



Publication Year	2016
Acceptance in OA	2020-04-30T16:08:37Z
Title	Return to [Log-]Normalcy: Rethinking Quenching, The Star Formation Main Sequence, and Perhaps Much More
Authors	Abramson, Louis E., Gladders, Michael D., Dressler, Alan, Oemler, Augustus, Jr., POGGIANTI, Bianca Maria, Vulcani, Benedetta
Publisher's version (DOI)	10.3847/0004-637X/832/1/7
Handle	http://hdl.handle.net/20.500.12386/24396
Journal	THE ASTROPHYSICAL JOURNAL
Volume	832



RETURN TO [Log-]NORMALCY: RETHINKING QUENCHING, THE STAR FORMATION MAIN SEQUENCE, AND PERHAPS MUCH MORE

LOUIS E. ABRAMSON¹, MICHAEL D. GLADDERS², ALAN DRESSLER³, AUGUSTUS OEMLER, JR.³,
 BIANCA POGGIANTI⁴, AND BENEDETTA VULCANI⁵

¹ Department of Physics & Astronomy, UCLA, 430 Portola Plaza, Los Angeles, CA 90095-1547, USA; labramson@astro.ucla.edu

² Department of Astronomy & Astrophysics, and Kavli Institute for Cosmological Physics, The University of Chicago, 5640 South Ellis Avenue, Chicago, IL 60637, USA

³ The Observatories of the Carnegie Institution for Science, 813 Santa Barbara Street, Pasadena, CA 91101, USA

⁴ INAF-Osservatorio Astronomico di Padova, Vicolo Osservatorio 5, I-35122 Padova, Italy

⁵ School of Physics, The University of Melbourne, VIC 3010, Australia

Received 2016 March 29; revised 2016 August 31; accepted 2016 September 1; published 2016 November 10

ABSTRACT

Knowledge of galaxy evolution rests on cross-sectional observations of different objects at different times. *Understanding* of galaxy evolution rests on longitudinal interpretations of how these data relate to individual objects moving through time. The connection between the two is often assumed to be clear, but we use a simple “physics-free” model to show that it is not and that exploring its nuances can yield new insights. Comprising nothing more than 2094 loosely constrained lognormal star formation histories (SFHs), the model faithfully reproduces the following data it was not designed to match: stellar mass functions at $z \leq 8$; the slope of the star formation rate/stellar mass relation (the SFR “Main Sequence”) at $z \leq 6$; the mean sSFR ($\equiv \text{SFR}/M_*$) of low-mass galaxies at $z \leq 7$; “fast-” and “slow-track” quenching; downsizing; and a correlation between formation timescale and sSFR (M_*, t) similar to results from simulations that provides a natural connection to bulge growth. We take these findings—which suggest that quenching is the natural downturn of all SFHs affecting galaxies at rates/times correlated with their densities—to mean that: (1) models in which galaxies are *diversified* on Hubble timescales by something like initial conditions rival the dominant grow-and-quench framework as good descriptions of the data; or (2) absent spatial information, many metrics of galaxy evolution are too indiscriminating—if not inherently misleading—to confirm a unique explanation. We outline future tests of our model but stress that, even if ultimately incorrect, it illustrates how exploring different paradigms can aid learning and, we hope, more detailed modeling efforts.

Key words: galaxies: evolution – galaxies: formation – galaxies: luminosity function, mass function – galaxies: star formation

1. INTRODUCTION

Measurements of galaxy sizes, masses, and star formation rates (SFRs) now span much of cosmic time. Yet, core uncertainties regarding how to interpret these data hinder the construction of a definitive physical narrative of galaxy evolution. This issue lies at the center of a debate between two schools of thought in this field:

1. Galaxy growth via in situ star formation is a uniform phenomenon *interrupted* by internal or external “quenching” processes (e.g., Peng et al. 2010; Steinhardt & Speagle 2014).
2. Galaxy growth is a heterogeneous phenomenon *diversified* by even more fundamental processes of which quenching is symptomatic, but not necessarily informative. (Tinsley 1968; Gladders et al. 2013; Kelson 2014).

This distinction is not semantic, but axiomatic: It reflects different ideas about the meaning of the data and ultimately leads to different basic questions in galaxy evolution—“What stops star formation?” or “What shapes star formation histories?” Hence, it also affects where and how one looks for the key physical mechanisms regulating galaxy growth.

A prolific body of work describes the quenching paradigm and its associated mechanisms, including: starvation from external gas supplies; heating of gas in dark matter halos; various galaxy-galaxy interactions; and feedback from

supernovae and/or active galactic nuclei (AGNs; e.g., Larson et al. 1980; Dekel & Silk 1986; Moore et al. 1999; Kereš et al. 2005, 2009; Springel et al. 2005; Croton et al. 2006; Oppenheimer et al. 2010; Hopkins et al. 2014, 2016).⁶

Quenching has the virtue of posing a precisely framed question: What agents truncate a galaxy’s otherwise “ordinary” life process of continual growth via star formation? A drawback is that many candidates act on scales too small to simulate with current techniques (e.g., Hopkins et al. 2014), have many possible parameterizations leading to widely varying effects (e.g., Elahi et al. 2016; Knebe et al. 2015), and rely on phenomena that are typically difficult to observe.

Such issues and the results of our own analyses have led us to an alternative paradigm. Over the course of a number of papers—Dressler et al. (2013), Oemler et al. (2013), Gladders et al. (2013), Abramson et al. (2013, 2014, 2015)—we have developed a framework in which galaxy star formation histories (SFHs; records of in situ stellar mass (M_*) production) are not maintained or squelched, but *extended* or *compressed*, rising and falling smoothly over time. The cessation of star formation is thus an extension of the *slowing* of star formation,

⁶ “Quenching” first described abrupt truncations of arbitrary levels of star formation (e.g., (post-)starbursts; Harker et al. 2006). The term has broadened to encompass gradual processes that push star formation to low levels before it is shut off (e.g., starvation; Peng et al. 2015). We argue that the latter are better thought of as “normal” (uninterrupted) galaxy evolution, but accept the usage because it is so widespread.

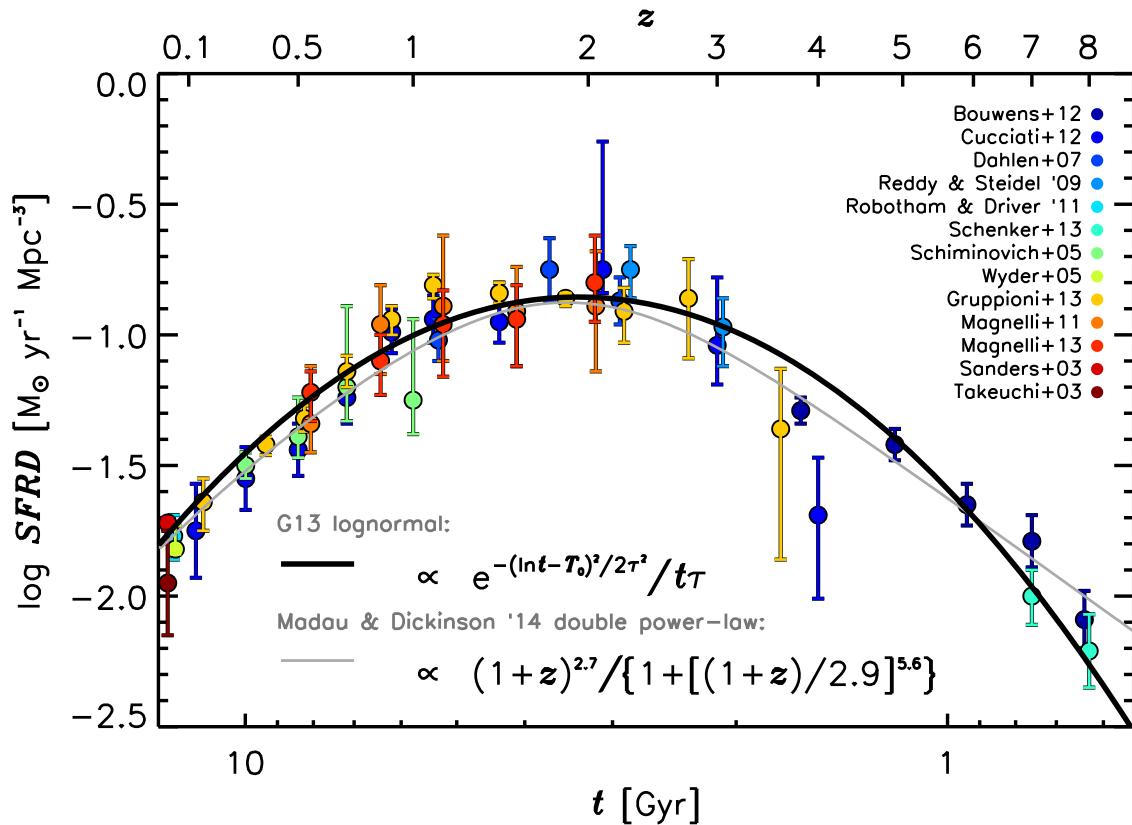


Figure 1. Cosmic $\text{SFRD}(t)$ is well-described by a simple lognormal in time. Data points are from the studies listed as compiled by Madau & Dickinson (2014). The thin gray line is the best-fit double power-law from that text. The thick black line is the best-fit lognormal: $(T_0, \tau)_{\text{Uni}} = (1.64, 0.66)$, or $(5.16, 1.93)$ Gyr (Equation (1); $A_{\text{Uni}} = 0.96 M_{\odot} \text{Mpc}^{-3}$). Due to fitting revised data, these values differ immaterially from those in G13 (Appendix A, Figure 16). While the two fits are of similar quality, the lognormal has one fewer free parameter.

and so the natural conclusion of the galaxy lifecycle, not the untimely termination of an otherwise healthy existence.

Seen in this light, focus shifts from local quenching mechanisms to some more-fundamental driver of *diversity* in growth histories. Provided the term can be defined to encompass gradual declines in star formation activity, we do not dismiss quenching *per se*, but suggest it to be a manifestation of this deeper organizing principle (Section 4.1).

Here, we develop this argument by exploring the simple mathematical description of galaxy evolution detailed in (Gladders et al. 2013, hereafter G13). Though it neither is nor is meant to be a substitute for physical explanations, we suggest that its success in reproducing suites of data it was not designed to match either recommend it as a viable *paradigmatic description* of galaxy evolution, or question the utility of a number of cornerstone observations in this field.

Below, Section 2 outlines the motivation and construction of the G13 model. Section 3.1 compares it to key observations— $\langle \text{sSFR}(z) \rangle$, $d \log \text{SFR} / d \log M_*$, stellar mass functions, “fast-” and “slow-track” quenching, the UVJ diagram, downsizing—and demonstrates its success at matching most of them. Sections 3.2 and 3.3 dissect the model to see why it behaves the way it does, concluding that (1) its interpretation of the scatter in the $\text{SFR}-M_*$ relation as *Hubble timescale* SFH diversity, and (2) the physical emphasis it places on SFH *widths* are critical. Section 4 presents these findings’ broader implications, namely that (1) “grow-and-quench” is an unnecessarily narrow framework of galaxy evolution, (2) initial conditions may be highly predictive of a galaxy’s SFH, or (3) different kinds of data than

are currently available are needed to show otherwise. Section 5 summarizes.

Readers interested in our core arguments can jump to Sections 3.2, 3.3, or 4. They might also see Appendices A–C, which flesh out key ideas presented in those Sections. Readers less interested in historical context can skip Section 2.1.

We begin by revisiting the inception of the G13 model and its achievements to date.

2. TOWARD AN EXPLICITLY QUENCHING-FREE WORLDVIEW

Of the papers mentioned in Section 1, (Oemler et al. 2013, hereafter O13) and G13 are most germane to the present discussion. These works summarize the motivation and construction of the G13 model, respectively.

2.1. Motivation: A Diversity of Smooth SFHs

O13 reached two conclusions based on analyses of specific star formation rate ($\text{sSFR} \equiv \text{SFR}/M_*$) distributions at $z \lesssim 1$.

The first was that canonical SFH parameterizations— τ and delayed- τ models (Tinsley 1972; Gavazzi et al. 2002)—could not explain the global decline in sSFRs observed over the above interval: they cannot generate both the tail of high sSFRs at $z = 1$ and the low values seen today. This problem—central to understanding how G13 bypasses explicit quenching (Section 3.2.1, Appendix A)—is not small: $\sim 25\%$ of $z \approx 0$ galaxies with $M_* \gtrsim 4 \times 10^{10} M_{\odot}$ —a fair fraction of the

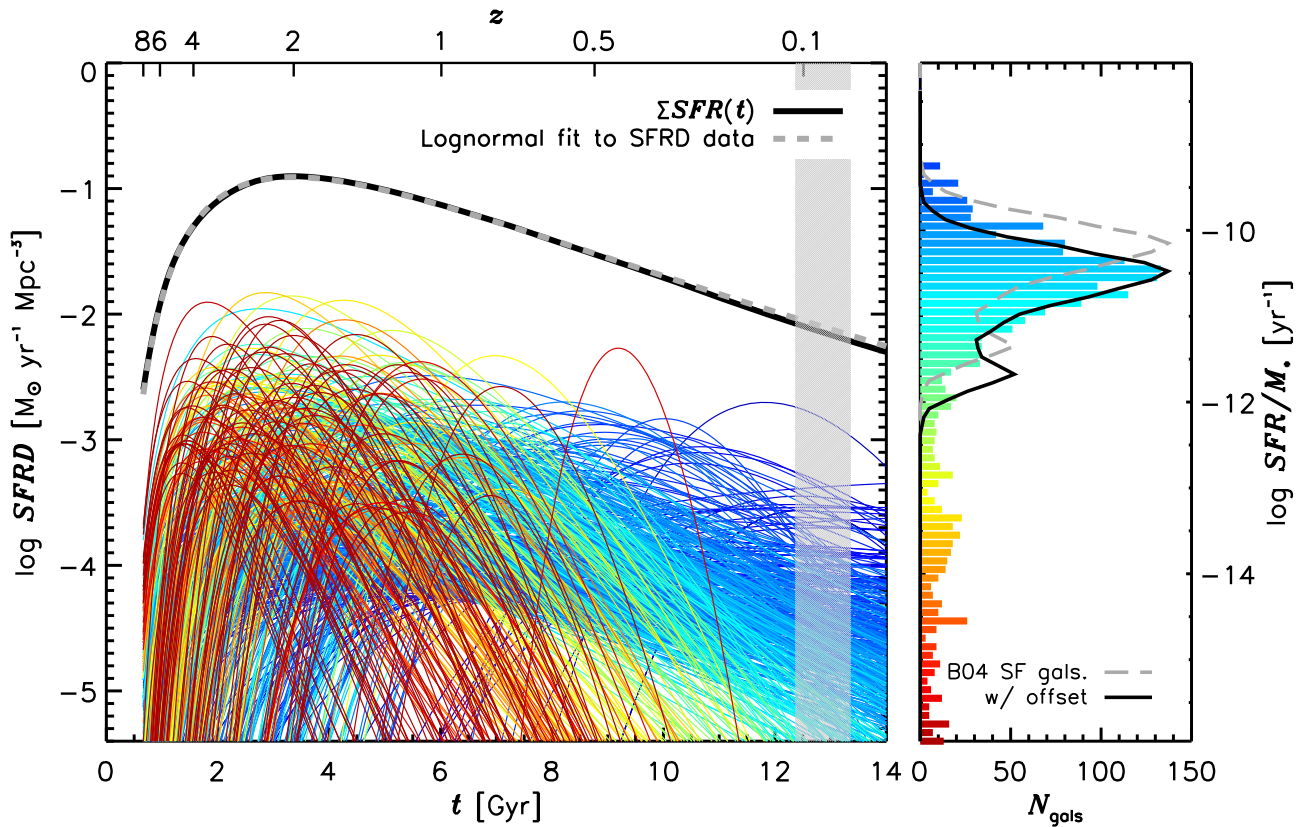


Figure 2. [G13](#) generates one lognormal SFH (colored lines at left; scaled by $10\times$ for visibility) for each galaxy in a volume-complete sample so as to reproduce their measured sSFRs (right; SDSS sSFR distribution overlaid (Brinchmann et al. 2004)), and the cosmic SFRD ($z \leq 8$) (solid black and dashed gray curves, left). Input galaxies span redshifts shown by the shaded vertical band. We assume (1) the functional form of the SFHs, and (2) that non-star-forming galaxies can have any SFR below some detection threshold ($0.05 M_{\odot} \text{ yr}^{-1}$; Figure 9, Equation (4); [G13](#)). Parameterizing the SFHs allows them to be projected to arbitrary lookback times. Any form can be used, but some are better than others (Section 4.3, Appendix A; [G13](#)). The choice of the lognormal was motivated mainly by the shape of the cosmic SFRD (Figure 1). There is an 0.3 dex sSFR calibration offset between the B04 SDSS measurements and those used to constrain the [G13](#) model. This is within systematic uncertainties (Figure 3) and is applied for visualization purposes in relevant comparisons (Figures 10 and 11).

universe’s stellar mass—cannot have evolved via the above prescriptions (see also Dressler et al. 2016).

The second conclusion was that large discontinuities—e.g., starbursts—cannot be invoked to remedy this discrepancy. Such events (briefly) modulate the sSFRs of individual galaxies, but they cannot drive the evolution of that quantity for whole populations over many Gyr (as is necessary): if some objects are in high states, others are in low ones, largely nulling their global effect ([O13](#) Figure 4).

[O13](#) thus demonstrated that many galaxies must have had SFHs that rose and fell rapidly at relatively late times (i.e., faster than the $\text{SFR}-M_*$ relation at $z \lesssim 1$) but were not starbursts.⁷ Taken with the results of, e.g., Rodighiero et al. (2011) and our other investigations (Abramson et al. 2013; Dressler et al. 2013), these findings strongly suggested that abrupt discontinuities—upwards in the form of starbursts, or downwards in the form of rapid quenching—were not the primary shapers of global star formation over the past $\gtrsim 7.5$ Gyr.

Instead, the smooth, long-timescale growth modes must encode important physics. Since [O13](#) showed that their previous τ /delayed- τ descriptions poorly captured this physics, a new SFH form was needed.

[G13](#) provided such a form and universalized the above conclusions into a general description of galaxy evolution.

2.2. Construction: A [Log-]Normal “Model” of Galaxy Evolution

[G13](#) recognized that the universe’s SFH—the evolution of the cosmic SFR density (SFRD; e.g., Madau & Dickinson 2014)—is well described by a lognormal in time (Figure 1). That is:

$$\text{SFRD}(t) = \frac{A_{\text{Uni}}}{\sqrt{2\pi\tau_{\text{Uni}}^2}} \frac{\exp\left[-\frac{(\ln t - T_0, \text{Uni})^2}{2\tau_{\text{Uni}}^2}\right]}{t}, \quad (1)$$

where T_0 , τ are the SFH’s half-mass-time and width (in units of $\ln(\text{time})$), and A is a scaling factor. [G13](#) then asked a simple question: What if galaxy SFHs shared this functional form?

To explore this question, [G13](#) took 2094 local galaxies with $\log M_* \geq 10$ from the Sloan Digital Sky Survey (SDSS; York et al. 2000) and Padova Millennium Galaxy and Group Catalog (PM2GC; Calvi et al. 2011) and assigned them lognormal SFHs by fitting for the 2094 (T_0, τ) pairs⁸ that best reproduced (1) each galaxy’s observed (M_*, SFR) while (2) ensuring that

⁷ The existence of such SFHs implies that star-forming galaxies of the same M_* do not grow up together/evolve homogeneously along the $\text{SFR}-M_*$ relation. We argue that the latter scenario is both overly constraining and inaccurate (Sections 3 and 4; Dressler et al. 2016).

