learned about the GC outflow when future X-ray spectrometers (such as *Athena*) can provide information on the velocity field of the hot plasma.

⁷ Unless the hot outflow is moving at supersonic speeds ($v \gg 500$ km s⁻¹), also the ram pressure of the hot outflow is one-to-two orders of magnitude smaller than the pressure of the molecular outflow.

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Appendix A: Comments on the recent suggestion that the northwestern edge of the chimney is a foreground feature

As discussed in §3.2, the line of sight distance toward the northwestern edge of the chimney is still highly debated. Indeed, recently several authors have proposed that the northwestern edge of the chimneys is a foreground feature located within $\sim 5\text{--}6$ kpc from us (Nagoshi et al. 2019; Tsuboi et al. 2020). These claims have been formulated on the basis of four points:

First: The surprisingly low radial velocities (from -4 to $+10$ km s $^{-1}$; Law et al. 2009; Nagoshi et al. 2019) of the radio recombination lines (RRL) along the edges of the northern chimney;

Second: The observation of the silhouette of the northern chimney against low-frequency emission;

Third: The observation of H α emission associated with the top of the northwestern chimney;

Fourth: The distribution of the visual extinction from the Digitized Sky Survey.

Indeed, Tsuboi et al. (2020) proposed that the northwestern edge of the chimney corresponds to a giant H II region in the Galactic disk. Additionally, Wang (2020) pointed out that the ionized plasma inside such an H II region might be the source of the scattering layer observed toward many GC sources, most notably the GC magnetar (Mori et al. 2013; Rea et al. 2013; Eatough et al. 2013; Bower et al. 2014).

While the case for a location of the northwestern edge of the chimney in front of the main emission from the GC appears solid, its location at only few kiloparsecs from us leaves a few open issues.

First: The first and most striking is related with the good degree of north-south symmetry of the chimney. In particular, the western edge of the chimney can be traced very well both in the X-ray and radio maps all the way from $b \sim -1.5$ to $\sim 0.8^\circ$ (Fig. 3). If the northwestern edge of the chimney is located only a few kiloparsecs away, the striking continuity of the northern and southern chimneys would likely imply that the southern counterpart is also relatively local, and the challenge becomes explaining such a morphologically peculiar feature. In that case, it is not obvious what could produce such an energetic feature with such a unique morphology, and it would have to be a truly remarkable coincidence that such a feature is centered on the GC.

Second: The northwestern edge of the chimney appears intimately related with AFGL 5376. AFGL 5376 is the most prominent part of a molecular cloud moving at very high positive velocities that strongly indicate its GC location. AFGL 5376 cannot be placed anywhere but in the central molecular zone (Serabyn & Morris 1994). Its large positive velocity ($v \gg 100$ km s $^{-1}$) has no known counterpart in foreground material, and there is no obvious mechanism to accelerate an individual cloud from the velocity range previously observed in this direction (from -60 km s $^{-1}$ to $+20$ km s $^{-1}$) up to such a high positive velocity.

Third: The northeastern side of the chimney is nestled up against the radio Arc and its extended filaments and polarized plumes, which are clearly located at the GC.

We note that one of the most relevant arguments to place the northwestern side of the chimney at few kpc from us are the low radial velocities of the RRL. We note that RRL can be inverted and produce weak maser amplification and that the low density and temperatures in the surface layers of foreground molecular

clouds in the Galactic disk could produce such inversions. In this scenario, the radial velocities of the RRL are characteristic of the foreground clouds, therefore they do not necessarily imply that the northwestern edge of the chimney must be located in the Galactic disk. We leave the detailed discussion of this scenario to a future publication (Morris et al. in prep.).

The observation of the silhouette of the northern chimney against low-frequency emission definitely places most of the northern chimney in front of the bulk of the low-frequency radio continuum produced at the GC. We propose that both the northern and southern chimneys are rooted at the GC, but they are significantly tilted toward us, such that the northern chimney is in front of the bulk of the radio emission arising in the GC.

The Galactic latitudes at which H α is seen are relatively large, so if the GCL is at the GC, the extinction at those latitudes might not be so large that the H α line would be completely extinguished. Therefore, the presence of H α is not a convincing argument for the foreground hypothesis (see Fig. 5 of Law 2010).

While these arguments do not address all the open issues (e.g., the visual extinction toward the northwestern edge, which supports the hypothesis that the northwestern edge of the chimneys is located in the foreground), we assume that the chimney is a coherent structure, with a significant tilt, and rooted at the GC. However, we caution the reader that some portion of the features that in projection appear superimposed on the chimney might have a different origin.