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LETTER TO THE EDITOR

The evolution of luminous red nova AT 2017jfs in NGC 4470[★]

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ABSTRACT

We present the results of our photometric and spectroscopic follow-up of the intermediate-luminosity optical transient AT 2017jfs. At peak, the object reaches an absolute magnitude of $M_g = -15.46 \pm 0.15$ mag and a bolometric luminosity of 5.5×10^{41} erg s⁻¹. Its light curve has the double-peak shape typical of luminous red novae (LRNe), with a narrow first peak bright in the blue bands, while the second peak is longer-lasting and more luminous in the red and near-infrared (NIR) bands. During the first peak, the spectrum shows a blue continuum with narrow emission lines of H and Fe II. During the second peak, the spectrum becomes cooler, resembling that of a K-type star, and the emission lines are replaced by a forest of narrow lines in absorption. About 5 months later, while the optical light curves are characterized by a fast linear decline, the NIR ones show a moderate rebrightening, observed until the transient disappears in solar conjunction. At these late epochs, the spectrum becomes reminiscent of that of M-type stars, with prominent molecular absorption bands. The late-time properties suggest the formation of some dust in the expanding common envelope or an IR echo from foreground pre-existing dust. We propose that the object is a common-envelope transient, possibly the outcome of a merging event in a massive binary, similar to NGC 4490–2011OT1.

Key words. binaries: close – stars: massive – supernovae: individual: AT 2017jfs – supernovae: individual: NGC 440–2011OT1 – stars: winds, outflows

1. Introduction

Red Novae (RNe) form a family of optical transients spanning an enormous range of luminosities. This includes faint objects with absolute peak magnitudes M_V from -4 to -6.5 mag, such as OGLE 2002-BLG-360 (Tyndel et al. 2013) and V1309 Sco (Mason et al. 2010; Tyndel et al. 2011), intermediate-luminosity events ($M_V \gtrsim -10$ mag) like V838 Mon (Munari et al. 2002; Goranskij et al. 2002; Kimeswenger et al. 2002; Crause et al. 2003), and relatively luminous objects such as NGC 4490–2011OT1 (Smith et al. 2016), that can reach $M_V \approx -15$ mag. Objects brighter than $M_V = -10$ mag are conventionally named luminous red novae¹ (LRNe; for a review, see Pastorello et al. 2019, and references therein).

Although the physical processes triggering these outbursts have been debated, there is growing evidence that RNe and their more luminous counterparts are produced by the coalescence of stars with different masses following a common-envelope phase (e.g., Kochanek et al. 2014; Pejcha et al. 2016, 2017; MacLeod et al. 2017, 2018). In particular, the inspiralling motion of the secondary was revealed by the long-term monitoring of V1309 Sco (Tyndel et al. 2011).

[★] Table A.1 is only available at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via <http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/625/L8>

¹ The alternative naming “luminous red variable” was also used in the past (e.g., Martini et al. 1999).

The recent discovery of LRNe suggests that common envelope ejections and/or merging events may also happen in more massive close binary systems (Smith et al. 2016; Mauerhan et al. 2018), with major implications for the evolution of the resulting merger. In this context, here we report the results of our follow-up campaign of a LRN recently discovered in the galaxy NGC 4470: AT 2017jfs.

2. AT 2017jfs, its host galaxy, and reddening

AT 2017jfs² was discovered by *Gaia* on 2017 December 26.13 UT (MJD = 58113.13, Delgado 2007) at a *Gaia* G-band magnitude 17.17 ± 0.20 . No source was detected at the transient position on 2017 November 30 down to a limiting magnitude of 21.5. The transient was located at $\alpha = 12^h 29^m 37.79^s$ and $\delta = +07^\circ 49' 35''.18$ (equinox J2000.0), in the almost face-on, early-type spiral galaxy NGC 4470.

The source was tentatively classified by the extended-Public ESO Spectroscopic Survey for Transient Objects (ePESSTO; Smartt et al. 2015) as a type-IIIn supernova (SN IIIn) or a SN impostor (Bufano et al. 2018), and for this reason it was designated with a SN name (SN 2017jfs). In this paper, we show it to be a LRN, hence we adopt the label AT 2017jfs.