⁸ For each galaxy, A is set such that $\int_{t_0}^{\text{obs}} \text{SFR}(t) dt = M_*^{\text{obs}}/f$, where f is a $\text{SFR} \mapsto M_*$ conversion factor (0.7 for a Salpeter IMF).

the ensemble of SFHs summed to the SFRD as far back as the data allowed ($z \lesssim 8$; see Figure 2).

G13 and Appendix A present technical details, but the above captures the model’s essence: a continuum of smooth SFHs rising and falling naturally over loosely constrained timescales and peaking at loosely constrained times. No physical prescriptions govern these behaviors, only the universe’s observed SFH and the end-states of a set of real objects.

In this sense, the G13 “model” is not a model at all, but a realization of possible lognormal SFHs for a sample of galaxies. As such, all physics must be inferred from the resulting (T_0, τ) distribution, or read into the form of the SFH. Sections 3.2, 3.3, and 4 revisit this point.

Note: Given its $\log M_*(z \approx 0) \geq 10$ mass limit, G13 does not describe the SFHs of local dwarfs (though neither does it contain any). Unlike giant systems (Rodighiero et al. 2011), these are known to grow via stochastic bursts (Weisz et al. 2014) which a single lognormal cannot capture. Hence, the following discussion is likely not applicable to them. Our principal interest lies in understanding what shapes the longer-term growth modes of galaxies that have consistently dominated the universe’s mass and SFR budget (Section 2.1).

2.3. Previous Descriptive Successes

In G13, we showed that the model naturally reproduced the population of galaxies O13 (and subsequently Dressler et al. 2016) found to have both late-peaking and narrow SFHs, something impossible to do using τ or delayed- τ SFH parameterizations (Figures 2 and 9; O13). This encouraged us that the G13 model represented not only an improved description, but perhaps a physically meaningful worldview.

More encouragement has come as the model continues to reproduce a range of observations it was neither constrained by nor intended to describe. We explore most of these results below, but review here what we have demonstrated to date:

1. Constrained only by the cosmic SFRD and $z \approx 0$ input galaxies, G13 (Figure 7) showed that the model faithfully reproduced sSFR distributions at four epochs at $z \lesssim 1$. As O13 did for τ /delayed- τ models, G13 showed that Gaussian SFHs failed this test badly.
2. Turning the $z \lesssim 1$ sSFR distributions into modeling constraints and refitting for each galaxy’s (T_0, τ) , G13 (Figure 13) showed that the model reproduced the $z \sim 2$ SFR- M_* relation (the “Star Formation Main Sequence,” or “SFMS”) from Daddi et al. (2007).⁹
3. Similarly constrained at $z \lesssim 1$, Abramson et al. (2015) showed that the model reproduced the evolution of the stellar mass functions of star-forming and non-star-forming galaxies at $z \lesssim 2.5$. This was despite the fact that no mass-sensitive constraints were ever imposed except via the intrinsic mass distribution of the $z \approx 0$ input data.

These achievements by no means make the G13 model uniquely successful: others are equally capable (e.g., Sparre et al. 2015). However, as G13 entails no explicit treatment of any physical process, we are spurred to investigate what makes

⁹ The agreement with $z \lesssim 1$ sSFR distributions achieved by the initial model suggests that the $z \sim 2$ SFMS could have been reproduced without turning these into constraints. We have not performed this test, however.

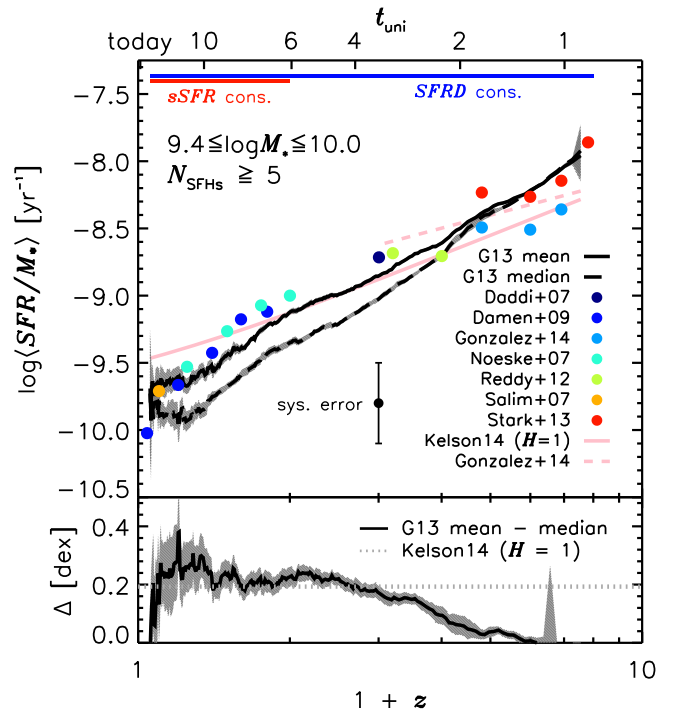


Figure 3. Top: mean/median sSFR ($9.4 \leq \log M_* \leq 10.0$) from G13 (solid/dashed black lines), describing systems on “the flat part” of the SFMS where sSFR is mass-independent. Model measurements span epochs where at least 5 SFHs support it. Data (colored points; estimates of the *mean*) and systematic uncertainties (black point) are from González et al. (2014) and Salim et al. (2007). Agreement is good everywhere, especially considering that the model is unconstrained at $z > 1$ except by the cosmic SFRD (blue bar at top), and that data from Stark et al. (2013) and González et al. (2014) were published after the model was created. Bottom: the mean–median offset, a probe of the SFMS scatter. To $z \sim 2$, the model reproduces the ~ 0.2 dex offset implied by K14’s preferred stochastic evolutionary scenario (see text; his Figure 3(a)), showing that paradigms with opposite assumptions can produce the same (non-) evolution in this quantity.

this simple mathematical description an accurate expression of complex, poorly understood underlying physics.

2.4. New Challenges

Here, we test the G13 model in a variety of new contexts to better ascertain its standing as a serious paradigm of galaxy evolution. We contend that its successes demonstrate its viability. But, at a minimum, they highlight issues at the descriptive level in this field that deserve real attention, or call for a reevaluation of some of its cornerstone observations.

2.5. A Note to the Reader

Some vocabulary used below is unconventional in astronomy. Two definitions are in order:

Longitudinal data—repeated observations of individual objects—and only those objects—over a period of time. G13 describes such data (Section 2.2). Unfortunately, given the timescales of galaxy evolution, astronomers have no access to this information. Instead, we approximate it by time-ordering *cross-sectional* data.

Cross-sectional data—observations of a sample of objects at a single epoch. We have no recourse but to assemble series of these data for different objects at different times to constrain galaxy evolution, so we must also use them to

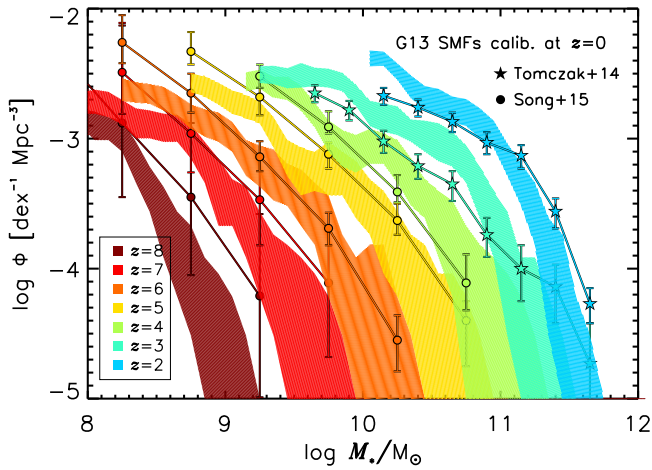


Figure 4. G13 stellar mass functions (colored bands; 90% credibility interval) at $z \leq 8$ compared to data from Song et al. (2016) and Tomczak et al. (2014) (points). The model is calibrated only to $z \approx 0$ data (Moustakas et al. 2013) over the range $\log M_* \in [10, 10.5]$ (Figure 18), and constrained only by the cosmic SFRD over all epochs plotted.

evaluate any and all models (Section 3.1). Sections 3.3 and Appendix B discuss implications of this fact.

3. RESULTS

We examine the G13 model from three different angles:

1. Section 3.1 deals with its ability to reproduce important metrics in galaxy evolution neither used in its construction nor discussed in our previous work.
2. Section 3.2 explores how its core premises/traits guide these behaviors.
3. Section 3.3 projects the model into new spaces and uses the results to build a physical interpretation of both G13’s behavior and its axiomatic underpinnings. It also presents some concrete predictions that may be testable with (future) observations to further assess the viability of the paradigm.

We split the analysis in this way to clarify which aspects of the results reflect “superficial” details of the model’s current realization and which are expressions of its underlying structure—its “DNA.” We assume $(H_0, \Omega_m, \Omega_\Lambda) = (73 \text{ km s}^{-1} \text{ Mpc}^{-1}, 0.27, 0.73)$ and a Salpeter (1955) initial mass function (IMF).

3.1. Comparing G13 to Data

3.1.1. The $\langle \text{sSFR} \rangle$ of Low-mass Galaxies Since $z = 7$

Large photometric surveys have provided a wealth of data on the mean sSFR of relatively low-mass galaxies— $9 \lesssim \log M_* \lesssim 10$ —at $z \lesssim 7$ (e.g., Daddi et al. 2007; Noeske et al. 2007; Damen et al. 2009; Reddy & Steidel 2009; Stark et al. 2013; González et al. 2014). These objects reside on the “flat part” of the sSFR– M_* relation, where $d \log \text{sSFR} / d \log M_* \simeq 0$ (or $d \log \text{SFR} / d \log M_* \simeq 1$; e.g., Whitaker et al. 2014). $\langle \text{sSFR} \rangle$ is thus akin to the SFMS zeropoint.

In this mass regime, the only G13 modeling constraint was the cosmic SFRD ($\dot{\rho}$); i.e., the evolution of the sum of SFRs of all galaxies. Additionally, our understanding of $\langle \text{sSFR}(z) \rangle$ has

evolved markedly since the construction of the model (Stark et al. 2013; González et al. 2014; Kelson 2014). For both reasons, these data provide a meaningful test of the accuracy of the G13 SFHs and therefore the validity of the approach.

Figure 3 presents this test. Here, we show how the $\langle \text{sSFR} \rangle$ of SFHs with $9.4 \leq \log M_*(z) \leq 10$ compares to data for galaxies in the same mass range taken from González et al. (2014) and Salim et al. (2007). Agreement is remarkably good at all $z \leq 7$ (where at least 5 SFHs support the model measurement), especially considering systematic uncertainties (black point), and that the points from Stark et al. (2013) and González et al. (2014) were published after the model was created. In fact, at the time of its construction, $\langle \text{sSFR} \rangle$ was thought to flatten at $z > 2$ (Stark et al. 2009); only later, when SED fits were adjusted to account for emission line fluxes (Stark et al. 2013), did those points show the monotonic increase anticipated by G13.

We neglect mergers here and in all following Sections. Results from Leitner (2012), Behroozi et al. (2013b), Abramson et al. (2015), and our own numerical tests suggest that this move does not affect the arguments below, but it is worth keeping in mind. For example, as presented (see Appendix A), the SFHs may average over the histories of multiple pre-merged systems at high- z . Assuming such pieces have similar sSFRs—as G13 did and is implied by the SFMS slope (see above and Section 3.1.3)—model galaxy counts would be impacted more than SFR-related metrics (Section 3.1.2).

Figure 3 is meant to be compared with Figure 3(a) of (Kelson 2014, hereafter K14). That model predicts a similar trend to G13 (pink line) but adopts a very different physical premise, positing that individual galaxies evolve via quasi-stochastic SFR discontinuities correlated on arbitrary time-scales denoted by H . $H = 1$ corresponds to *Hubble timescale* correlations and fits the data well (K14 Section 3). It implies a *median* $\text{sSFR}(t) \approx (H + 1)/t = 2/t$, with a mean ~ 0.2 dex higher. That offset is a measure of the SFMS scatter, σ_{SFMS} .

Notably, Figure 3, bottom, shows that, to $z = 2$ —three Gyr prior to the earliest sSFR constraint—the G13 model *also* produces a ~ 0.2 dex mean/median offset. This is intriguing because K14 and G13 are based on not only different, but *opposite* assumptions: G13 assumes that SFHs are smooth, parametrizable, and diversified quasi-deterministically; K14 assumes that they are discontinuous, unparametrizable, and diversified by stochastic changes in equilibrium conditions. Section 4.2 revisits the potential ramifications of such comparably accurate yet physically distinct descriptions.

Moving to higher- z , the G13 mean and median sSFRs converge, implying a decrease in σ_{SFMS} . This result is inconsistent with recent analyses that show it to be roughly constant instead of shrinking (e.g., Schreiber et al. 2015; Speagle et al. 2014). Yet deciphering this discrepancy—and some others that will arise later—is non-trivial.

First, it may be superficial, not structural, at least over the timespan probed. The model offset is sensitive to the location of sSFR distributional constraints, and so could be changed by imposing additional ones at $z > 1$.

Second, it may be accurate: recall that the data reflect cross-sectional samples (different objects at different times; Section 2.5), which, at $z > 0$, will almost certainly contain objects that do not represent progenitors of the G13 input sample (Sections 2.2) and so are not described by the model. Hence, it may be that, when true progenitors are isolated a priori, their sSFRs do converge at early times, with non-

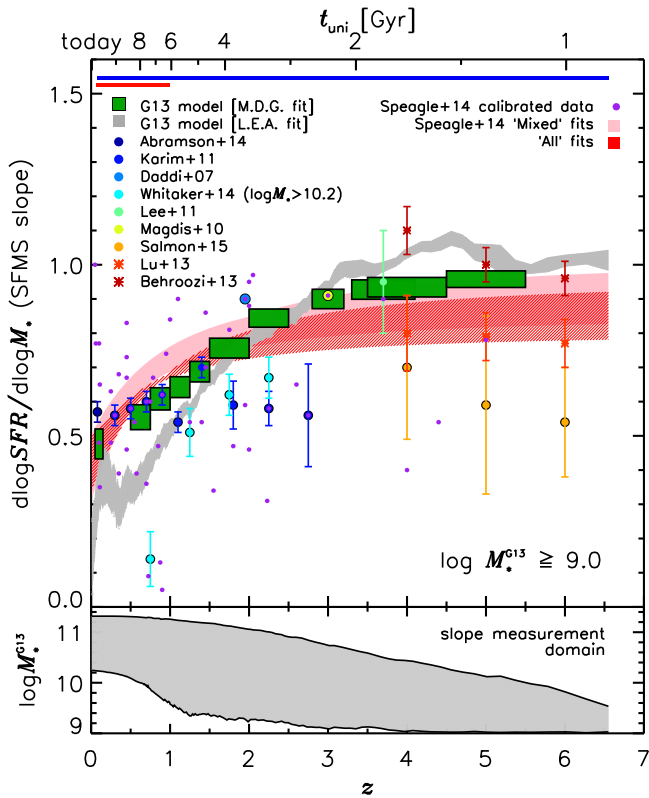


Figure 5. Top: **G13** predictions for the evolution of the slope of $SFR(M_*)$ (green boxes, gray shaded band) compared to measurements from the literature (colored circles) and semi-analytic models (stars). Systematic uncertainties are large, but a general trend toward more-negative values at later times is apparent in both the data and **G13** (see Speagle et al. 2014’s extensive analysis; pink band). Bottom: the 95% **G13** $M_*(t)$ range considered by LEA, who continuously fit the SFHs as they evolved. MDG bootstrapped them in redshift windows similar to those in the data and used a different σ -clipping scheme. The fact that there is a sensible model trend at all is unexpected.

ancestral galaxies “padding-out” the scatter in the data (Figure 11).

That said, at some point a decrease in σ_{SFMS} (but not necessarily that shown in Figure 3) must become structural: At $t \ll \langle \tau \rangle \sim 1.6$ Gyr—the mean **G13** characteristic timescale for the ensemble of SFHs—**G13** “galaxies” share approximately the same exponential $SFR(t)$, leading to a unique SFH for each $M_*(t)$ (which becomes a scale factor). So, if σ_{SFMS} is ultimately shown *not* to decrease at those times ($z \gtrsim 3$ –4)—or if the addition of sSFR constraints at $z > 2$ cannot remedy any real discrepancies there—we will learn something important about the limitations of smooth (lognormal) SFHs.

3.1.2. The Evolution of the Stellar Mass Function Since $z = 8$

Besides the evolution of the mean sSFR of (low-mass) galaxies, another critical observation to reproduce is the evolution of the stellar mass function. Abramson et al. (2015) demonstrated the success of the **G13** model in this context at relatively low redshifts ($z \lesssim 2.5$). We now extend this test to earlier epochs.

Figure 4 plots the **G13** prediction for the evolution of the total stellar mass function at $2 < z \lesssim 8$, as far back as data are available. The only calibration applied to the model is an absolute y-offset to match the $z \approx 0$ stellar mass function of Moustakas et al. (2013) at $10 \leq \log M_* \leq 10.5$ (10 Gyr away from the earliest epoch plotted), and the removal of some over-

represented high-mass galaxies in the **G13** input sample (Appendix D, Figure 18).

Considering that no mass functions were used to constrain the model (except through the intrinsic mass distribution of the $z \approx 0$ input sample), agreement is remarkably good over more than a decade in M_* , two decades in space density, and lookback times approaching 13 Gyr. This finding substantially extends those of Abramson et al. (2015) regarding the accuracy of **G13** SFHs at the mass-bin level.

A systematic offset may exist at $z \sim 8$, but we caution against over-interpreting it: At all epochs shown, the model is constrained only by the sum of all SFHs. Hence, there is no obvious reason to expect the shape or normalization of the stellar mass functions to look anything like the data (something we have confirmed through numerical tests). Additionally, we have not verified that the highest- z measurements—convolved with the appropriate $SFR-M_*$ relation—are in fact consistent with the SFRD used to derive the model, which only runs to $z = 8$. Finally, assuming the data are reliable and consistent, and that the offset is meaningful, this tension is again probably superficial; it should be eased by introducing further constraints at $z > 1$.¹⁰ Indeed, a comparison of **G13** Figures 5 and 9 shows that adding sSFR constraints tends to push (subsets of) SFHs toward earlier peak-times, precisely as the data suggest.

For all of these reasons, we view this result as an unanticipated success of the **G13** model.

At low-masses, the **G13** stellar mass functions are truncated where they turn over (Appendix D, Figure 18). Because of the longitudinal nature of our sample (it tracks the same group of objects through time; Section 2.5), this completeness threshold evolves *downward* with increasing redshift. At high-masses, a cut-off is imposed by the location of the most-massive model “galaxy” at any epoch. Increasing the number of input galaxies to which SFHs are assigned will modulate both bounds, as will incorporating mergers, which remove high-mass and add low-mass systems. Figure 18 provides more context for how these effects might enhance or degrade agreement with data.

A final note: The evolution of the SFMS—and therefore $\langle sSFR(z) \rangle$ —is linked to that of the cosmic SFRD via the galaxy stellar mass function (at least for star-forming systems; Abramson et al. 2015). Since the **G13** model is designed to match the SFRD, and, as shown in Section 3.1.1, also captures the evolution of $\langle sSFR(z) \rangle$ over most of the interval spanned in Figure 4, a concern might be that it is guaranteed to capture the evolution of the stellar mass function. This is only true, however, over the mass range $9.4 \lesssim \log M_* \lesssim 10$ at epochs where such galaxies dominate the SFRD. There is nothing in the modeling procedure that guarantees matching the rest of the stellar mass function at those epochs, or any of it at other epochs. That we accomplish the latter suggests that the model must correctly reproduce the evolution of the entire SFMS. We explore this next.