² The object is known by multiple survey designations, including *Gaia*17dtkh, PS17fqp, ATLAS18aat.

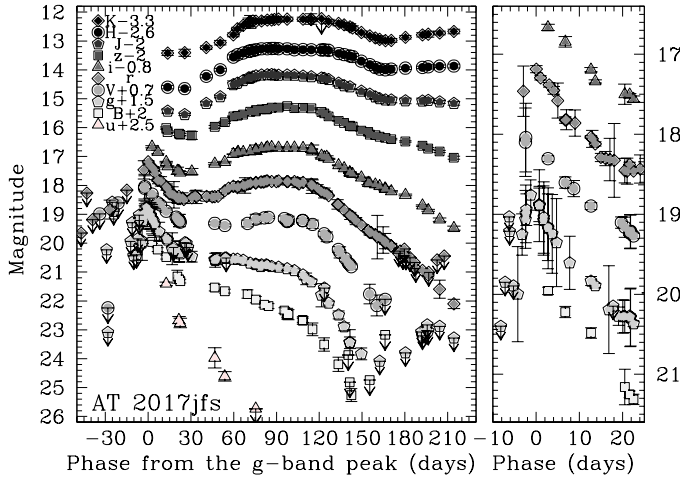


Fig. 1. *Left:* optical and NIR light curves of AT 2017jfs. The phases are calculated from the first g -band peak, on MJD = 58114.8 ± 1.8 . Only the most significant detection limits are shown. *Right:* B , V , g , r and i light curves during the first (blue) peak.

The distance to NGC 4470 is somewhat controversial, and the NASA/IPAC Extragalactic Database (NED)³ gives a number of discrepant estimates based on the Tully-Fisher method, ranging from 11.6 to 34.6 Mpc, with an average value of 18.76 ± 6.6 Mpc (corresponding to a distance modulus $\mu = 31.25 \pm 0.70$ mag). Given the uncertainty in the Tully-Fisher estimates, we prefer to adopt a kinematic distance $d = 35.2 \pm 2.7$ Mpc (corrected for Virgo Infall and estimated adopting a standard cosmology with $H_0 = 73 \text{ km s}^{-1} \text{ Mpc}^{-1}$), hence $\mu = 32.73 \pm 0.15$ mag. This estimate also agrees with the distance $d \sim 34.7$ Mpc reported by Koliopanos et al. (2017).

The Galactic line-of-sight reddening is modest, $E(B - V) = 0.022$ mag (Schlafly & Finkbeiner 2011). Our early spectra do not have high signal to noise ratios (S/Ns), and therefore the host galaxy reddening cannot be well constrained. However, prominent absorption features of Na ID are not visible, suggesting a modest host galaxy reddening contribution. Later spectra with higher S/N show only a narrow Na ID in absorption centered at $5884 \text{ \AA}$ (rest wavelength), hence likely a feature intrinsic to AT 2017jfs. For this reason, we adopt $E(B - V) = 0.022$ mag as the total reddening towards AT 2017jfs.

3. Photometric evolution

The follow-up campaign started soon after the classification of AT 2017jfs, and continued for about 7 months. Photometry data were reduced following standard prescriptions (see, e.g., Cai et al. 2018), using the SNOoPy package (Cappellaro 2014). The magnitudes are listed in Table A.1, available at the CDS, which contains the following information: Column 1 lists the date of the observation, Col. 2 lists the MJD, Cols. 3–11 give the optical and near-infrared (NIR) magnitudes, and Col. 12 reports a numeric code for the instrumental configuration. The multi-band light curves are shown in Fig. 1. The Sloan- u light curve shows a monotonic decline after maximum, with an average rate of $6.5 \pm 0.9 \text{ mag (100 d)}^{-1}$. The photometric evolution in the other bands is somewhat different. The g -band maximum is constrained to MJD = 58114.8 ± 1.8 (2017 December 27.8 UT) through a low-order polynomial fit (at $g = 17.35 \pm 0.02$ mag, hence $M_g = -15.46 \pm 0.15$ mag).

The g -band light curve has a rise time to maximum of about 4 d, followed by a rapid decline ($6.5 \pm 0.2 \text{ mag (100 d)}^{-1}$) until ~ 50 d. The light curve follows a plateau-like evolution until ~ 110 d when it begins a faster decline ($7.0 \pm 0.4 \text{ mag (100 d)}^{-1}$) that lasts until it has faded below the detection threshold. The evolution in the Johnson B and V bands is similar to the g band, although the V -band light curve shows a low-contrast second peak, broader than the early one.

The light curve is remarkably different in the red and NIR bands. The transient reaches a peak at $r = 17.19 \pm 0.05$ mag on MJD = 58115.6, followed by a fast decline ($6.3 \pm 0.3 \text{ mag (100 d)}^{-1}$) lasting three weeks and reaching a minimum at $r = 18.32 \pm 0.07$ mag. Subsequently, from about 50 d after maximum, the r -band luminosity rises again and reaches a second maximum on MJD = 58209.0, at $r = 17.68 \pm 0.03$ mag. This second peak is much broader than the early one. Later on, from 110 to 200 d after the first peak, the r -band light curve declines with a rate of $3.90 \pm 0.04 \text{ mag (100 d)}^{-1}$. The i -band light curve is very similar, with the two peaks reaching comparable luminosities.

The NIR light curves have a second maximum, brighter than the early one. As an additional feature, GROND (Greiner et al. 2008) observations reveal a moderate rebrightening of the NIR light curves from ~ 170 to 220 d. Although we do not have very late spectroscopic observations to support this (Sect. 4), a late-time NIR luminosity excess can be associated with the formation of new dust or IR echoes, occasionally observed in LRNe at late phases (see, e.g., Banerjee et al. 2015; Exter 2016). The late NIR brightening may also be a consequence of the transition to the brown (L-type) supergiant stage, as happened for V838 Mon (Evans et al. 2003; Munari et al. 2007), although this scenario does not comfortably explain the late blue-shift of the $H\alpha$ emission observed in the late spectra of AT 2017jfs (see Sect. 4).

4. Spectral evolution

We collected 14 epochs of optical spectroscopy, spanning about six months of the evolution of AT 2017jfs. Information on the instrumental configurations is given in Table A.2. Our spectral sequence of AT 2017jfs is presented in Fig. 2, while the comparison with a few LRNe at similar epochs and the line identification are shown in Fig. 3. We remark that the transient lies in a crowded region of NGC 4470, rich in nearby sources. As a consequence, the late spectra show some contamination from host galaxy lines.

The spectral evolution of AT 2017jfs follows a three-phase behavior, as observed in other extra-galactic LRNe. In particular, we note a remarkable similarity with NGC 4490–2011OT1 (Smith et al. 2016; Pastorello et al. 2019) at all phases. At early epochs (until 3–4 weeks after the first g -band peak) the spectrum of AT 2017jfs shows a blue continuum, dominated by prominent H lines in emission, with a Lorentzian profile and a full width at half maximum velocity $v_{FWHM} \sim 700 \text{ km s}^{-1}$ (corrected for spectral resolution). In this period, the temperature inferred from a black-body fit to the spectral continuum, T_{BB} , decreases from about $7800 \pm 700 \text{ K}$ (in the +9.5 d spectrum) to $6000 \pm 600 \text{ K}$ (in the +23.4 d spectrum). Emission lines from a number of Fe II multiplets are also detected, along with O I. The Ca II NIR triplet is also identified in emission, while the H&K feature, which is usually prominent in absorption in other LRNe (see Fig. 3, top panel), is marginally detected in AT 2017jfs.

With time, the continuum becomes redder and the spectrum experiences an evident metamorphosis. During the second peak, from ~ 50 d to 4 months, the red spectrum ($T_{bb} = 4300 \pm 700 \text{ K}$ at 82.3 d) is dominated by a forest of metal lines in absorption.