3.1.3. The Evolution of the SFMS Slope Since $z = 6$

Figure 3 tracked galaxies with $\log M_*(z) < 10$, where effectively all stellar mass is associated with star formation

¹⁰ The volume probed by, e.g., Tomczak et al. (2014) is $\sim 10\times$ that of PM2GC (used by **G13**). Our results are thus subject to shot-noise in terms of the area of (T_0, τ) space the data *must* span: The input sample may contain too few massive (red) galaxies to guarantee very-early-peaking output SFHs. Rerunning the model on a larger volume might therefore also remedy this issue, but is beyond the scope of this work.

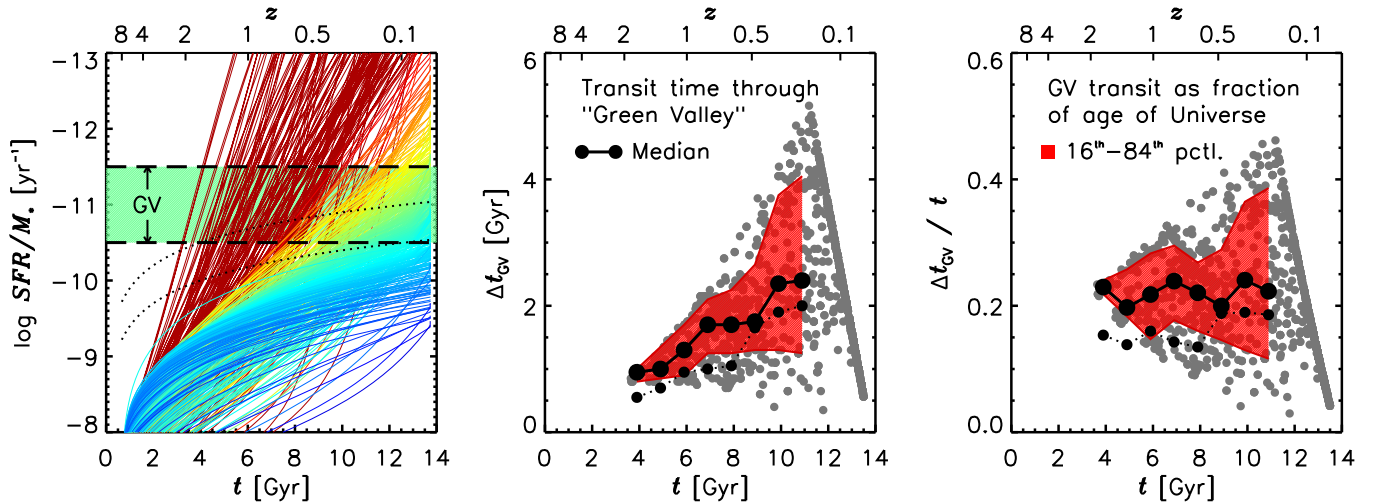


Figure 6. G13 naturally describes the shift from so-called “fast-” to “slow-track” quenching. Left: model sSFRs as a function of time colored by present-day values (red = low, blue = high; y-axis oriented as in Barro et al. (2013) and the color/mass diagram). A fiducial transition zone between SF and quiescent galaxies—the “Green Valley” (GV)—is highlighted by the green shaded band with an alternative—evolving $\propto 1/t$ with half the width in $\log SFR$ —plotted as dotted lines. Slopes of SFHs crossing the GV are steep at high- z but shallower today, implying quenching timescales were rapid in the past and are slower now. The middle panel plots GV transit time— Δt_{GV} —against cosmic time to quantify this trend; small dots show medians using the $1/t$ GV definition. Interestingly, while G13 recovers fast- and slow-track quenching, once normalized by the age of the universe at crossing (right), the absolute trend disappears: galaxies spend $\sim 20\%$ of their life in the GV. This jibes with the scale-free interpretation of K14 and also supports the idea that non-star-forming galaxies have compressed SFHs, but have *not* been subject to “special” physics as a class. (See also Figure 11 of Vulcani et al. 2015a.)

(e.g., in a disk) and $d \log sSFR / d \log M_* \approx 0$. These define the SFMS’s zeropoint. Now we examine the evolution of its slope.

To some extent, the SFMS is affected by the presence of bulges in high-mass galaxies, which add M_* at fixed SFR. This effect would tend to push the SFMS slope to more-negative values (Abramson et al. 2014; Whitaker et al. 2015; A. Oemler et al. 2016, in preparation, Figure 13, below; but cf. Schreiber et al. 2016, who attribute the effect to reduced SF efficiency). As it is now believed that bulge growth could be a multi-phase process, with some galaxies forming them at late times (e.g., Lang et al. 2014; Huertas-Company et al. 2015), the slope of $SFR(M_*)$ should monotonically evolve from ~ 1 at high- z to lower values today if fit by a single power-law.

Figure 5 shows data (points), drawn mainly from Table 4 of Speagle et al. (2014). As substantiated by those authors’ thorough meta-analysis (pink/red bands), the slope indeed evolves subtly from $d \log SFR / d \log M_* \sim 0.8$ at $z \approx 4$ to ~ 0.4 today.

As for G13’s predictions, independent fits based on (1) continuous measurements of the SFHs (gray band; performed by LEA) or (2) bootstrapping in redshift windows defined by the data (green boxes; performed by MDG) show the slope to evolve similarly, dropping from near unity at $z \approx 4$ also to ~ 0.4 today.¹¹ Beyond $z \approx 4$, the G13 model is broadly consistent with both the Speagle et al. (2014) projections, and results from at least one semi-analytic model based on dark matter halo growth (Behroozi et al. 2013b; perhaps less so Lu et al. 2014; both plotted as colored stars).

Clearly, as evident from comparing points at similar epochs, systematic uncertainties in the data are large at all z due, e.g., to differences in mass ranges and definitions of “star-forming.” These remain despite attempts to calibrate-out variations in

measurement techniques (purple points; see Table 6 of Speagle et al. 2014). Also, the (scant) extant measurements at $z > 4$ (Steinhardt et al. 2014; Salmon et al. 2015) are generally lower than model predictions, so more and better information is needed to fully assess any findings here.

Nevertheless, a core point is robust: The G13 model neither contains physical prescriptions (for bulge growth or star formation efficiency) nor is bound by constraints that would inform its behavior in Figure 5. In fact, at $z > 1$, it “knows” only about the cosmic SFRD. As the mass-insensitive integral of the SFRs of all galaxies, this is perhaps as far removed from the SFMS slope—sensitive to mass differentials between individual systems—as a constraint can be. Hence, as above, that G13 qualitatively reflects the data—and indeed might be quantitatively consistent with state-of-the-art analyses and semi-analytic models—either speaks to its validity, or comments on the physical information content of the metrics examined. Sections 3.3 and 4.3 revisit this point.

3.1.4. The Congruity of Fast- and Slow-track Quenching

So far, we have dealt with star-forming galaxies. These dominate the stellar mass density of the universe at $z \gtrsim 1$ (Muzzin et al. 2013; Tomczak et al. 2014), so limiting the discussion to them is not too biasing. However, since we have set the G13 model in opposition to those specifically designed to create *non*-star-forming galaxies, it behooves us to discuss quenching in this framework. We do so now and in the next section.

Recently, a picture has emerged in which galaxies transitioned from a star-forming to non-star-forming state more rapidly in the past than they do today (on average). This dichotomy between “fast-track” and “slow-track” quenching (Barro et al. 2013; Schawinski et al. 2014) has been used to suggest an interesting change in the mechanisms inducing this transformation over cosmic time. The quenching timescale may therefore be a critical datum in galaxy evolution.

¹¹ Both approaches limited fits to SFHs with $\log M_*(z) \geq 9$ —a lower bound encompassing almost all data sets. Discrepancies here illustrate the difficulty of measuring this quantity robustly.

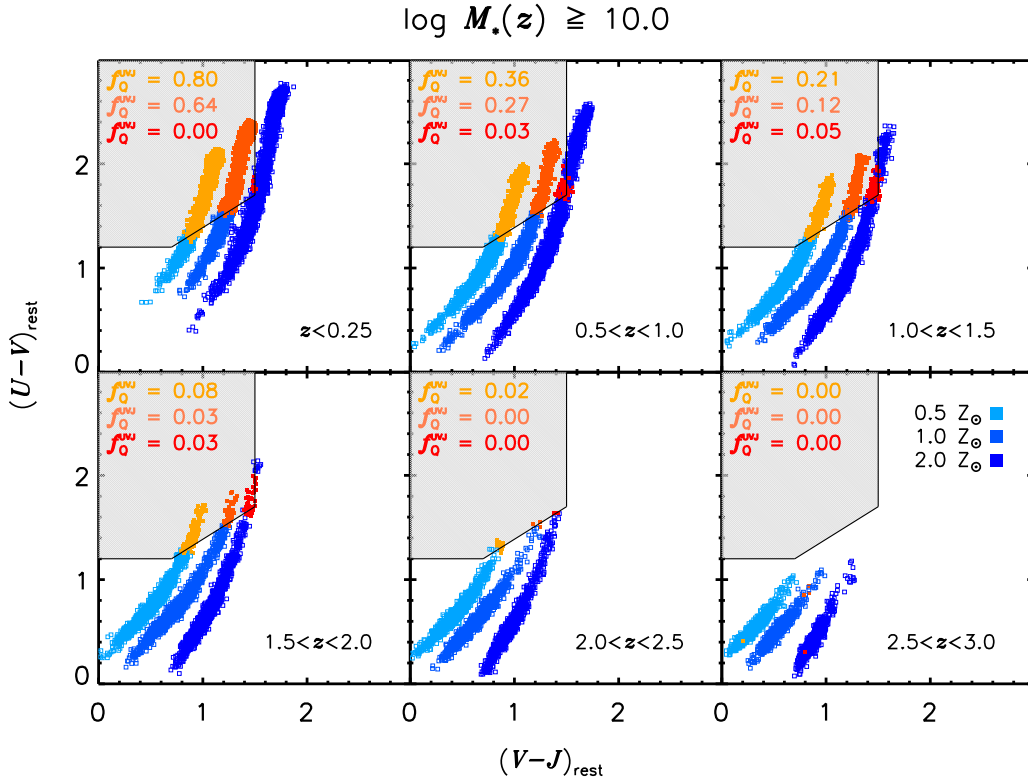


Figure 7. Synthetic color–color diagrams derived from the G13 SFHs and the models of Bruzual & Charlot (2003) shown as in Tomczak et al. (2014). Blue/orange points show SFHs that would be classified as star-forming/quiescent, respectively, using the UVJ criterion of that work. Each SFH was run assuming a fixed metallicity of $Z = \{0.5, 1.0, 2.0\} \times Z_\odot$ (light, medium, and dark shades, respectively). Only SFHs with $\log M_*(t) \geq 10$ are shown to correspond with the mass range probed by Abramson et al. (2015). As printed in the top-left of each panel, at face value, the model underproduces quiescent systems at $z \gtrsim 1.5$; integrating the Tomczak et al. (2014) mass functions over this mass range yields $f_Q^{\text{obs}} \sim 24\%–28\%$ at these epochs. Yet, the model matches data if an SFMS-relative sSFR cut is used (Abramson et al. 2015), implying that this shortcoming could be ameliorated by introducing additional sSFR constraints.

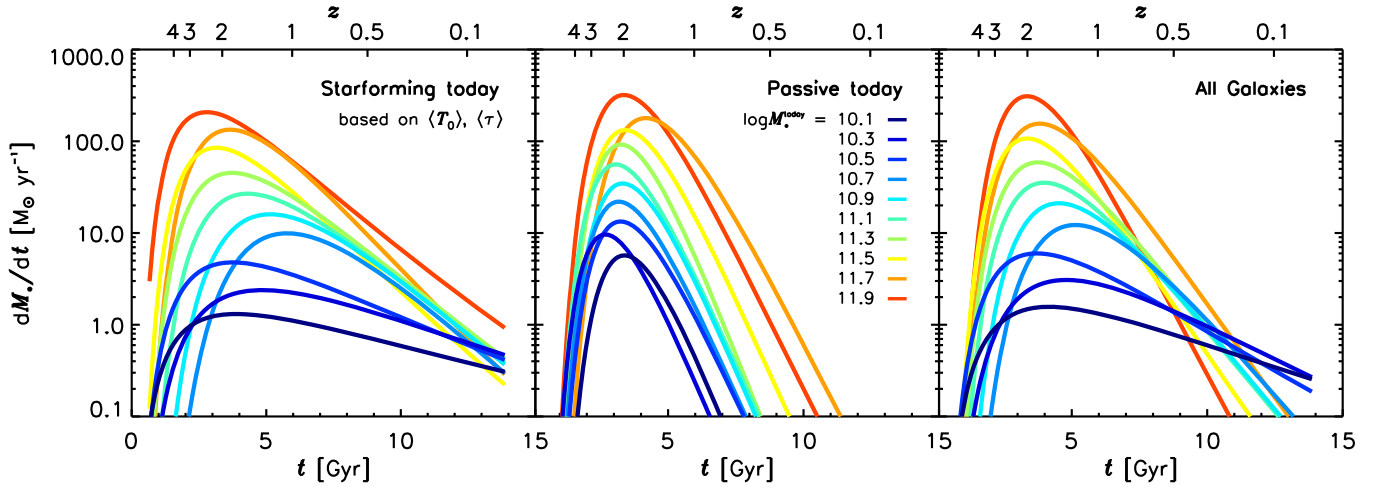


Figure 8. Lognormal SFHs derived from the median (T_0, τ) parameters in bins of present-day M_* . Panels show results considering all galaxies (right), or only those currently classified as star-forming (left) or non-star-forming (middle; cut at $\text{sSFR} = 10^{-11} \text{ yr}^{-1}$). A “downsizing” trend is clear in the star-forming population—driving that seen for all galaxies—with today’s lower-mass systems forming later/with more extended SFHs than higher-mass systems. That trend is not visible in present-day passive G13 galaxies because these have effectively no sSFR constraint at late times. Hence, the SFHs are guided mainly by the cosmic SFRD, and so follow its shape closely regardless of final mass (note the consistent peak at $z \gtrsim 2$). In reality, some of these systems would have formed earlier than this. Such behavior leads to the under-abundance of objects in the passive region of the $z \gtrsim 1.5$ UVJ diagram (Figure 7).

The G13 model suggests otherwise, however. Figure 6, left, shows tracks of $\text{sSFR}(t)$ for a representative subset of SFHs. As in Barro et al. (2013), the y-axis is oriented to correspond with the color-mass diagrams in which the active/passive transition is often discussed (red/low-sSFR up). The green horizontal swath and dotted tracks in sSFR approximate two definitions of

the transition regime (see below); i.e., the “Green Valley” (GV). The middle panel shows the time each SFH spends in the GV; i.e., the quenching timescale, Δt_{GV} .

Two things are clear. First, the G13-derived quenching timescale is indeed an increasing function of time, with transformations being faster in the past than they are today

(middle panel). Crossing times depend quantitatively on the definition of the GV, but—as shown by choosing one that evolves $\propto 1/t$ (small dots)—the trend will not (for reasonable choices). Thus, **G13** naturally reproduces the fast-track/slow-track quenching transition. However, this *need not correspond to a change in quenching physics* because the model contains no explicit quenching mechanism.

Instead, Figure 6, left, makes it clear that the trend is purely a function of the slope of the SFHs. Upon reflection, this result is inevitable: To reach their final mass early—i.e., to become a passive system at high- z —galaxies *must grow rapidly*. Hence, their SFHs *must* rise and fall steeply, and Δt_{GV} must be short. Conversely, a galaxy crossing the GV today is likely to have had a much more extended SFH, and therefore a much more slowly evolving $s\text{SFR}(t)$. Thus, it will naturally linger in the GV with larger Δt_{GV} relative to systems that quenched earlier (see also Pacifici et al. 2013).

This phenomenon reflects a deeper point: different parts of (T_0, τ) space support the SFRD at different times (Appendix C). Hence, the SFHs corresponding to red or blue galaxies at *fixed* (M_*, t) can derive from disjoint loci in **G13**’s underlying unimodal (T_0, τ) distribution: the former hail from regions dominating the past SFRD, the latter from regions dominating the present (Figures 15 and 17). Thus, continuous descriptions can manifest quenched/star-forming bimodalities.

This interpretation—that the transition from fast- to slow-track quenching does not represent a change in mechanism, but the natural shift in the GV population from galaxies that are rapidly growing to those that are slowly growing—is driven home by Figure 6, right. There, we replot Δt_{GV} , but as a fraction of the age of the universe mid-way through the crossing ($\Delta t_{\text{GV}}/t$). Seen in this light, there is no evolution in the crossing time at all. Scatter between SFHs is large, but the median stays at roughly 20% of the age of the universe (effectively the age of a galaxy) forever. (Zolotov et al. 2015 find a similar 35% using hydrodynamic simulations.) As such, quenching appears to be a scale-free phenomenon qualitatively consistent with the picture of K14 (and perhaps Stringer et al. 2014). If so, quenching timescales may yield little insight into evolutionary processes.

Barro et al. (2013) tied fast- versus slow-track quenching to galaxy sizes, with compact galaxies evolving via the fast route. As we show in Sections 3.3.1 and 3.3.3, the **G13** model naturally establishes this connection while maintaining an underlying scale-free nature to quenching through τ .

3.1.5. The UVJ Diagram and Color Selection

With the advent of large, high-redshift surveys, the division between star-forming and non-star-forming/quiescent/passive galaxies must be made according to photometric and not spectroscopic criteria. A popular distinction is based on the distribution of galaxies in $U - V/V - J$ color-color space—the “UVJ diagram” (e.g., Williams et al. 2009). Hence, another useful test of the **G13** model is to examine its predictions for the evolution of galaxies in this plane.

Figure 7 shows the results of running the stellar population synthesis (SPS) code of Bruzual & Charlot (2003) to extract synthetic colors for each SFH in different redshift intervals assuming three different metallicities— $Z \in \{0.5, 1.0, 2.0\} \times Z_\odot$. We limit the plotting to galaxies with $\log M_*(z) \geq 10$ at any epoch to remain consistent with Abramson et al. (2015).

Taking the results at face value, the **G13** model fares worse in this context than in the others explored so far. While, accounting for metallicity spreads, our results appear roughly consistent with, e.g., the Tomczak et al. (2014) data at $z \lesssim 1.5$ ($0.28 \lesssim f_Q \lesssim 0.47$), the model underproduces photometrically quiescent systems at higher- z , returning percentages in the single-digits as compared to the 24%–28% from that work. As such, had we split the galaxy population using this technique as opposed to an evolving SFMS-relative $s\text{SFR}$ cut, we would not have correctly reproduced the evolution of the quiescent mass function shown in Abramson et al. (2015).

However, because the model *does* reproduce the correct stellar mass function evolution using an evolving $s\text{SFR}$ cut, the lack of *red* objects would seem to reflect a mismatch in the mean shapes of SFHs leading to color-defined passive galaxies at those epochs, not a real deficit of “quenched” systems. That is, the comparatively low- $s\text{SFR}$ systems we identified were “quenching,” just not rapidly enough to produce colors photometrically identifying them as such at early epochs.

Thus, Figure 7 may again reveal a superficial issue related to **G13**’s data constraints: if a $z > 1$ $s\text{SFR}$ constraint was introduced—forcing some SFHs to lower (T_0, τ) and thus steeper declines—a model with identical assumptions may well produce enough systems reaching low enough $s\text{SFR}$ s at early enough times to match the photometry.

The next section elaborates on this statement to get at *why* the model deviates from the data in this context.

3.1.6. Average Galaxy SFHs and Downsizing

First insights come from Figure 8, showing the average SFHs of galaxies split by present-day M_* and SFR.

When examining either all SFHs (right), or the subset that lead to “star-forming” systems today (left; $\log s\text{SFR} \geq -11$), a clear trend emerges for more-massive systems (redder curves) to grow/peak earlier than less massive systems (bluer curves). This model trend corresponds with the now well-established observational paradigm of galaxy “downsizing” (e.g., Cowie et al. 1996; Neistein et al. 2006; Pacifici et al. 2013), and is directly related to the apparent slow-down in GV crossing times discussed in Section 3.1.4.

The middle panel, however, reveals no such trend for **G13** SFHs leading to today’s non-star-forming/passive galaxies [$s\text{SFR}(z=0) < 10^{-11} \text{ yr}^{-1}$]. These histories are basically identical when scaled by their present-day M_* .

The reason the model produces this uniformity in passive galaxy SFHs has to do with the amount of $s\text{SFR}$ “real estate” they can occupy compared to their star-forming counterparts. As Figure 2 shows, star-forming galaxies today, by definition, must land on a relatively narrow $s\text{SFR}$ target (~ 1 dex) defined by the $z \approx 0$ SFMS. Passive galaxies, on the other hand, can occupy an effectively infinite space below this locus. Thus, they have essentially no late-time $s\text{SFR}$ constraint. As the only other constraint in the model is the cosmic SFRD, the SFHs for these systems naturally tend toward that history, rising quickly to peak at $z \sim 2$ and then falling as seen in Figure 8. Since colors cannot discriminate between objects that last formed stars, e.g., 3 versus 10 Gyr ago—but are very good at identifying those that formed them within the past 0.1–1 Gyr—the ability for these SFHs to take arbitrarily low $z \approx 0$ $s\text{SFR}$ s need not translate to red colors at $z \gtrsim 2$, allowing the corresponding area in UVJ-space to remain vacant.

The upshot is that, absent additional sSFR constraints, it is difficult for the G13 approach to produce large numbers of SFHs reaching low sSFRs before the cosmic SFRD peaks (Figure 2). This leads to the dearth of red *UVJ* galaxies seen at $z > 1.5$ (Figure 7) and probably contributes to any real underestimation of the $z \approx 8$ stellar mass function (Figure 4). As such, this behavior will again be modulated by adding sSFR constraints at higher- z , or by rerunning the fitting on a larger number of systems that presumably more-fully span (T_0, τ) space.

A final note about downsizing: As is empirically well-established and shown in Figure 3, the SFMS has fallen consistently over cosmic time. This means that galaxies doubled in mass (e.g.) faster in the past. As is also well-established and shown in Figure 5, the slope of the (high-mass end of the) SFMS is sub-unity for most of cosmic time. This means that low-mass galaxies always double in mass faster than high-mass galaxies. These facts combine to ensure that galaxies growing-up at earlier times *rushed to states of comparative inefficiency* (“raced to the bottom”) more quickly than galaxies growing-up at later times. This is downsizing, and it is a generic prediction of any model leading to/ incorporating $d \log \text{SFR} / d \log M_* < 1$. Why the (high-mass) SFMS slope is < 1 is an interesting question, but it is unclear if the details of downsizing will be edifying in that context.

We discuss other such aspects related to the structure/“genetics” of the G13 model in the next section.

3.2. Exploring G13’s Structure

Section 3.1 dealt with G13’s appearance in frames defined by a range of cross-sectional data. This “phenotype”—the outward manifestations of a suite of smooth, continuous, lognormal SFHs—compares favorably to most the observations studied, but because the model lacks explicit physical prescriptions, interpreting this fact is non-trivial.

To derive a physical interpretation, we now turn to G13’s “DNA”—the structural properties/axiomatic underpinnings/intrinsic features that lead to the above results. Our confidence in the model depends more on the accuracy of ideas in this Section than tests in the last.

3.2.1. The G13 Interpretation of the SFMS Scatter

The most salient structural aspect of G13 is its interpretation of the SFMS. As illustrated by Figure 9, it reads σ_{SFMS} —the scatter in (s)SFRs at fixed M_* —primarily as a gradient in the lognormal half-mass-time parameter, T_0 . That is, *the scatter in the SFMS reflects an age gradient in this model*.

Indeed, the model suggests that even the “small” observed dispersion of ~ 0.3 dex can accommodate a ~ 7 Gyr spread in half-mass times. Intriguingly, using a totally different approach, Speagle et al. (2014) derive a similar interpretation for σ_{SFMS} , and an almost identical ~ 6 Gyr (2σ) spread. As with K14, however (Section 3.1.1), the models differ meaningfully in how σ_{SFMS} is generated/maintained (see below).

We stress that no aspect of the modeling procedure explicitly induces such gradients. Instead this behavior is emergent; the addition of sSFR constraints at any number of epochs likely cannot alter it.

This fact is illustrated by Figure 10 (top), which shows that the age gradients are maintained to $z = 2$. Indeed, they persist

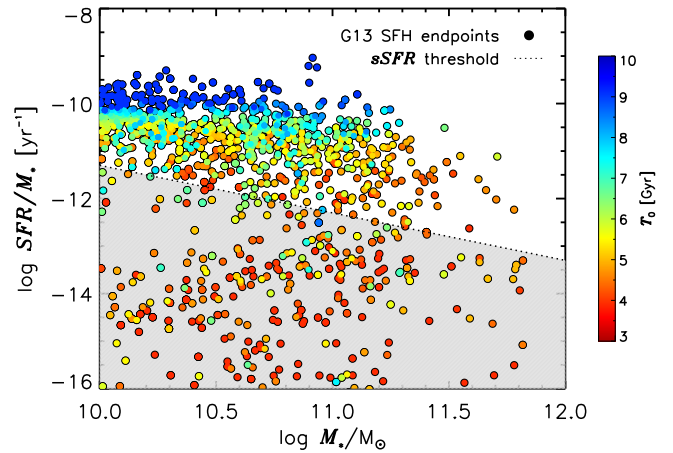


Figure 9. G13 $z \approx 0$ SFMS colored by half-mass times, T_0 . Gradients at fixed mass thus reflect age gradients. These persist at all M_* over many epochs (Figure 10) and so are key to the model’s interpretation of the SFMS. They need not have emerged and cannot be imposed without violating other constraints (O13). The dotted line shows the SFR limit ($0.05 M_\odot \text{ yr}^{-1}$; see G13) above or below which the model reproduces a galaxy’s measured M_* and SFR, or measured M_* but arbitrarily low SFR, respectively.

to $z \gtrsim 4$ (not shown), where horizontal stratification begins to dominate over the vertical structure illustrated here.

Two points are notable:

First, the appearance of T_0 gradients means that G13 sees σ_{SFMS} as encoding *long-term differentiation* of galaxy SFHs. That is, σ_{SFMS} is not (or at least need not be) due to short-timescale, $t \ll t_{\text{Hubble}}$, “weather-like” variations in SFRs. Instead, as in Speagle et al. (2014), even star-forming galaxies at fixed $M_*(t)$ have—and have always had—a wide range of histories: Some reached half that mass many Gyr ago, others quite close to the epoch of observation. Thus, *equal-mass galaxies on the SFMS need not form an evolutionary cohort*. Figure 11 (detailed below) reinforces this claim.

This finding accords qualitatively and quantitatively with new SFH measurements by Dressler et al. (2016). It is also similar to hydrodynamic numerical results from Tacchella et al. (2016) showing the SFMS scatter to reflect SFR variations of $\sim 0.5 t_{\text{Hubble}}$ (albeit at epochs when t_{Hubble} is only a few Gyr). We return to that work in Section 3.3.1.

Of course, in reality, starbursts and observational errors contribute to σ_{SFMS} . Hence, G13 yields something of an upper-bound to the amount of differentiation that quantity encodes, and thus the timescale of SFH segregation. Regardless, G13-like models suggest that *a significant portion of the SFMS scatter reflects galaxy diversification on Hubble-like timescales* (as seen by Dressler et al. 2016), and is therefore a manifestation of a mechanism operating globally and probably at very early times (e.g., initial conditions; Section 4.1).

The second point is that while G13 sets the SFMS scatter mostly as an age/half-mass-time gradient, it does not do so exclusively: There are late-forming galaxies lying well below the SFMS in Figure 9, and early-forming ones lying on it (though none at the very top). This subtle-but-real violation of the age-gradient interpretation—i.e., the breaking of age/mass rank-ordering—is critical, and where G13 departs from Speagle et al. (2014).

The SFHs that most dramatically drive this phenomenon are precisely those posited by O13 and subsequently spectro-photometrically identified by Dressler et al. (2016). They

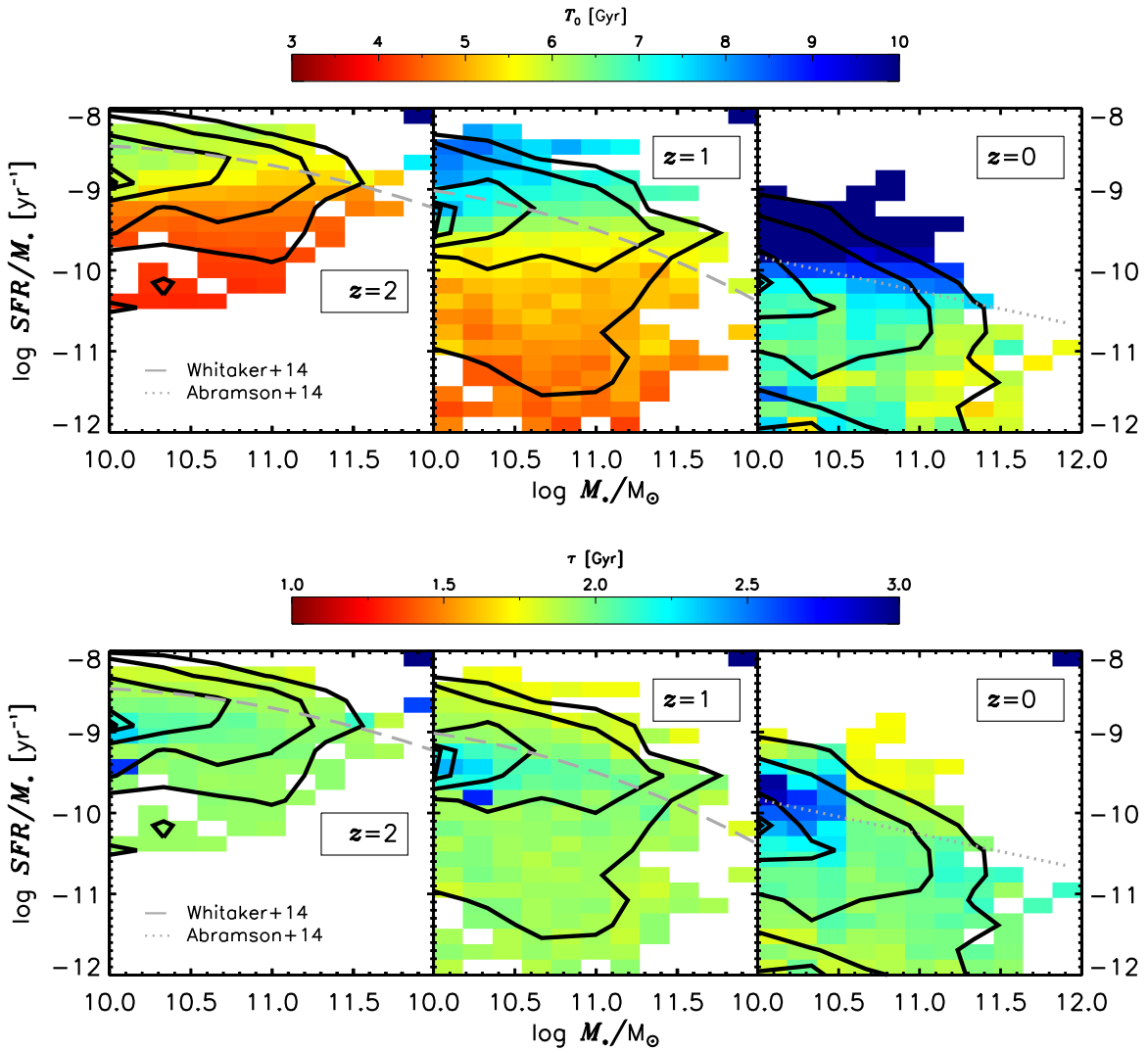


Figure 10. From left: **G13** model projections onto the sSFR– M_* plane at $z \in \{2, 1, 0\}$. Density contours—5%, 16%, 50%, 84%, 95% of max—roughly highlight the SFMS. For comparison’s sake, SFMS fits from (Whitaker et al. 2014, dashed) and (Abramson et al. 2014, dotted line) are overplotted in gray. Colors in the top and bottom panels reflect the mean T_0 and τ values of SFHs in 0.2×0.2 dex boxes, respectively. Rainbows at fixed M_* (top) reveal that the age gradients seen in Figure 9 (top-right panel here) persist to $z = 2$, implying that long-term SFH differentiation is and has always been a major source of the SFMS scatter. Interestingly, τ shows no such cross-sectional stratification, suggesting that SF *timescale* is not a source of σ_{SFMS} . However, this apparent decorrelation is due to the mixing on the SFMS of SFHs terminating at different $M_*(z=0)$ at all $M_*(t)$. This washes out real longitudinal trends in τ that are critical to diversifying true evolutionary lineages (Section 3.3; Figures 11–13). An 0.3 dex cross-calibration offset has been applied to the **G13** sSFRs in all panels (see Figure 2). This Figure is analogous to Figure 2 of Tacchella et al. (2016), left.

correspond to the $\sim 25\%$ of galaxies with present-day $M_* \gtrsim 4 \times 10^{10} M_\odot$ that had late-rising but also narrow SFHs. That is, SFHs that rose and fell *faster* than the SFMS.

Single-parameter SFH models cannot describe this behavior: they cannot attain sufficiently high sSFRs ($> 1/t$) at $z \sim 1$ and low ones today. To meet the first criterion, they must be constant or late-and-continually rising, automatically precluding them from meeting the second (O13 Figure 1). The same mathematics implies that σ_{SFMS} could *only* reflect an age gradient in a universe containing only those SFH forms (or ones confining galaxies to the SFMS; Speagle et al. 2014).

O13 showed that we live largely—but not entirely—in that universe. To make up the difference, a two-parameter form with independent characteristic times and timescales was required. Lognormals fit this bill (through T_0 and τ , respectively) with significant physical implications: The resultant ability for SFHs to fall faster than the SFMS translates to allowing galaxies to exit

the star-forming population *without imposing explicit quenching*. Hence, this feature is one of the model’s core axioms (Appendix A), and the fact that it produces the general trend predicted by single-parameter SFHs (accurate for many galaxies) but also known violations (Dressler et al. 2016) is one of its key—and perhaps most informative—successes.

The model cannot speak to the environments of the galaxies to which late-peaking, narrow (i.e., *young*) SFHs are ascribed, but O13 showed that they are probably isolated. Dressler et al. (2016) later verified that that, while present, their abundances decrease in dense regions, complementing Poggianti et al. (2013a)’s finding of higher fractions of older objects in clusters. These results suggest that SFH diversity is fundamentally related to the (initial) configuration of the matter density field. This seems compatible with Λ CDM structure growth, where a halo’s initial density/concentration maps closely to its mass accretion history (e.g., Lacey & Cole 1993; Avila-Reese

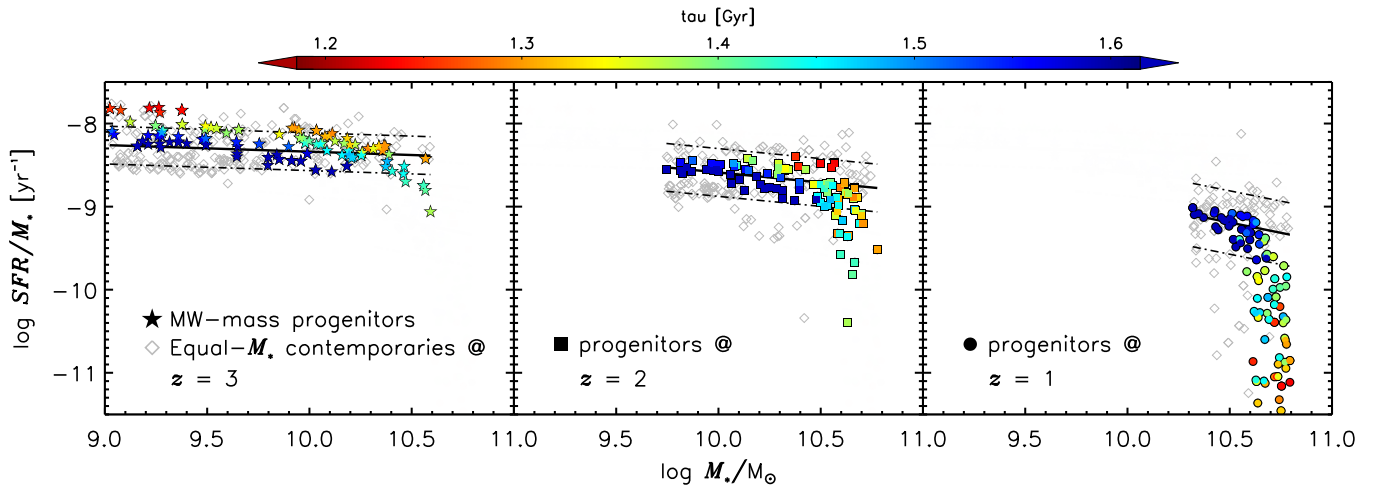


Figure 11. Colored points: **G13** predictions for the locations of galaxies that actually evolve into today’s MW-mass systems in $\text{sSFR}-M_*$ space at $z = 3, 2$, and 1 (from left; stars, squares, circles, respectively). Color coding reflects each SFH’s τ value. In this *longitudinal* view, a clear relationship emerges between τ and sSFR : galaxies with compressed histories (small- τ) start high on the SFMS and then fall below it, while galaxies with extended histories (large- τ) remain at or slightly below its midline. This correlation—reminiscent of observed and simulated trends with gas and/or stellar surface density (Ostriker & Shetty 2011; Genzel et al. 2015; Tacchella et al. 2016; Zolotov et al. 2015)—vanishes when the full SFMS is plotted (Figure 10, bottom). This is because galaxies that *will not* evolve into MW-mass systems cohabit the locus with true progenitors at all M_* (gray diamonds). Those systems display different τ trends depending on their destination, so mixing them with true MW-mass progenitors washes-out the signal shown by the colored points here. We replot all loci in all panels (faded) to show the full time evolution of the SFHs/SFMS (cf. Tacchella et al. (2016), Figures 2 (right) and 7).

et al. 1998; Lemson & Kauffmann 1999; Wechsler et al. 2002; Diemer & Kravtsov 2015). Section 4 revisits this point.

3.2.2. What about τ ?

As just mentioned, T_0 —the lognormal half-mass-time—is important to generating σ_{SFMS} in the **G13** scheme. Yet, as also just mentioned, a key feature of the model is that it treats τ —the SFH width—independently. So, what does it do?

Interestingly, the bottom panel of Figure 10 shows almost no correlation between τ and sSFR at fixed stellar mass. Thus, the **G13** model suggests that formation timescale—unlike formation time—does not play a large role in generating σ_{SFMS} , at least not when the locus is examined in the frame of the observer (Section 3.3).