³ <https://ned.ipac.caltech.edu/>

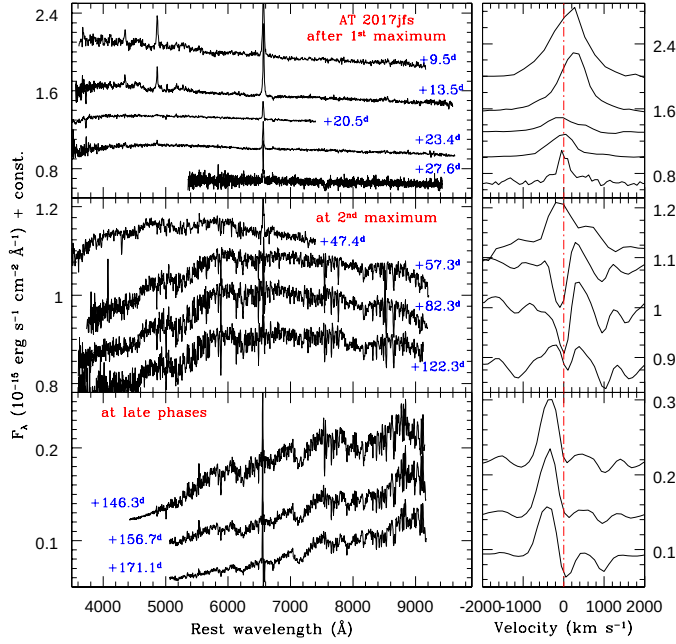


Fig. 2. *Left:* spectral evolution of AT 2017jfs soon after the first blue maximum (*top*), around the second, red maximum (*center*), and at late phases (*bottom*). *Right:* expanded image of the $H\alpha$ region in the velocity space. The vertical line marks the rest wavelength position of $H\alpha$ at the same phases as the left panels (after correcting for the host galaxy redshift $z = 0.007809$ as reported in NED). Only the highest-S/N spectra are shown. The absorptions visible at velocities 0 and $+1000 \text{ km s}^{-1}$ in late spectra can be partly due to an over-subtraction of the unresolved $H\alpha$ and $[\text{N II}] \lambda 6583$ from the host galaxy background.

The H lines become much weaker, showing now a P Cygni profile (see Fig. 2, left panel), although an over-subtraction of the narrow $H\alpha$ emission from a nearby H II region may affect the apparent strength of the absorption. The spectrum in this phase is reminiscent of intermediate-type stars (e.g., late G to K types). We identify a number of metal lines (from Fe II, Ti II, Sc II and Ba II multiplets), along with the Na I 5889, 5895 Å doublet (see Fig. 3, mid panel). The O I and the NIR Ca II triplet are now seen in absorption. The velocity of the narrow Fe II lines deduced from the wavelengths of absorptions is about 450 km s^{-1} . Some of the absorption lines visible at this stage are likely due to neutral metals, in particular Fe I at red wavelengths.

From about 5 months after maximum (hence during the steep, late luminosity decline; see Sect. 3), the spectrum changes again, becoming much redder ($T_{\text{bb}} \approx 2950 \pm 150 \text{ K}$ at 157 d) and closer to that of an M-type star. $H\alpha$ is now mostly in emission, with an evident blue-shift of its peak (by about 400 km s^{-1} ; see Fig. 2, right), and a deep absorption at the rest velocity. While an over-subtraction of the contaminant H II region can be responsible in part for this absorption, the strong blue-shift of the emission is real. The $H\alpha$ profile is similar to that of LRN NGC 4490–2011OT1 about 200 days after maximum (Smith et al. 2016), whose bluest emission peak was shifted by -280 km s^{-1} . Following Smith et al. (2016), the development of a blue-shifted component in emission would be consistent with an expanding, shock-heated line-forming region, possibly with aspherical symmetry. A peculiar geometry or, alternatively, the formation of dust hiding the rear emitting region (or both) may explain the late $H\alpha$ profile for both NGC 4490–2011OT1 and AT 2017jfs. Higher-resolution spectra and a wider temporal monitoring would help in discriminating the different scenarios.