This behavior is interesting because correlations between, e.g., galaxy structural parameters and their location on the SFMS at fixed M_* appear weak in some works (Wuyts et al. 2011; Bluck et al. 2014; Tacchella et al. 2016, but see Figure 13, below).¹² Hence, ironically, the lack of a strong gradient in τ at fixed M_* across the SFMS suggests that it, and not T_0 , might best-link the **G13** SFHs to observed galaxy structural properties. We explore this idea and the consequences of moving out of the observer’s frame in the next section.

3.3. New Observational Spaces

So far, we have studied how the **G13** model expresses itself in the context of canonical observations in galaxy evolution. Even when discussing its structure, we have shown projections into observed frames.

This approach is necessary to validate the model. Yet, because it is longitudinal—containing information not present in the data on the (theoretical) time-evolution of individual galaxies (Section 2.5)—we need not limit ourselves to cross-

sectional snapshots—to “connecting the dots” between different systems at different epochs—to learn from **G13**. Indeed, we need not limit ourselves to data at all.

In this section, we first show how moving out of the frame of the data reveals new physical insights. We then apply those insights to a new space—galaxy bulge mass fractions—and obtain results remarkably in-line with observations. If these exercises do not illustrate successes of the **G13** model, they provide good avenues by which to assess its ultimate validity.

3.3.1. A Longitudinal View of τ

Figure 10 (bottom) shows no τ trends at fixed mass along the SFMS. Yet, this presentation has limited virtue in elucidating how τ —and therefore what physics—differentiates galaxies because not all equal- M_* galaxies on the SFMS are evolving toward the same destination (Section 3.2.1). When galaxies that *will* ultimately evolve to a similar final mass are selected a priori, the **G13** τ trends change dramatically.

Figure 11 shows the evolution of all SFHs terminating near the Milky Way’s (MW’s) stellar mass ($10.6 \leq \log M_* \leq 10.8$) that had $\log M_*(z=3) \geq 9$ (to simplify the illustration). At early times, there is a clear preference for systems with low τ —i.e., narrow SFHs—to lie high on the SFMS, while those with high τ —elongated SFHs—lie low. At late times, however, this trend reverses, such that low- τ galaxies lie well below their high- τ counterparts. So, if such true progenitors could be isolated, what appeared as a decorrelation in cross-sectional data would transform into a clear trend.

On reflection, this trend *must* be present (Section 3.1.4, Appendix C): If a galaxy is to reach a given final mass before others do, it must grow rapidly early on—surging ahead of its eventual peers (middle panel)—and then stop, moving from a high to low position on (and then off) the SFMS (consistent with the findings of Marchesini et al. 2014). However, this trend is not apparent in the data because galaxies that will ultimately evolve into systems of different final masses

¹² This is not true of gas properties such as SFR surface density (e.g., Wuyts et al. 2011; Nelson et al. 2016; Tacchella et al. 2016), or when considering galaxies of all sSFRs (e.g., Omand et al. 2014).

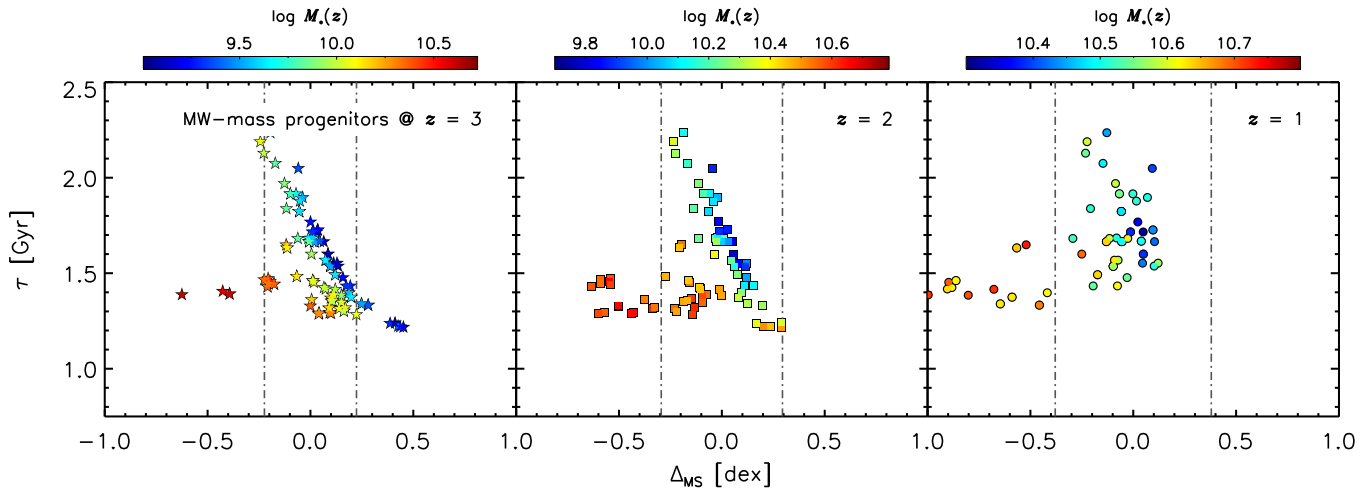


Figure 12. Relationship between τ and distance from the SFMS midline (Δ_{MS}) at the epochs in Figure 11. Points correspond to that figure but color-coding reflects M_* here, with redder points being closer to their $z = 0$ mass. This is the mathematical definition of quenching, but such systems also lie toward the left (underproductive) region of the diagram—the observational definition of quenching. Conversely, bluer points, which will grow substantially in the interim, span all positions on the SFMS. Their τ values tend to anti-correlate with Δ_{MS} such that more rapidly growing (low- τ) star-forming systems lie high while more slowly growing (high- τ) systems lie low. This trend is similar to those seen when τ is replaced by (inverse) Sérsic index or SF/stellar mass surface density (e.g., Schiminovich et al. 2007 (cf. their Figure 17), Wuyts et al. 2011 (cf. their Figure 4)), or gas depletion timescale (Tacchella et al. 2016, cf. their Figure 7). Combined with results shown below (Figure 13), this suggests τ encodes a density-like property. Since systems do not evolve in τ (i.e., along the y-axis) in G13, a prime candidate is the *initial* (baryonic) density in some physical aperture defined at a galaxy’s birth.

commingle with the true evolutionary cohort and obliterate this signal (gray diamonds in Figure 11).¹³

Note, however, how much real estate is covered by what will ultimately be galaxies of comparable stellar mass: At $z = 3$, such systems span *more than a factor of 30* in that quantity. This number is not inconsistent with theoretical results based on abundance-matching. Using the tool supplied by Behroozi et al. (2013a), we find the 2- σ spread of *halo* masses to be a factor of ~ 16 for the $z = 3$ progenitors of today’s MW-mass halos [$M_{\text{halo}}(z = 0) \simeq 10^{12} M_{\odot}$; e.g., Watkins et al. 2010]. This grows to a factor of ~ 18 in *stellar* mass assuming a 0.2 dex scatter in $M_*(M_{\text{halo}})$ (Behroozi et al. 2013b). As such, the factor of ~ 30 suggested by the model seems not so large as to rule it out on number-density conservation grounds (i.e., by requiring a wildly high merger rate).

Regardless, Figure 11 depicts a simple encapsulation of G13’s core astrophysical proposition: *the interesting physics in galaxy evolution is revealed by contractions of factors of 30 in M_* to eventual factors of < 2* . Because mean-based techniques such as SFMS integration (Peng et al. 2010; Leitner 2012) average over *that* process, these approaches may neglect critical parts of the story, blinding themselves to sources of galaxy diversity beyond quenching. In the extreme, the G13 model suggests that there is little-to-no astrophysical utility in describing the average SFH of galaxies at fixed $M_*(t)$.

A final note in this vein: Examining a plot similar to Figure 11 but color-coded by T_0 shows less (or no) vertical stratification, but instead greater differentiation along the mass axis. This trend is also expected—more-massive galaxies at any epoch should have earlier half-mass times almost by definition—but it is *orthogonal* to that seen in the full G13 SFMS

(Figures 9 and 10 top), again emphasizing the limitations of that locus as a window on evolutionary lineages.

3.3.2. Toward Physics

When examined from the perspective of galaxies that will evolve into similar-mass systems, Figure 11 shows τ to be a potentially important hook for physics. Figure 12 attempts to clarify this by moving into the frame of the SFMS. Here, we show the offset of each SFH from the SFMS midline as a function of τ at $z \in \{3, 2, 1\}$. The points are coded now by $M_*(z)$, revealing how close they are to “quenched” (more-properly “finished,” since the mathematical condition is $M_* = M_{\text{final}}^*$) at any epoch.

The trend described in Section 3.3.1 is more obvious here: galaxies with extended SFHs—large τ —sit near the mean SFMS ($\Delta_{\text{MS}} = 0$), while those with narrow SFHs—small τ —occupy all positions across it with M_* being the determining factor. The tail of low- τ high- M_* galaxies drifting below the SFMS (to the left in Figure 12) at any epoch might be seen as quenching, but G13 suggests it is better seen as a direct consequence of the fact that, if a galaxy is to reach its final mass early, it must have a compressed SFH (Wellons et al. 2015).

The most interesting objects are those with low- τ and low- M_* (bottommost blue points in Figure 12). These objects are rapidly shooting across the top of the SFMS in Figure 11. This behavior is potentially instructive.

If we make the reasonable assumption that τ —the “span” of the SFH—to some extent reflects gas consumption timescales, which are perhaps *set* by gas surface densities (or, in a globally averaged sense, filling factors Kennicutt 1998; Bigiel et al. 2008; Genzel et al. 2010; Ostriker & Shetty 2011), this figure is evocative of the findings of Genzel et al. (2015), Tacchella et al. (2016), and compatible with those of (Zolotov et al. 2015, see their Figures 12 and 13). That is, galaxies with compressed SFHs are those that consume gas rapidly perhaps *because* it is in an extremely dense configuration.

¹³ The SFMS is a train: people in different cars get off at the same station, people in the same car get off at different stations. Figure 10 (bottom) tracks the train and reveals a τ jumble. Figure 11 tracks only passengers going to the same stop and this reveals a meaningful trend. If a third parameter were found that mapped better to “time on train” than $M_*(t)$, τ should correlate with it in a 3D version of Figure 10.

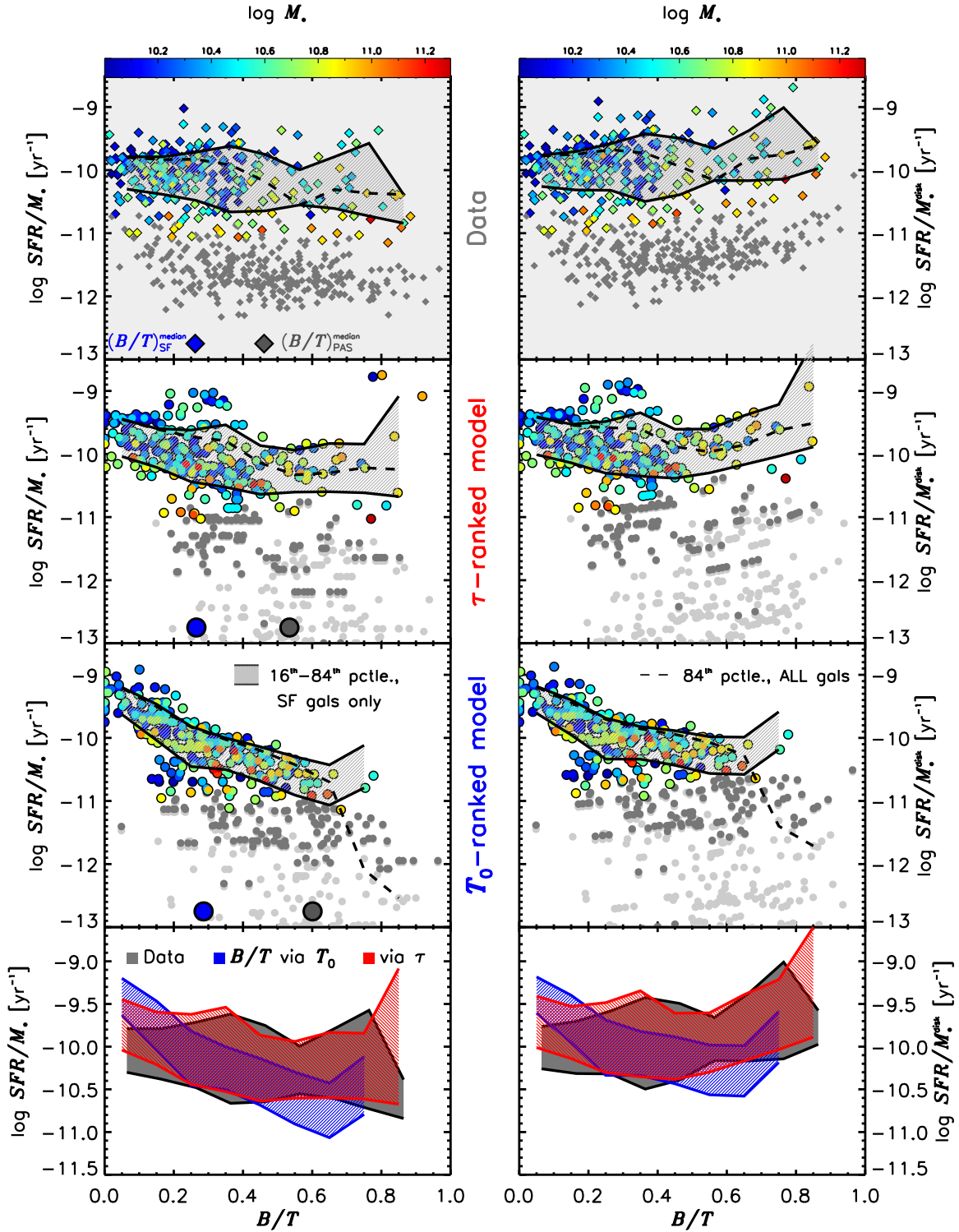


Figure 13. Projections of the $z = 0$ SFR/M_* - and $SFR/M_*^{\text{disk}}-B/T-M_*$ planes (left/right columns, respectively). Data are shown at top (gray panels; Brinchmann et al. 2004; Gadotti 2009). Recreations using τ or T_0 to assign B/T s from the top panel to the G13 SFHs are shown in the second and third rows, respectively. Colored and gray points show “star-forming” and “non-star-forming” galaxies/SFHs, respectively ($SFR < SFMS - 2\sigma_{SFMS}$; light gray points show G13 SFHs with $SFR(z_{\text{obs}}) < 0.05 M_\odot \text{ yr}^{-1}$; Equation (4)). Shaded regions highlight the $1-\sigma$ spread of the star-forming systems only. Dashed lines in the top six panels show the 84th percentile of all galaxies (including sSFR upper-limits; attempts to capture the “shower rod” from which the “curtain” of galaxies descends). Point colors reflect present-day M_* . As summarized in the bottom panels, τ -ranking—assigning the narrowest G13 SFH per mass-bin to the highest B/T , and vice-versa—reproduces the data trends far better than T_0 -ranking—where the earliest-forming G13 SFH gets the highest B/T . This is true in terms of (1) the slope and width of the $SFR/M_*^{\text{disk}}-B/T$ relation for SF galaxies; (2) the maximal B/T of SF galaxies; (3) the location of the upper-envelope of all galaxies; and (4) the median B/T for the SF and non-SF populations (blue/gray symbols, respectively, on the top three x-axes at left).

This scenario, which echoes (Holmberg 1965, see his Figure 6), naturally explains trends from, e.g., Franx et al. (2008), Williams et al. (2010), Valentiniuzzi et al. (2010), Fang et al. (2013), Omand et al. (2014), Barro et al. (2015), and Morishita et al. (2016): Dense *star-forming* galaxies evolve rapidly, so, at any epoch, galaxies that appear to be quenching will be dense, leaving dense stellar configurations behind to be preferentially identified as passive galaxies in a future epoch. However, because G13 contains no such physics, galaxies do not evolve toward denser states in this model. As such, the above picture need not entail a “compaction” event (though it could; Barro et al. 2015; Zolotov et al. 2015). Instead, Figures 11 and 12 suggest quenching is the “peeling-off” of the densest tail of the star-forming population at any epoch, explaining why the (M_* and) surface density threshold of (Barro et al. 2015, see their Figures 1 and 7) falls with time.

Ultimately, both scenarios—compaction, in which SF galaxies evolve toward a denser state before quenching, and “peeling,” in which the densest SF galaxies simply “finish first”—may be correct. Our fundamental argument is that, because the model has no density parameter, yet captures something like the density-related trends seen in data through τ , the fundamental physics of interest is whatever sets the τ spectrum: it is that which determines if and when a galaxy will undergo either of the above processes.

An attempt below to tie G13 to data similarly outside its formal scope supports this interpretation.

3.3.3. Bulges and Disks

G13’s stance on τ is that it should appear unrelated to a galaxy’s position on the SFMS in cross-sectional data—i.e., observations—but should still be important to diversifying *actual members* of an evolutionary lineage. If we could identify an observable that maps to τ we could define cuts in, e.g., the sSFR– M_* plane that isolate true progenitors but are otherwise invisible in the data.

Oddly, the lack of an apparent correlation of star formation activity with τ in cross-sectional samples may hint at which observable(s) it is linked to. For example, at similar masses, Wuyts et al. (2011) showed that trends in galaxy Sérsic indices are also weak at fixed M_* across most of the SFMS. Fang et al. (2013) and Barro et al. (2015) (mentioned above) observed something similar for stellar mass surface density, which Tacchella et al. (2016) also see in simulations. These weak correlations are well-reflected by τ ’s behavior across the SFMS (Figure 10, bottom), but not T_0 ’s, which shows strong gradients at fixed mass (Figure 10, top).

Going one step further, if we associate Sérsic index or stellar mass surface density with bulge mass fractions, B/T , then the connection posited above between τ and a gas consumption or compaction timescale would also encourage us to link τ to B/T , since compact galaxies with short formation timescales likely evolve into bulge-dominated systems (Schiminovich et al. 2007; Barro et al. 2013, 2015; Steinhardt & Speagle 2014; Zolotov et al. 2015; Tacchella et al. 2016).

We test this association in Figure 13. This Figure is complex, so we review it in some detail.¹⁴

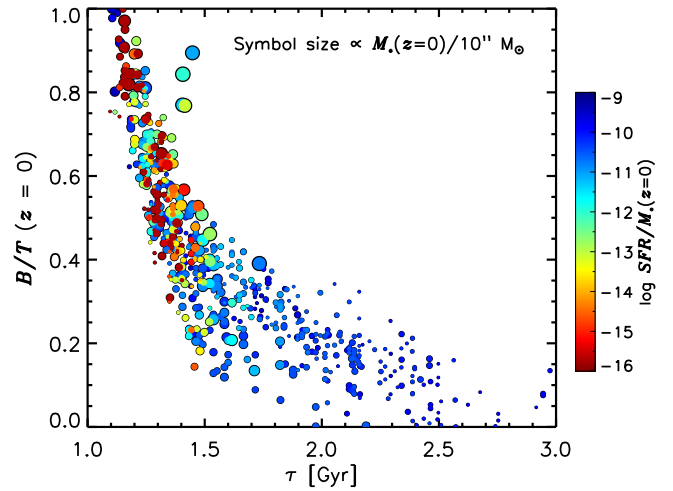


Figure 14. Present-day $\tau \mapsto B/T$ mapping that produces the second row of Figure 13. Symbol colors and sizes reflect each SFH’s $z = 0$ sSFR and M_* , respectively. The B/T assignment guarantees the correct symbol size– y -axis location relationship, but color-highlighted trends are emergent. These reveal G13 to reproduce known tendencies for star-forming and passive galaxies to span a range of B/T , but for the former be disk- and the latter almost free of pure disks (Schiminovich et al. 2007, A. Oemler et al. 2016, in preparation), implying that short formation timescales are key to bulge-building.

The first step in its construction is to assign each G13 SFH a B/T ratio. We do this in a manner similar to the “age matching” prescription of Hearin & Watson (2013).

Using a sample of high-quality $z \approx 0$ B/T measurements (Gadotti 2009) and moving in bins of M_* , we rank-order the data by B/T and the SFHs by either τ (second row) or T_0 (third row). We then simply map the lowest- τ or lowest- T_0 SFH to the highest B/T and continue down both lists. Physically, this corresponds to assuming either rapidly (low- τ) or early forming galaxies (low- T_0) become bulge-dominated, while slowly/late forming ones become disk-dominated.

We do this without respect to SFR. As such, the known $z = 0$ B/T – M_* relation—slope and scatter—is reproduced by construction, but the SF properties of high- and low- B/T systems is not. Examining either the trend of sSFR (SFR/M_*) or disk-mass normalized SFR ($\text{SFR}/M_*^{\text{disk}}$) with B/T (left and right columns, respectively) may therefore yield hints as to which G13 model parameter is most closely associated with bulge growth—or stellar surface density, or Sérsic index, or “compaction”—and enhance our physical interpretation.

The top row in Figure 13 (gray panels) shows the data using aperture-corrected SFRs from Brinchmann et al. (2004) (DR7 release). Colored points (coded by M_*) are “star-forming” galaxies defined by a $2\text{-}\sigma$ cut below the SFMS. As expected, high-mass systems have reduced SFR/M_* compared to low-mass systems, and high- B/T systems have reduced sSFR compared to low- B/T ones (highlighted by the gray shaded bands, showing the 16th–84th percentile spread of the star-forming population). As also anticipated (Abramson et al. 2014), high-mass, high- B/T systems have $\text{SFR}/M_*^{\text{disk}}$ more comparable to lower-mass or diskier systems. Note that the distribution of either quantity at fixed B/T is broad.

Comparing the data to the model shows that T_0 -ranking (third row) poorly reproduces these trends; both SFR/M_* and $\text{SFR}/M_*^{\text{disk}}$ fall too steeply with B/T , and loci at fixed B/T are too narrow. The former is true considering only star-forming galaxies, or the upper-limb of all galaxies (dashed black lines in

¹⁴ Lilly & Carollo (2016) provide an alternative test to what follows. While they work from an independent perspective, there is nothing obviously incompatible between our approaches or conclusions.

all panels). Further, the median B/T for *non*-star-forming galaxies is $\sim 15\%$ higher than that in the data (large symbols near abscissas).

In all respects, τ -ranking (second row) does a far better job. As summarized in the bottom row of Figure 13, the slopes and widths of the observed trends for star-forming (or all) galaxies are much more faithfully reproduced using this method, as is the median non-star-forming B/T ratio.

This outcome is precisely what is expected if τ is indeed linked to gas consumption timescales, or gas/stellar mass surface densities as Figures 11 and 12 suggest. As such, it deepens our sense that τ encodes the action of a key organizational principle in galaxy evolution.

Note: Figure 13 and the above discussion are not intended to convince the reader that there is, in reality, a 1:1 mapping between τ and a galaxy’s present-day B/T . Neither are we suggesting that we have found—or even attempted to find—an accurate description of the data using the G13 model. We are suggesting, however, that a model with no physics combined with very simple prescriptions to extend it to domains quite far from anything it was intended to describe can yield highly suggestive and potentially informative results.

We believe this fact alone is weighty. However, the results in Figure 13 combined with the phenomenological resemblance of τ ’s relationship to the SFMS and that of the above phenomena embolden us to posit that we are not looking at a mere description of the data in these cases, but in fact an explanation. That is, τ —or some optimal combination of T_0 and τ we have yet to discover—is not only correlated with, but in fact *controls* the eventual gas/stellar surface density, bulge fraction, or size of a galaxy, and therefore its probability of being “quenched” at any epoch/as a function of time. As such, we will use what we have learned to make a prediction.

Figure 14 shows the G13 $\tau \mapsto B/T$ mapping as functions of a galaxy’s present-day stellar mass (symbol size) and sSFR (symbol color). As demanded, low- τ objects have high- B/T and vice versa, but the shape of this relationship and its coloring are telling. The steepness at the low- τ end—dominated by more-massive, more-passive galaxies—is reminiscent of the spread in B/T s exhibited by early Hubble type objects (Dressler 1980; Kelson 1998; Fabricant et al. 2000). Meanwhile, the shallower trend at higher- τ shows the general diskiness of all late-type star-forming galaxies today (though they span many morphological subclasses, perhaps reflecting a range of formation timescales). Furthermore, this mapping leads to *some* high- B/T star-forming galaxies, *but no* (very) low- B/T passive ones (i.e., red points). This is true even though there are plenty of passive objects at masses dominated by pure-disks. These last subtleties are observed *facts* (e.g., Schiminovich et al. 2007; A. Oemler et al. 2016, in preparation); that the model produces them along with the grosser trends mentioned above further supports its validity.

Since τ describes a formation timescale, it should also correlate with α -element enhancement, for example—or, again, perhaps Hubble type. As such, we encourage other investigators to assemble such data and compare their observations to the predicted trend in Figure 14. We also suggest that, now or in the (not-too-distant) future, advanced SPS modeling may allow reliable determinations of τ . Hence, others might also wish to run such spectral/SED fitting using lognormal SFHs and compare their results to this diagram (or to assemble the relevant data when such technology is more robust). Finally,

we would encourage numerical simulators to produce a similar diagram to provide an independent prediction from which deeper physical causes can perhaps be identified.

We finalize our physical interpretation of G13 below.

4. DISCUSSION

The phenomenological slog above was intended simply to persuade the reader that the G13 model—a collection of just 2094 lognormal SFHs—is sufficiently good at reproducing sufficiently many observations across sufficiently diverse domains and sufficiently long stretches of time that its deeper physical and methodological consequences merit serious consideration. There are three possible consequences:

Astrophysical—Galaxy growth trajectories can diverge radically from inferences based on the evolution of scaling laws (e.g., the SFMS). They are not characterized by large discontinuities, more likely are (largely) defined by initial conditions, and may (largely) share a common terminating force. Or...

Metaphysical—Current data can be described from orthogonal—if not opposite—astrophysical perspectives. Such disagreements must be settled or understood before tests of more explanatory models can be fully contextualized. Or...

Epistemological—The information content of the observations examined above is (very) low.

We now explore each of these statements in turn.

4.1. Astrophysical Implications

4.1.1. Quenching is Uninformative

For those who view the G13 model as a (sufficiently) good description of the data, our immediate conclusion is that quenching—defined any way that suggests the removal from the star-forming population of galaxies with arbitrarily high sSFRs—is a needlessly limiting physical framework, if not an overly mechanistic interpretation of the facts. We say this because G13 entails no quenching physics, and indeed no star-forming “states” for galaxies to exist in at all—only a diversity of smooth SFHs evolving along independent trajectories that naturally rise and fall.

This statement is broader than it might first appear. For example, it can be rephrased as “Galaxies do not evolve *along* the SFMS.” As shown in Figure 11, there is no need for star-forming galaxies observed at fixed M_* at one epoch to appear as a similarly confined cohort at any other epoch. The SFMS therefore *shows you* galaxies evolving, but it does not *tell you* (much) about why they are doing it. In this sense, it and similar trends are like swarms of starlings: they emerge from flight paths of individual birds, but reveal few details about those trajectories (Section 4.2), and therefore little about how/why birds fly (or stay grounded). The histories tell the story; the loci do not. As such, the utility of inferring physics from the (evolution of) scaling relations is dubious. If galaxies are seen as evolving through the SFMS, the need for an ad hoc mechanism to derail them from that trend is removed.

Our argument is slightly deeper than this. The invocation of grow-and-quench scenarios is natural given cross-sectional data: when galaxy properties at a single epoch are collapsed onto 2D planes, the obvious correlations that result encourage causal interpretations—or at least encourage us to think we are seeing evolutionary trends. Yet this cannot be true: star-

forming galaxies at a given epoch are the progenitors of precisely *none* of the passive galaxies at that epoch. Hence, such conclusions are valid if and only if quenching must be rapid, a scenario we and others have here and elsewhere shown need not be the case (e.g., Dressler et al. 2013; Schawinski et al. 2014; Peng et al. 2015). If a distinction between longitudinal and (series of) cross-sectional samples is maintained, the need for/idea of quenching changes, or disappears.

One could argue at this point that the falling part of our histories *is* quenching. We would agree, so long as “quenching” means “the smooth crossing of an arbitrarily low sSFR threshold”; i.e., *normal galaxy evolution*.¹⁵ This might entail a combination of gas exhaustion and perhaps a mechanism to prevent some gas from condensing to fuel star formation (Conroy et al. 2015; Voit et al. 2015). The simplest interpretation of our model would be that the SFH declines—the phenomena requiring explanation—are due purely to the former: the gradual reduction in cold gas fractions over time (Geach et al. 2011; Popping et al. 2015). However, as populations of non-star-forming galaxies exist with non-negligible amounts of neutral gas (Chen et al. 2001; Young et al. 2011; Thom et al. 2012; Johnson et al. 2015), something akin to the latter is probably also at work. We emphasize that neither option is terribly exotic: although, e.g., supernovae and rapid AGN feedback must be important in some contexts, our findings suggest that they are generally sub-dominant to the above (concurring with Feldmann & Mayer 2015).

Again, one could argue that galaxies are known to suffer large discontinuities in star formation induced by, e.g., environmental gas stripping or starbursts (e.g., Gunn & Gott 1972; Dressler & Gunn 1983; Poggianti et al. 2016), so “fast-track” quenching clearly exists. We admit that the current G13 model can neither capture nor constrain the prevalence of such events. However, we would argue that the results above—and the data (Zabludoff et al. 1996; Quintero et al. 2004; Rodighiero et al. 2011; Peng et al. 2015; Dressler et al. 2013; Abramson et al. 2013; Dressler & Abramson 2015)—suggest that these are perturbations to an underlying narrative based on smooth/long-timescale SFR(t) correlations, and are therefore corrections to a G13-like scheme.¹⁶

4.1.2. Galaxy Evolution Is Driven by Initial Conditions

“Quenching is uninformative” can be rephrased in yet another way: “Environmental and internal effects may be inseparable” (see also De Lucia et al. 2012). This is because our model makes no such distinction; it simply assigns each galaxy a (T_0 , τ) pair at birth. If these quantities—“clocks” that set the speed of each galaxy’s evolution—were somehow correlated on supergalactic scales early on, the causal effects of internal and external influences on SFHs would be linked.

The immediate interpretation of the above (shared by Speagle et al. 2014) is that *initial conditions must substantially determine the course of a galaxy’s evolution*. We believe that

this is the fundamental implication of the G13 model, and that our conclusion regarding quenching is a consequence of it. If we are correct, from a paradigmatic standpoint, the most critical aspect of galaxy evolution is which initial conditions set (something like) T_0 and τ .

An obvious candidate is the configuration of the matter density field, specifically the initial baryonic overdensity in some optimal aperture. This choice is motivated by Figures 9–13, in which T_0 and τ (mostly τ) are seen to play roles similar to gas/SFR surface densities or gas consumption timescales as assessed by, e.g., Wuyts et al. (2011), Zolotov et al. (2015), and Tacchella et al. (2016). We believe that our model is also consistent with the work of Hearin & Watson (2013), which suggests that a halo’s age—the time at which it passes a mass, mass-accretion, or halo-membership threshold (see their Sections 4.1.2 and 6.3)—is a critical parameter, in addition to, and separate from, its total mass. As it seems reasonable for halo baryon overdensity to be well-correlated with such a threshold, these may ultimately be two ways to describe the same phenomenon.

Regardless, the action of both theoretical mechanisms is substantiated by data, either in the densities of many star-forming galaxies at high redshift (Barro et al. 2013, 2015), or the older stellar ages of passive galaxies at fixed M_* in clusters compared to those in the field (Thomas et al. 2005; Poggianti et al. 2013a; McDermid et al. 2015; Morishita et al. 2016), or the accelerated decline in the SFRD in overdense environments (Guglielmo et al. 2015).

To summarize, it is not that passive galaxies are “quenched,” but that they are “finished,” having had an accelerated assembly/formation history compared to star-forming counterparts because they grew up in denser regions of the universe. This scenario—which is not new, but a phenomenological rediscovery/synthesis of ideas from Holmberg (1965), Tinsley (1968), and Dressler (1980), and seems deeply related to the concept of assembly bias (e.g., Wechsler et al. 2006; Wetzel et al. 2007; Dalal et al. 2008; Feldmann et al. 2016; Saito et al. 2016; and relatedly Aragon-Calvo et al. 2016)—naturally links environmental density, internal galactic density, and the characteristics of galaxies in overdensities as observed at effectively all epochs. From that starting point, everything else emerges.

So, while it would be helpful to know if AGN-related activity ultimately produces the (spread of) SFH declines we have parametrized in a physics-agnostic manner, we would still ask if the action of that mechanism is a symptom of an underlying organizational principle such as the above, and suggest that confirming *that link* is important. Indeed, simulators may already have the information to do so: We would encourage them to see which aspects of their models set (parameters analogous to) T_0 and τ , and whether identifying trends at that level might enhance their understanding of deeper physics.

4.2. Metaphysical Implications: Too Many Descriptions to Permit Explanations

Issues will remain, however, independent of the outcome of such tests. Given results from, e.g., Peng et al. (2010), Behroozi et al. (2013b), and K14, multiple paradigms of galaxy evolution would seem to (at least have the potential to) describe all of the data we have examined as well as ours. Those works outline two other perhaps equally effective yet

¹⁵ In some sense, we are suggesting that a universal null quenching scenario be adopted entailing Hubble-timescale processes—which are rapid at high- z —that should be rejected before other explanations are sought.

¹⁶ We acknowledge that sufficiently narrow, smooth SFHs leading to red galaxies at $z \gtrsim 2$ may be difficult to distinguish observationally from abrupt truncations of SFMS-driven growth or powerful starbursts. Comparisons of the spatial distributions of galaxies binned by M_* and SFR might distinguish these scenarios in the era of *WFIRST*, however. (see H. Ferguson’s talk at https://conference.ipac.caltech.edu/wfirst2016/system/media_files/binaries/19/original/ferguson_wfirst2016.pdf?1457376009).

largely incompatible frameworks operating at the descriptive level of the G13 model.

The first is the grow-and-quench paradigm against which we have principally contrasted our model (e.g., Peng et al. 2010). O13 and Dressler et al. (2016) demonstrate that there must be some galaxies that do not follow SFMS-defined growth histories (Section 2.1). However, if those exceptions are not seen as overly damaging, it seems likely that a model in which SFMS or $M_*(M_{\text{halo}})$ evolution reveals *governing principles* behind galaxy growth—halted by some agent uncorrelated with initial conditions—is compatible with everything we have so far explored. As such, we do not rule it out as viable description of galaxy evolution.

The second is the quasi-stochastic model of K14. In some aspects, this paradigm is quite similar to ours, yet in others it could not be more different.

The most notable similarities are (1) that key physics is encoded by the *scatter* in, e.g., the SFMS; (2) that the cessation of star formation comes from the same phenomenon that sustains it (there is no explicit quenching); and (3) that controlling mechanisms must be inferred at a high level from the formalism of the model itself (in K14, the distribution of $\delta\text{SFR}/\delta t$ updates (σ); in G13, the distribution of (T_0, τ)).

Yet, K14 reaches these conclusions by positing that most of galaxy evolution is governed by non-deterministic SFH “updates” to equilibrium conditions. These can be arbitrarily discontinuous and explicitly forbid parametrization. That is, the act of parametrizing the K14 SFHs—e.g., by a lognormal—destroys the ability to learn about the core physical process, which is whatever sets the spectrum of quasi-stochastic updates. K14 showed that these features make that model a good description of the histories of local dwarfs (as directly measured by Weisz et al. 2014), which our model does not contain, but likely also does not describe (Section 2.2). In these senses, it is *opposite* the G13 model.

That said, there is some suggestion that the spectrum of σ s in K14 is closely tied to initial halo conditions, specifically density (D. Kelson 2016, private communication). If so, the *prima facie* lack of a way to correlate/corral SFHs, e.g., spatially—an observed fact and something we believe is encoded by (T_0, τ) —may be alleviated, or at least is likely to derive from the same mechanism we are espousing. One can see how this might work: Where G13 suggests a galaxy is “finished,” K14 posits that it “last updated a long time ago.” Yet, both would agree, e.g., that passive (old) galaxies at any epoch should be denser than star-forming (young) counterparts (Martig et al. 2009; Valentinuzzi et al. 2010; Fang et al. 2013; Poggianti et al. 2013a, 2013b; Barro et al. 2015; Morishita et al. 2016) because they reflect the denser nature of the universe at the time they were last active.¹⁷

The upshot is that descriptions in which galaxies:

1. evolve together along a physically informative SFMS and are quenched (e.g., Peng et al. 2010);
2. have diverse histories that naturally rise and fall as determined by initial conditions (G13);

3. grow quasi-stochastically, such that the SFMS reflects a foregone mathematical conclusion (K14);

may be equally compatible with many key statistical observations in galaxy evolution. This suggests either that the data are unconstraining (see below), or that *formalism substantially influences physics*: If a model posits that galaxies evolve *along* the SFMS [or $\langle M_*(M_{\text{halo}}) \rangle$], it cannot escape invoking what models with emergent SFMSs (G13, K14) see as extraneous quenching prescriptions. It must also neglect potentially revealing SFH diversity because it contains no formal resource but a description of the mean (without an *ansatz* for scatter).

At root, we are arguing that there is a meaningful lack of understanding (or at least consensus) regarding *how to describe* galaxy evolution (see Taylor et al. 2015 for extensive further discussion in the context of the $z \approx 0$ stellar mass function). As such, even the most basic questions—How predictive of a galaxy’s SFR(t) is its initial halo overdensity? Does the SFMS describe an emergent or driving phenomenon? How far does stochastic growth take you?—are not well (or widely) understood. If so, such ignorance may supersede questions at the level of “Is strangulation or AGN feedback the dominant quenching mechanism?”

This leads us to the third implication of our results.

4.3. Epistemological Implications

4.3.1. Many Core Metrics Have Limited Discriminatory Power

If nothing else, the results presented here, and in G13, Abramson et al. (2015), and Dressler et al. (2016) suggest that a loosely constrained collection of SFHs governed by no explicit physical prescriptions does a fair job at describing many observations generally approached using more sophisticated techniques, such as semi-analytical modeling or hydrodynamic simulation. This is in no way to diminish the latter—they are the only routes to a first-principles explanation of the details of galaxy evolution. However, as very few such details *must be understood* to reproduce the data explored here, the latter would seem insufficient to confirm/falsify these models.

For example, at least at $z \lesssim 2$, σ_{SFMS} has been shown to reflect: (1) Hubble timescale differentiation of the galaxy population (G13; this work); (2) burstiness on 10^7 years timescales (e.g., Hopkins et al. 2014); and (3) fluctuations on arbitrary timescales (K14). If this metric accommodates *all possible timescales* for SFH variability, then it alone is a poor test of physical models with implications for that timescale.¹⁸

Ultimately, we are suggesting that the data examined here and in many other works may constrain some basic mathematical properties of SFHs—e.g., that they rise and fall asymmetrically in time with independent half-mass-times and widths (Appendix A)—but not much else. To learn more, different information is needed.

4.3.2. Toward More Discriminating Tests

To progress, the dimensionality of the problem must be increased on both small and large scales.