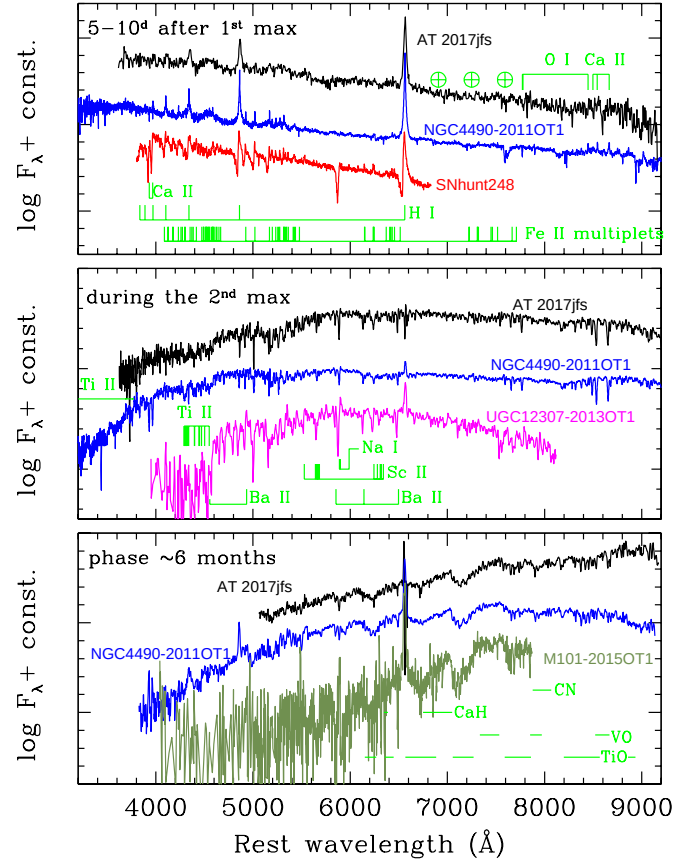


Fig. 3. Line identification in the spectra of AT 2017jfs and a few comparison LRNe at three representative phases: a few days after the blue peak (*top panel*), around the red peak (*middle panel*), and at late phases (about 5–6 months after the blue peak; *bottom panel*). The spectra of the comparison objects are taken from Pastorello et al. (2019), Blagorodnova et al. (2017), and Kankare et al. (2015). The identification of the molecular bands is performed following Kirkpatrick et al. (1991), Valenti et al. (1998), Martini et al. (1999), and Barsukova et al. (2014).

As observed in similar transients (e.g., Pastorello et al. 2019), the late-time spectrum is also characterized by broad absorption bands. The features are generally identified as being due to molecules, in particular TiO and VO, although CN and CaH are not ruled out (see Fig. 3, top panel).

5. Evolution of the temperature and the radius

The double-peaked light-curve evolution and the major spectroscopic transition from an SN IIn-like spectrum to that of a late-type star are two remarkable properties of LRNe (Pastorello et al. 2019). The photometric information in particular can be used to study how the spectral energy distribution (SED), the effective temperature, and the photospheric radius evolve with time.

To this aim, the SED is computed for a few representative epochs, and the observed data are fitted by a single black-body function. The early-time (near the blue maximum) SEDs do not contain u , i , z -band and NIR observations, while u -band data are not available from about two months after peak. Finally, the fluxes in the blue optical bands are not available at very late phases, because the object was below the detection thresholds in those filters. The resulting black-body fits are shown in Fig. 4 (left panel). Until about three months after maximum,