Regarding the former, it may be that critical astrophysics are hyper-local (e.g., Bigiel et al. 2008), so integrated quantities

¹⁷ G13 and K14 may be intimately related by the central limit theorem. K14 uses it explicitly to show that the shape and evolution of the low-mass SFMS may reveal nothing but mathematical inevitabilities. G13 invokes the theorem implicitly—lognormals are the limit for random multiplicative processes—but, taking our *ansatz* seriously and letting K14 describe *parts* of galaxies, there might be a deep formal connection between the two models.

¹⁸ Though future, e.g., deep spectroscopic observations at many z may substantially increase the potency of this constraint by providing additional information on stellar populations (Section 4.3.2).

—SFR, M_* , Z , r_e , etc.—are too low-resolution to be discriminating. Beyond diluting information, reducing galaxies to these numbers may cripple our ability to connect progenitors to descendants. Accounting for at least the fact that galaxies comprise bulges and disks changes the interpretation of the SFMS (Abramson et al. 2014). It must also change assessments of which galaxies evolve into which others. Where else will similar moves be even more edifying?

The acquisition of high spatial-resolution IFU-like spectroscopy for even a few thousand galaxies over a fair range of cosmic time could be decisive. *James Webb Space Telescope* (JWST) and forthcoming extremely large ground-based telescopes (ELTs) will provide these data. Access to, e.g., gas/stellar metallicity *gradients* at many epochs will not only support much more sensitive tests of the phenomena that control gas processing—and thus galaxy growth—but may also provide better ways to link progenitors and descendants and thereby stitch together different cross-sectional samples. Both will greatly enhance our knowledge of the details of galaxy evolution.

Pilot surveys of this nature are underway—GLASS (Jones et al. 2015; Treu et al. 2015; Vulcani et al. 2015b), GASP (B. Poggianti et al. 2016, in preparation), ATLAS3D (Cappellari et al. 2011), 3D-HST (Brammer et al. 2012; Nelson et al. 2016), CALIFA (Sánchez et al. 2012), MANGA (Bundy et al. 2015), SAMI (Allen et al. 2015), KROSS (Magdis et al. 2016)—and proving quite powerful. Hence, the “dimensionality problem” at small scales may itself soon be reduced.

Doing so would amount to winning most, but crucially not all of the battle. To achieve this, any physics inferred/ancestral mapping provided by JWST and the ELTs must be contextualized on *large scales* using clustering/“two-point” statistics. This is because models must not only generate galaxies with the right (M_* , SFR), but also put them in the right places at the right times. There is strong theoretical and empirical evidence for correlations between baryonic and halo properties at single epochs (e.g., Weinmann et al. 2006; Guo et al. 2010; Hearin & Watson 2013; Kawinwanichakij et al. 2014; Tinker 2016). Truly viable models must also recover the *evolution* of these correlations (e.g., Conroy et al. 2006) over large dynamic ranges in M_* , SFR, and time, thereby overcoming the substantial diversity in halo or stellar mass growth histories leading to the same (M_* , SFR) coordinate at any one moment (Behroozi et al. 2013a; Dressler et al. 2016; Tinker 2016).

This will require deep, wide surveys that sample and tomographically track the growth of the full range of galaxy environments over much of time. *WFIRST* (Spergel et al. 2013, 2015) will provide such data, enabling the clustering measurements needed to discriminate between scenarios that equally reproduce M_* and SFR distributions, but are based on different mappings to underlying halo activity.

For example, the G13 framework implies that, at some early time, the galaxy correlation function must be biased toward systems of the *highest*—as opposed to lowest—SFRs (see also Feldmann et al. 2016). If this is not seen, quite simply, the model is wrong. Additionally, for the same reason, our model suggests that the UV sources *observed* during reionization may differ substantially from those *powering* it: the former live in overdensities and are therefore mature, while the latter live in less-dense regions and are thus much younger (and might have different spectra; Furlanetto et al. 2004a, 2004b; Wise et al.

2014; Davies & Furlanetto 2016; Stark et al. 2016). This need not be the case in SFMS-driven or quasi-stochastic models, where environmental effects either might not exist or might be washed out. Hence, high- z observations sampling a high dynamic range of environments should enable powerful tests of these paradigms (see Mirocha et al. 2016 for quantitative explorations in this vein).

5. SUMMARY

We have shown that the following observational metrics of galaxy evolution:

1. the evolution of $\langle \text{sSFR} \rangle$ for low-mass ($\log M_* \in [9.4, 10]$) galaxies since $z = 7$ (Section 3.1.1, Figure 3);
2. the evolution of the galaxy stellar mass function since $z = 8$ (Section 3.1.2; Figure 4);
3. the evolution of the slope of the SFR- M_* relation (the SF “Main Sequence”; SFMS) since $z = 6$ (Section 3.1.3, Figure 5);
4. the transition from “fast-” to “slow-track” quenching over cosmic time (Section 3.1.4, Figure 6);
5. galaxy downsizing (Section 3.1.6, Figure 8);
6. the bulge mass fractions of local galaxies as a function of SFR/ M_* or SFR/ M_*^{disk} (Section 3.3.3, Figures 13 and 14);

are naturally and well reproduced by a purely mathematical, explicitly quenching-free description introduced in (Gladders et al. 2013, “G13”) *that was not designed to match any of them*. In this model, galaxies have continuous, smooth, lognormal star formation histories (SFHs) constrained mainly by the evolution of the cosmic SFRD (Section 2.2; Figures 1 and 2). The cessation of star formation is thus linked strictly and identically to the same process causing SFHs to generically transition from rising to falling states, suggesting that—for most galaxies—no additional terminating process is necessary (see also Aragon-Calvo et al. 2016).

1. The model’s emergent interpretation of the SFMS scatter as *Hubble timescale diversification* of galaxy SFHs combined with the lognormal’s ability to rise and fall faster than this locus enable its successes (Section 3.2; Figures 9 and 10). They allow G13 to recast binary “star-forming” and “non-star-forming” labels as a spectrum of *extended* and *compressed* SFHs, none of which must be subject to qualitatively different physics (Section 3.1.4; Figure 6), or any beyond that which causes SFHs to turn over. This implies:
 - i. Progenitors of galaxies within a factor of 2 of the Milky Way’s current stellar mass may span a factor of ~ 30 at $z = 3$; i.e., galaxies do not evolve “along” the SFMS, but “through” it (Section 3.3.1, Figures 11 and 12).
 - ii. If so, the utility of that scaling relation in characterizing evolutionary trajectories or gleaning physics is dubious: the mean SFH (M_* , SFR; t) is unrepresentative of individual SFHs and insensitive to potentially critical mechanisms that differentially accelerate galaxy growth.
 - iii. Trends displayed by (series of) single-epoch/cross-sectional data may be *orthogonal* to true

evolutionary/longitudinal trends and may therefore support misleading inferences.

- iv. Star formation *timescales*, τ , present a key manifestation of this issue, appearing uncorrelated with SFRs at fixed M_* when examining the full galaxy population at $z \lesssim 2$, but showing strong trends once systems cohabiting the SFMS with true evolutionary cohorts but destined for different final masses are removed (Sections 3.3 and 3.3.2; cf. Figures 10 (bottom) and 11).
- v. We interpret (iv) to mean that understanding whatever sets the *width* of (lognormal) SFHs is critical to placing observations in the context of a comprehensive narrative of galaxy evolution.
2. Toward that end, we posit that τ is closely tied to baryonic surface densities, gas consumption timescales, or “compaction” phenomena based on highly suggestive comparisons between trends found here and those in the literature (Sections 3.3.2 and 3.3.3; Figures 12 and 14).
3. Since G13 fixes τ once for each SFH for all time, we infer from the above that *initial conditions are powerfully predictive of a galaxy’s SFH*, and can serve as a broader organizing principle in which quenching—here, a Hubble-timescale phenomenon—might be more instructively couched (Section 4.1). We suggest the baryonic overdensity in some optimal aperture at $z > 8$ (the G13 start time) as a potential characterization of those conditions.
4. If so, in many instances, there need not be a distinction between environmental and internal mechanisms for quenching, which is better set as the “peeling off” of the densest tail of the star-forming population at any epoch. Internal galactic (gas) densities correlate with those on super-galactic scales, leading to spatially confined cohorts of equally rapidly evolving, overdense (i.e., bulge-dominated) galaxies that all “finish” their SFHs before their less dense (i.e., diskier) counterparts in lower-density regions (Section 4.1.2). This scenario seems naturally linked to the core Λ CDM concept of “assembly bias,” and echoes the ideas of Holmberg (1965), Tinsley (1968), and Dressler (1980).

We contend that all of the above makes G13 a descriptively powerful and physically informative paradigm of galaxy evolution that is meaningfully different from the dominant “grow-and-quench” interpretation. However, we acknowledge that models set in the latter framework—and at least one other based on *opposite* assumptions to our own (Kelson 2014)—can reasonably claim to match some or all of the observations examined here as well as G13 (Section 4.2). As such, independent of G13’s ultimate accuracy, we also contend that an important lack of consensus exists regarding basic aspects of what a “good” description of galaxy evolution looks like. These would stymie any attempt to set physical phenomena in a broader unifying explanatory theory.

We attribute much of this ambiguity to a lack of discriminating power in current data (Section 4.3). This may be due to the requirement to project galaxies onto planes defined by spatially/temporally integrated quantities. Such issues will be remedied by the next generation of telescopes—*JWST*, *WFIRST*, and thirty-meter class ground-based facilities—which will enable us to better connect progenitors to descendants based on spatially resolved spectroscopy over

large swaths of cosmic time, contextualize those findings in diverse environments, and therefore better assess both the physics driving and the story of galaxy evolution.

L.E.A. thanks Dan Kelson, Jordan Mirocha, and Lindsay Young for intense, informative, and inspiring conversations. He also thanks Sean Johnson for software contributions; Josh Speagle for edifying feedback; Tommaso Treu for his patience and guidance during the completion of this work; and Hsiao-Wen Chen, Andrey Kravtsov, and Rich Kron for their helpful comments during its inception. Finally, he thanks Barbara Zera Abramson for valuable editorial and organizational assistance, and the anonymous referee for his/her insightful remarks. B.V. acknowledges support from an Australian Research Council Discovery Early Career Researcher Award (PD0028506).

APPENDIX A AXIOMATIC UNDERPINNINGS OF G13

A.1. Defining the G13 Class

A key aim of this paper is to explore how axioms of SFH models affect physical interpretations of data (Section 4). The vehicle for this exploration has been a model presented in Gladders et al. (2013) that is sufficiently descriptive of the data to make this exercise worthwhile (Section 3.1). Yet, it is just one member of a class of models representing instantiations of similar ideas. We are therefore interested in the potential of the G13 *model class* beyond the specific realization published in the above text. To clarify what this means, we must clarify what the G13 model class is.

Gladders et al. (2013) describes a “species” of the model family defined by the following traits:

1. The basic modeling unit is the individual SFH as assigned to each member of an input set of (real) galaxies;
2. The SFHs are continuous and smooth;
3. The SFHs rise and fall;
4. The SFHs are constrained mainly in the ensemble (e.g., by the time derivative of their *sum*).

In terms of other models mentioned in-text, (1) is unique to this family; (2) is shared by quenching-based/SFMS-driven models (up to the quenching event; e.g., Peng et al. 2010); and (4) is shared by the diversification-based stochastic model of Kelson (2014). Trait (3) *appears* in all, but is not required by K14, and SFHs that fall faster than the SFMS (see below) cannot be produced by SFMS-based approaches, driving their requirement for additional quenching mechanisms. Thus, though akin to a “basis-set” in “paradigm-space,” these perspectives are not quite orthogonal.

Trait (1) sets the basic aspect of the G13 family: a longitudinal survey of the input galaxy sample. This point is important and the subject of Appendix B. We do not know if the G13 instantiation *specifically* (see below) is a singularly good representative of this family, but we have shown that at least one “cousin”—a collection of Gaussians—is not viable (see G13).

Extrapolating from this and the results of O13, the *class* of viable models to which G13 belongs exhibits two traits beyond those listed above:

5. The SFHs are *time-asymmetric*, such that they rise faster than they fall;

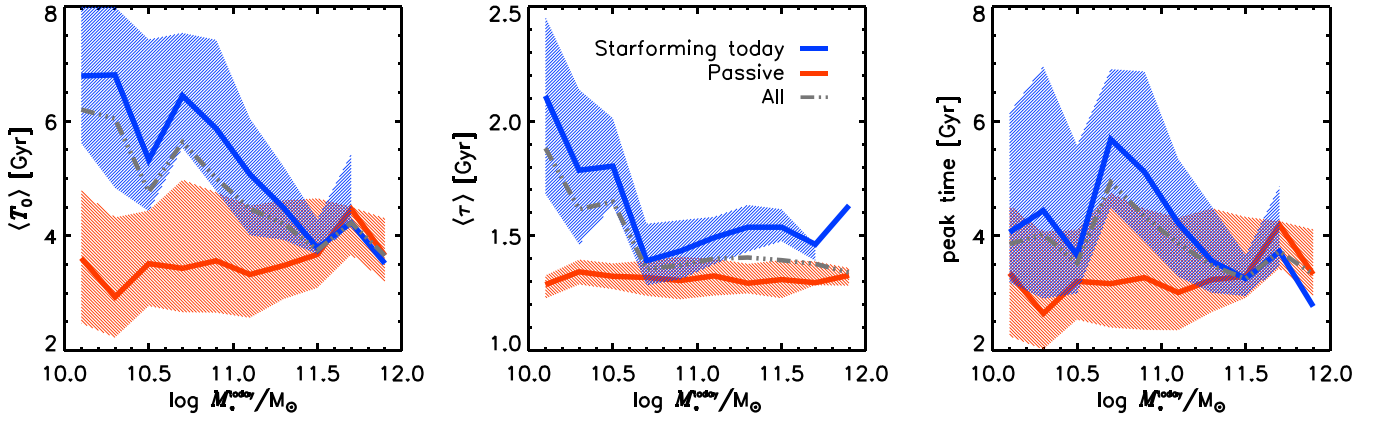


Figure 15. From left: median T_0 , τ , and $t_{\text{peak}}-M_*^{\text{obs}}$ trends as output by G13 for “star-forming” (blue) and “non-star-forming” (orange) present-day galaxies (split at $\log \text{sSFR} = -11$). Band widths denote 25%–75% parameter spreads; gray dotted–dashed lines show trends for all galaxies. These quantities show different mass and population dependencies. That $\langle T_0 \rangle$ monotonically decreases while $\langle \tau \rangle$ stays flat at $\log M_* \gtrsim 10.5$ for today’s star-forming systems demonstrates the independence of these parameters [trait (6)] needed to produce the large- T_0 , small- τ objects identified by O13. Note that the decorrelation of all quantities with M_* for the passive population reflects the same lack of sSFR constraints discussed in the context of the UVJ diagram and downsizing (Sections 3.1.5 and 3.1.6; Figures 7 and 8). The apparent bimodality in $\langle \tau \rangle$ at $\log M_* \sim 10$ may also be driven in part by this issue, but much of it may not (Appendix C), demonstrating that population bifurcation can emerge from models where SFHs are drawn from a unimodal parent distribution (Figure 17).

6. The SFHs are characterized by at least two parameters, such that their half-mass-times are formally independent of their widths (Figure 15).

The ability of this class to produce the late-peaking, narrow SFHs demanded by O13 was a result of (6). This is a consequence, however, of the more basic fact that (6) allows SFHs to fall faster than the SFMS—i.e., to exit the star-forming population without invoking any mechanism beyond whatever causes SFHs to decline, generally. As such, (6) is *also* the reason G13 can avoid explicit quenching (see discussion of Figure 9 in Section 3.2.1). Yet, because Gaussian SFHs exhibit this trait and are *not* good descriptions, this criterion alone is insufficient. The lognormal fulfills both (5) and (6), though it was not selected for this reason.

The G13 paradigm would be invalidated if the following *class-level* failures were identified. From most to least damaging:

1. Large, abrupt discontinuities are required to explain a significant fraction of SFHs for systems dominating the universe’s mass and star formation budget over a significant fraction of cosmic time (smooth, continuous parameterizations cannot describe most SFHs);
2. There is no interesting physical interpretation of (the independence of) SFH half-mass times and widths;
3. There exists no physical quantity at, e.g., $z \gg 8$ to which the parameters of smooth, continuous SFHs could be tied that correlates meaningfully with observed properties at, e.g., $z \ll 8$.

Such findings would demonstrate flaws at the core of the G13 modeling framework and cause us to abandon it.

A.2. Defining the G13 Instantiation

Though nominally central to the model—and perhaps ultimately physically meaningful—the lognormal form is really just one of a suite of mathematical details that distinguish G13 *as currently realized* as an instantiation within the class just described. Its defining structures are displayed below.

The lognormal SFH of the i th galaxy is given by:

$$\text{SFR}(t)_i \propto \frac{\exp\left[-\frac{(\ln t - T_{0,i})^2}{2\tau_i^2}\right]}{t\sqrt{2\pi\tau_i^2}}, \quad (2)$$

under the boundary conditions:

$$f \int_{t_0}^{t^{\text{obs}}} \text{SFR}(t)_i dt = M_{*,i}^{\text{obs}}, \quad (3)$$

$$\text{SFR}(t^{\text{obs}})_i = \begin{cases} \text{SFR}_i^{\text{obs}} & \text{if } \text{SFR}_i^{\text{obs}} > \text{thresh} \\ (-\infty, \text{thresh}] & \text{otherwise,} \end{cases} \quad (4)$$

and

$$\frac{1}{V} \sum_{i=1}^N \text{SFR}(t)_i = \text{SFRD}(t). \quad (5)$$

That is, the current G13 model is a set of N ($=2094$) half-mass-time/width parameter pairs, $\{(T_0, \tau)_i\}$, that ensure (1) the i th galaxy’s lognormal SFH (Equation (2)) leads to its observed $(M_*, \text{SFR})^{\text{obs}}$ at the correct epoch (3) and (4); and (2) the ensemble of SFHs sums to the correct SFRD at all times (5) (V is the comoving volume of the input sample). We assume that SFHs for galaxies with SFR^{obs} upper-limits can take any value below that threshold at t^{obs} ($0.05 M_\odot \text{ yr}^{-1}$; G13). We also adopt a constant IMF-dependent $\text{SFR} \mapsto \dot{M}_*$ conversion factor, $f \sim 0.7$ (Salpeter). Graphically, Figures 2 and 9 depict these equations.

The original treatment also entailed a merger prescription (see G13), but such events were neglected in the calculations used here. Given results from Leitner (2012), Behroozi et al. (2013b), Abramson et al. (2015), and our own numerical experiments, this appears safe, though it may impact number counts at $z \gtrsim 2.5$ and $\log M_* < 10$ (Figures 4 and 18; cf. Abramson et al. 2015, Figure 1) and likely implies that G13 SFHs represent mean histories for ensembles of merging sub-systems at similar or higher z .