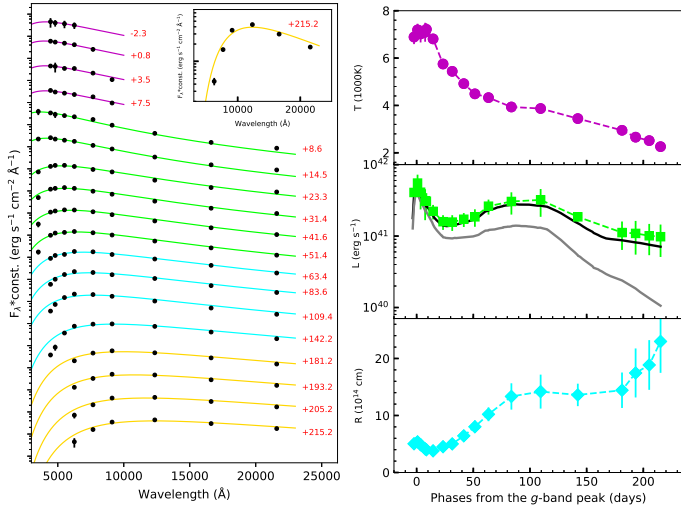


Fig. 4. *Left:* evolution of the SED at some selected epochs spanning the entire evolution of AT 2017jfs. *Top-right:* evolution of the effective temperature. *Middle-right:* bolometric light curve of AT 2017jfs (green squares and dashed line), compared with the *uvoir* (black solid line) and *opt* (gray solid line) pseudo-bolometric curves (see text). *Bottom-right:* evolution of the photospheric radius of AT 2017jfs.

observations are well modeled by black-body fits, although from ~ 3 – 4 weeks the line-blanketed *u*-band sits below the adopted models. After the red peak, a single black body is not sufficient to accurately represent the observed SEDs in the blue region (see the inset in the left panel of Fig. 4). This happens when the NIR light curves of AT 2017jfs start a new rise before the object is in heliacal conjunction. This is possibly due to the contribution of a second black-body component peaking at longer wavelengths that cannot be properly fitted because of the inadequate wavelength coverage of our observed SED, in particular towards the mid- and far-infrared domains. The nature of this putative cold component is unclear. It is possibly due to an IR echo from distant pre-existing dust or, more likely, to the condensation of newly formed dust, as suggested by the early appearance of molecular bands in the spectra and the strong blue-shift of the late $H\alpha$ emission (Sect. 4).

The evolution of the effective temperature is shown in the top-right panel of Fig. 4. The temperature remains roughly constant at about 7000 K during the blue peak. Soon after maximum, the temperature declines very rapidly, reaching ~ 4500 K at 50 d. Later on, the temperature fades more slowly, down to about 2300 K at 215 d, although this value is uncertain, as it was inferred from a poor, single black-body fit (see inset in Fig. 4, left panel).

The temporal evolution of the bolometric luminosity of AT 2017jfs, inferred by integrating the black-body fluxes over the entire wavelength range, is shown in Fig. 4 (mid-right panel), and is compared with the pseudo-bolometric curves obtained by accounting for the contribution of the optical plus NIR bands (*uvoir*), and the optical bands (*opt*) only. For the first peak we obtain a bolometric luminosity $L_{\text{bol}} \sim 5.5 \times 10^{41} \text{ erg s}^{-1}$, which is comparable to those of other intermediate-luminosity optical transients (Berger et al. 2009; Soker & Kashi 2012) or faint core-collapse SNe (Pastorello et al. 2004; Spiro et al. 2014). After the post-peak decline (with a minimum of $L_{\text{bol}} \sim 1.5 \times 10^{41} \text{ erg s}^{-1}$), the bolometric light curve rises to the second peak with $L_{\text{bol}} \sim 3.2 \times 10^{41} \text{ erg s}^{-1}$, and then declines again. After ~ 150 d, in coincidence with the late NIR brightening, the bolometric light curve flattens to $L_{\text{bol}} \sim 10^{41} \text{ erg s}^{-1}$. We note that

the bolometric luminosity of the blue peak in AT 2017jfs is twice that of the red peak. This is a major difference with NGC 4490–2011OT1, where the red peak was twice as luminous as the early blue peak (see Pastorello et al. 2019, their Fig. 11). This discrepancy is likely a consequence of the large UV contribution during the blue peak that was not accounted for in the pseudo-bolometric light curve of NGC 4490–2011OT1.

Using the Stefan-Boltzmann law, with the luminosities and temperatures estimated above, we infer the evolution of the radius at the photosphere for AT 2017jfs (Fig. 4, bottom-right panel). The radius R at blue peak slightly exceeds $5 \times 10^{14} \text{ cm}$ ($R \approx 7700 R_{\odot}$). After a modest decline, the radius rapidly increases reaching $\sim 19000 R_{\odot}$ at about 80 d, and then remains roughly constant until ~ 6 months. During the last month of the monitoring campaign of AT 2017jfs, we observe a further fast increase in the photospheric radius, which exceeds $33000 R_{\odot}$ at 215 d. This rise in the photospheric radius and the dramatic decline of the effective temperature at very late phases favor the formation of new dust, like in RN V838 Mon (Bond 2003). This is also consistent with the blueshift of the $H\alpha$ emission shown in Fig. 2 (left).

6. Discussion and conclusions

Pastorello et al. (2019) presented optical data for a wide sample of extra-galactic LRNe, all of them showing double-peaked light curves with maximum absolute magnitudes M_V in the range -12.5 to -15 mag. They discuss the observational similarity of LRNe with fainter ($M_V > -10$ mag) RNe discovered in the Milky Way, and agree with Kochanek et al. (2014) and Smith et al. (2016) that all these transients are explained in a similar binary-interaction framework. Most likely, they result from stellar merging events that occurred after the ejection of the common envelope. Lipunov et al. (2017) and MacLeod et al. (2017) discuss the structured light curves of LRNe. A plausible scenario for the double-peak light curve of AT 2017jfs invokes an initial mass outflow as a consequence of the merging event, followed by a later interaction with the common envelope. This would produce the first luminosity peak and the spectra resembling those of type-IIIn SNe. During the second peak, the photospheric radius ($\sim 2 \times 10^4 R_{\odot}$) is likely coincident with that of the ejected common envelope. With the temperature decline, the H recombines, and the released radiation determines the broad red maximum. According to Metzger & Pejcha (2017), the double-peak light curve of LRNe is explained with a modest mass ejection following the coalescence, with the early peak being due to the release of thermal energy from the fast ejecta in free expansion along the polar axes. The late red peak would result from shock-powered emission in the collision between the fast shell and pre-existing material in the equatorial plane. This would also generate a cool dense shell, which is an ideal site for late dust formation, as likely observed in AT 2017jfs. Barsukova et al. (2014) provided a somewhat different interpretation. The rapid coalescence generates a violent forward shock which leads the photospheric temperature to largely increase, producing the blue light curve peak. This phase is followed by the fast adiabatic expansion of the envelope with thermal energy carried out with some delay to the outer layers producing the broad red maximum.

Kochanek et al. (2014) proposed that the wide range of peak luminosities observed in RN/LRN events (over 4 orders of magnitudes in luminosity) is tightly connected with the total mass of the binary system, with faint RNe having progenitor systems of the order of $1 M_{\odot}$ and intermediate-luminosity events like V838 Mon of $\lesssim 10 M_{\odot}$. Luminous transients such as AT 2017jfs,

SNhunt248 (Mauerhan et al. 2015; Kankare et al. 2015), and NGC 4490–2011OT1 (Smith et al. 2016; Pastorello et al. 2019) likely arise from more massive binaries (up to 50–60 $M_{\odot}$, Mauerhan et al. 2018). While for AT 2017jfs we do not have any direct information on the progenitor system and the pre-outburst light-curve evolution, its luminous light curve would favor a massive binary as precursor of AT 2017jfs.

A possible correlation between outflow velocities and light-curve peak luminosities for merger candidates is presented in Mauerhan et al. (2018), but includes intermediate-luminosity red transients similar to SN 2008S and M85-OT whose nature is debated (e.g., Botticella et al. 2009; Kasliwal et al. 2011; Kulkarni et al. 2007; Pastorello et al. 2007). Since for AT 2017jfs we measure an expansion velocity $v_{FWHM} \approx 700 \text{ km s}^{-1}$ and $\log(L_{\text{peak}}/L_{\odot}) \approx 7.2$, the object is positioned very close to NGC 4490–2011OT1 in their Fig. 12, hence supporting the parameters trend discussed in Mauerhan et al. (2018).

While RNe from relatively low-mass stars are expected to be quite common, luminous events are extremely rare (Kochanek et al. 2014). In particular, following Kochanek et al. (2014, their Fig. 3), AT 2017jfs-like events would occur at a rate of $<10^{-4} \text{ yr}^{-1}$ within 1 Mpc. Therefore, within a volume of 40 Mpc in radius, we should find about three events per year, which is roughly consistent with observations. In fact, while we observed at least four RNe with $M_V \gtrsim -10$ mag in the Milky Way in the past two decades (V4332 Sgr, V838 Mon, V1309 Sco and OGLE-2002-BLG-360), LRNe brighter than $M_V < -10$ mag were never discovered in our Galaxy, with only less than ten objects observed within 40 Mpc in the past few years.

Due to the limited number of objects discovered so far and incomplete data sets, RNe/LRNe are still not fully understood. Well-sampled, multi-band light curves extending to longer wavelengths and high-S/N spectra with good resolution are essential tools for improving their characterization. Discovering new LRNe at larger distances and RNe outside the Local Group is crucial for understanding the physics of these objects, and for providing reliable intrinsic rates. These are key objectives of the Large Synoptic Survey Telescope (LSST Science Collaboration 2009) and other future-generation surveys.

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Appendix A: Additional data

Table A.2. General information on the spectra of AT 2017jfs.

Date	MJD	Phase (days)	Instrumental configuration	Exptime (s)	Res (Å)	Range (Å)
2018 Jan 6	58124.31	+9.5	NTT+EFOSC2+gm13	600	18	3650–9200
2018 Jan 10	58128.26	+13.5	NOT+ALFOSC+gm4	3600	18	3600–9650
2018 Jan 17	58135.33	+20.5	NTT+EFOSC2+gm11	2 × 3600	14	3350–7450
2018 Jan 20	58138.16	+23.4	NOT+ALFOSC+gm4	3600	14	3500–9700
2018 Jan 24	58142.41	+27.6	P200+DBSP+gt316/7500	1200	5.5	5400–9500
2018 Feb 13	58162.22	+47.4	NTT+EFOSC2+gm11	2 × 3600	14	3350–7450
2018 Feb 23	58172.10	+57.3	GTC+OSIRIS+R1000B+R1000R	1800 + 1800	7.0 + 7.8	3850–9250
2018 Mar 20	58197.10	+82.3	GTC+OSIRIS+R1000B+R1000R	1800 + 1800	7.0 + 7.8	3650–9350
2018 Apr 22	58230.94	+116.1	BTA+SCORPIO+VPHG550G	2 × 1200	13	3650–7900
2018 Apr 29	58237.05	+122.3	GTC+OSIRIS+R1000B+R1000R	2 × 900 + 2 × 900	7.0 + 7.8	3650–9200
2018 May 23	58261.06	+146.3	GTC+OSIRIS+R1000B+R1000R	2 × 1600 + 1600	7.0 + 7.8	3650–9250
2018 Jun 1	58270.98	+156.2	GTC+OSIRIS+R1000R	2 × 1800	7.8	5100–9350
2018 Jun 2	58271.97	+157.2	GTC+OSIRIS+R1000R	2 × 1800	7.8	5100–9350
2018 Jun 16	58285.93	+171.1	GTC+OSIRIS+R1000R	4 × 1500	7.8	5100–9250

Notes. The phases are from the *g*-band maximum. NTT = 3.58 m New Technology Telescope (ESO-La Silla, Chile); NOT = 2.56 m Nordic Optical Telescope (La Palma, Canary Islands, Spain); P200 = 5.1 m (200-inch) Hale Telescope (Mt. Palomar, California, USA); GTC = 10.4 m Gran Telescopio Canarias (La Palma, Canary Islands, Spain); BTA = 6.05 m Bolshoi Teleskop Alt-azimutalnyi (Special Astrophysical Observatory, Karachay-Cherkessian Republic, Russia).