Various G13 “siblings” are characterized by lognormal or double-lognormal forms for $\text{SFRD}(t)$ (Figures 1, 16), and the

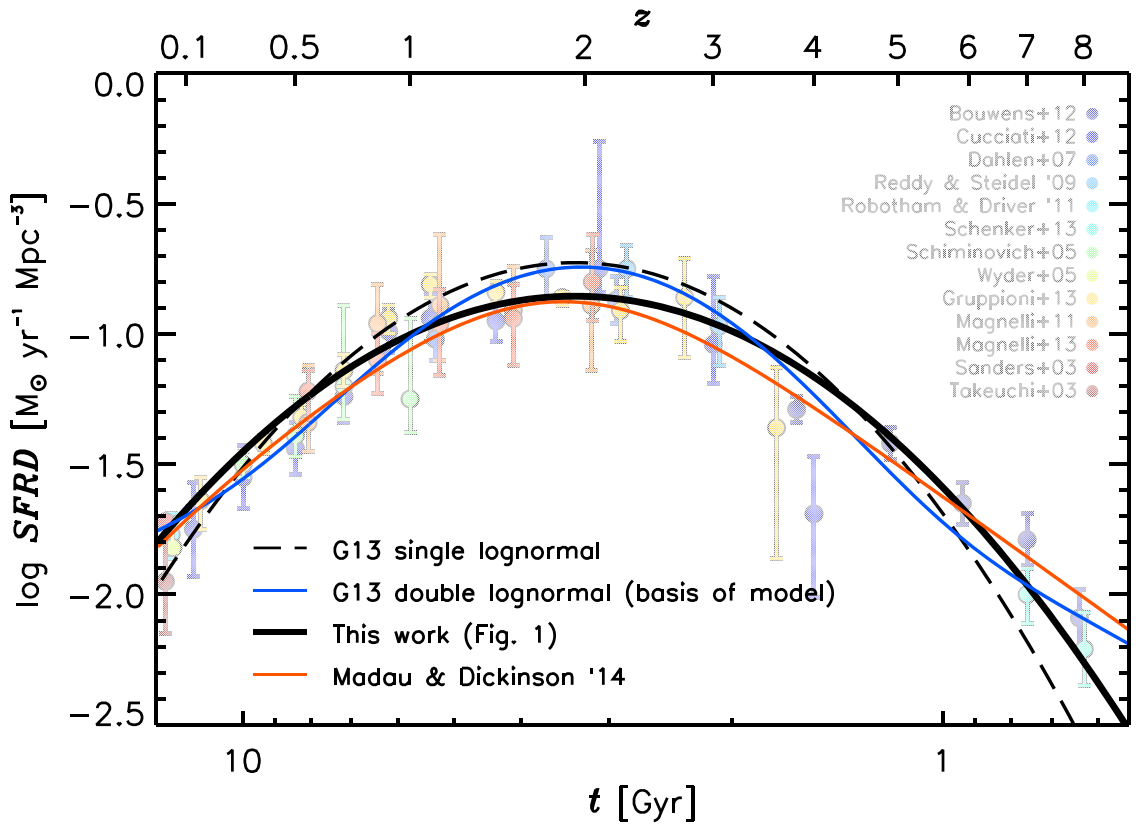


Figure 16. Parameterizations of the cosmic $SFRD(t)$. Solid black line: new best-fit single lognormal from Figure 1. Dashed line and solid blue line: the single- and double-lognormal fits from the original G13 paper, respectively. The former $[(T_0, \tau)_{\text{Uni}} = (1.54, 0.57)]$ was the inspiration for that work; the latter $[(T_{0,1}, T_{0,2}; \tau_1, \tau_2)_{\text{Uni}} = (1.39, 2.80; 0.46, 1.19)]$ constrained the SFHs analyzed here (and in Abramson et al. 2015). Orange line: the best-fit double power-law from Madau & Dickinson (2014). The fits are similarly good, though differences at the highest- z may affect, e.g., the evolution of the stellar mass function (Figures 4 and 18), and UVJ diagram (Figure 7).

use of sSFR distributions at $\langle z \rangle \simeq 0.3, 0.5, 0.7$, and 0.9 to further constrain $\{(T_0, \tau)_i\}$. The results in this paper were derived from the maximally constrained (double-lognormal $SFRD + sSFR$ distributions) realization, which is qualitatively similar to the others, though likely differs in quantitative detail. Models generated using a different $SFRD$ form and/or a different SFH parameterization may ultimately be shown to exhibit cross-sectional projections (Section 2.5, Appendix B) so closely aligned with G13’s as to be indistinguishable. Conversely, we might find that imposing additional (e.g.) sSFR constraints on the current model cannot remedy the “superficial” issues mentioned in-text (Section 3.1). In either case, we will learn something important about the nature of smooth (lognormal) SFHs, or identify regimes where they are inappropriate. We are obtaining more discriminating data to enable such tests of how the G13 SFHs relate to those of individual galaxies (Dressler et al. 2016, A. Oemler et al. 2016, in preparation).

APPENDIX B

FURTHER EPISTEMOLOGICAL IMPLICATIONS OF G13

The G13 family comprises theoretical longitudinal surveys of real galaxy samples. This fact formally limits any such model to describing *only* true (theoretical) progenitors/descendants. G13 is therefore subject to a different kind of incompleteness than that typically discussed in astronomy.

A sample is “complete” in the usual sense if it fills a representative volume of $(M_*, SFR, \dots)$ space *at the epoch from*

which it is drawn. This is “cross-sectional” completeness, and investigators go to great lengths to compile and compare such samples at multiple epochs.

G13 does not describe these data. It does not produce a time-ordered series of complete cross-sectional samples, but *cross-sectional views of a longitudinally complete model* (i.e., one containing the entire set of progenitors of a single cross-sectional sample). By comparing these two distinct entities, this work has perpetrated a subtle but significant sleight-of-hand. Because there should be objects in the data that do not represent progenitors of our input sample—and so should not be described by the model’s projections—any such comparisons may have been inappropriate.

Yet, the fact that the G13 model performs as well as it does in this (formally dubious) exercise suggests otherwise: In many contexts, none of the above seems to matter. One context where it might is the (non-)evolution of the scatter in the SFMS (σ_{SFMS}), but because we believe we can apply additional constraints to match this using the current G13 approach, even this might not be the case (up to some ultimate redshift limit; Section 3.1.1). If so, it would seem that a complete set of galaxies (in the usual, cross-sectional sense) at one epoch corresponds to a complete set of progenitors at all epochs, at least over interesting areas of parameter space (see, e.g., Figure 18).

Note: *this does not mean that cross-sectional trends are equivalent to evolutionary trends*, just that complete cross-sectional data at earlier epochs contain/represent analogs of

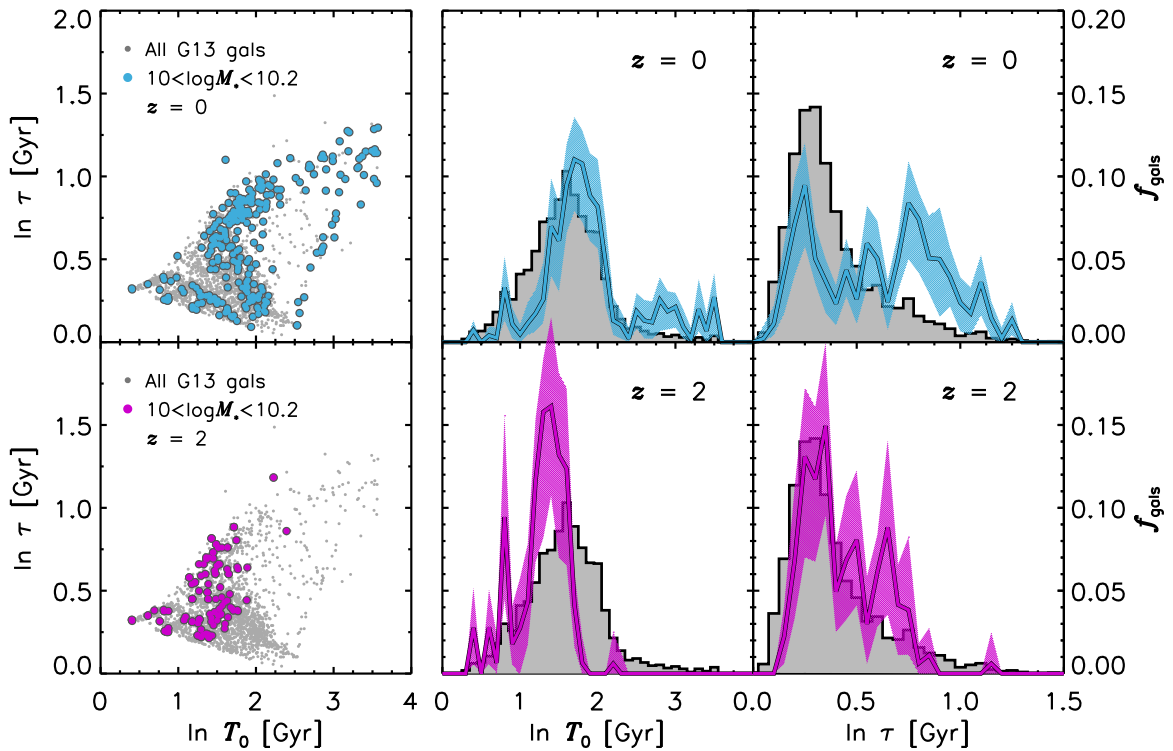


Figure 17. Top, from left: the joint (T_0, τ) , and 1D T_0 and τ distributions for all G13 SFHs (gray points/histograms; shading is 95% credibility) and those with $z = 0$ $\log M_* \sim 10$ (blue). Bottom: the same, but at $z = 2$ (purple). A τ bimodality emerges at $z = 0$ as anticipated from Figure 15, but not at $z = 2$, illustrating how similar (M_*, SFR) regimes correspond to different (T_0, τ) regimes different times. As a consequence, *contemporaneous*, equal- M_* passive and active galaxies will also lie in different parts of (T_0, τ) space: the former in regions that corresponded to “star-forming” long ago, the latter in those that do so at the time of observation. This naturally allows such classes (defined at a single epoch) to bifurcate in (T_0, τ) , though they all hail from a unimodal parent distribution.

many/all of the progenitors of complete cross-sectional samples at later epochs. Proving (under which conditions) this statement is true would be a boon to the science.

APPENDIX C FROM ONE, MANY: BIMODALITY AND THE $(T_0, \tau; t)$ CONTINUUM

Figure 15 suggests a bimodal (T_0, τ) distribution for today’s low-mass star-forming and non-star-forming galaxies. This would seem to contradict our central claim that quenching is not a “special” process, but corresponds to the declining side of *all* SFHs which are drawn from a (T_0, τ) continuum (see also Section 3.1.4). Figure 17 illustrates why this is not the case.

From left to right, this Figure shows the joint (T_0, τ) , and 1D T_0 and τ distributions for all G13 SFHs (gray), and those with $z \sim 0$ (top) and $z \sim 2$ (bottom) $10 \leq \log M_* \leq 10.2$ (colors). At either redshift, T_0 shows no bimodality, but at $z \sim 0$, there is some in τ as anticipated and reflected in the joint distribution. However, this signal vanishes at $z \sim 2$, belying a key point: Depending on the epoch, any SFH derived from a coordinate in the (T_0, τ) continuum will be above or below a given sSFR threshold and therefore appear as a star-forming or non-star-forming galaxy (cf. lines of constant sSFR in G13, Figure 9, with those in Dressler et al. 2016, Figure 10). That is, the SFRD is supported by different parts of (T_0, τ) space at different redshifts.

Comparing the purple to blue distributions in Figure 17 shows this: Galaxies of a given mass (and star-forming class) occupy different locations in (T_0, τ) space at different times.

Hence, at any given time, star-forming and non-star-forming galaxies will, by definition, lie in different parts of the plane: one is coming into dominance, the other is (far) removed from it. Indeed, at early times, some portion of the plane represents SFHs that mathematically cannot achieve a given (M_*, SFR) state and will therefore be unpopulated (lower-left panel). Conversely, at late times, much more (T_0, τ) space can lead to a given (M_*, SFR) —there is more time for slowly evolving SFHs to mature—more fully populating the diagram (upper-left panel). Hence, the age of the universe inevitably influences which portions of (T_0, τ) space dominate which portions of (M_*, SFR) space. This can naturally lead to the appearance of bimodality even though all systems were actually drawn from the same unimodal parent distribution.

A further practical consideration likely exacerbates any such natural segregation. As with all of our results (see Section 3), the details of G13s description of (T_0, τ) space depend on the data used to constrain the model. Since we used only sSFR distributions and the SFRD—and thus have no handle on *when* a galaxy quenched provided it was long-enough ago to lie below the $z \sim 0$ SFR threshold (Equation (4))—there can be significant separation between the passive (i.e., once-star-forming) population, and that with SFRs currently required to support the SFRD: SFHs with no $z \sim 0$ sSFR constraint tend toward the SFRD (Section 3.1.6, Figure 8), which peaks sufficiently early for those histories to finish while leaving plenty of time before present-day $\log M_* \sim 10$ star-forming galaxies must become important to the SFRD. The model has no mechanism to encourage it to fill-in this space, which would also contribute to the bifurcation discussed above.

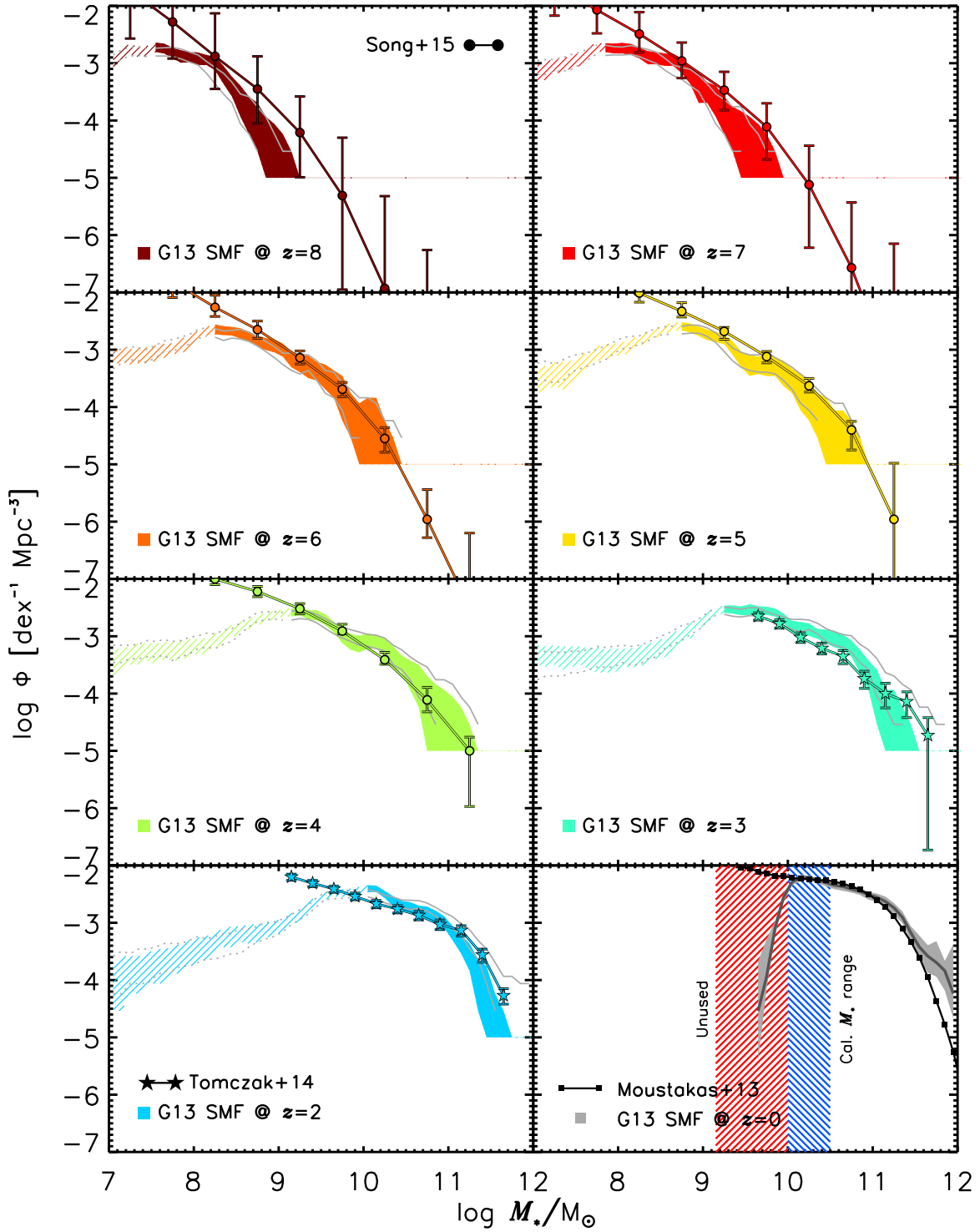


Figure 18. G13 stellar mass functions (Figure 4) broken-out by redshift to better illustrate cross-sectional incompleteness (hashed regions) and the effect of including the mild overabundance of massive galaxies in the G13 input sample compared to the full SDSS (Moustakas et al. 2013; bottom-right; tracebacks shown by solid/dotted gray lines in other panels). Colors, data points, and filled bands are those in Figure 4.

APPENDIX D

A NOTE ON THE GALAXY STELLAR MASS FUNCTION

For the sake of clarity, incomplete M_* regimes (in the usual, cross-sectional sense) were omitted in Figure 4, showing the G13 predicted stellar mass function evolution. We show the full G13 stellar mass function in Figure 18. Color-coding corresponds with the in-text figure.

Also shown is the $z \approx 0$ data (Moustakas et al. 2013, converted to a Salpeter IMF) used to calibrate G13 to an absolute scale (bottom-right). The gray curve shows the histogram of the input galaxy masses that serve as boundary conditions for the SFHs (Equation (3)). The G13 sample evidently contains an overabundance of galaxies at the highest masses compared to the full SDSS data, representative of the general population of the local universe.

To ensure tracing-back a fair galaxy sample, the G13 input data were randomly resampled ($100\times$) to match the Moustakas et al. (2013) stellar mass function before deriving results at other epochs (colored bands). If the input data are not resampled, one obtains the light gray lines in each panel. Changes are not dramatic, though tension is eased slightly in some instances and increased in others.

REFERENCES

- Abramson, L. E., Dressler, A., Gladders, M. D., et al. 2013, *ApJ*, **777**, 124
- Abramson, L. E., Gladders, M. D., Dressler, A., et al. 2015, *ApJL*, **801**, L12
- Abramson, L. E., Kelson, D. D., Dressler, A., et al. 2014, *ApJL*, **785**, L36
- Allen, J. T., Croom, S. M., Konstantopoulos, I. S., et al. 2015, *MNRAS*, **446**, 1567
- Aragon-Calvo, M. A., Neyrinck, M. C., & Silk, J. 2016, *MNRAS*, submitted (arXiv:1607.07881)
- Avila-Reese, V., Firmani, C., & Hernández, X. 1998, *ApJ*, **505**, 37
- Barro, G., Faber, S. M., Koo, D. C., et al. 2015, *ApJ*, submitted (arXiv:1509.00469)
- Barro, G., Faber, S. M., Pérez-González, P. G., et al. 2013, *ApJ*, **765**, 104
- Behroozi, P. S., Marchesini, D., Wechsler, R. H., et al. 2013a, *ApJL*, **777**, L10
- Behroozi, P. S., Wechsler, R. H., & Conroy, C. 2013b, *ApJ*, **770**, 57
- Bigiel, F., Leroy, A., Walter, F., et al. 2008, *AJ*, **136**, 2846
- Bluck, A. F. L., Mendel, J. T., Ellison, S. L., et al. 2014, *MNRAS*, **441**, 599
- Brammer, G. B., van Dokkum, P. G., Franx, M., et al. 2012, *ApJS*, **200**, 13
- Brinchmann, J., Charlot, S., White, S. D. M., et al. 2004, *MNRAS*, **351**, 1151
- Bruzual, G., & Charlot, S. 2003, *MNRAS*, **344**, 1000
- Bundy, K., Bershad, M. A., Law, D. R., et al. 2015, *ApJ*, **798**, 7
- Calvi, R., Poggianti, B. M., & Vulcani, B. 2011, *MNRAS*, **416**, 727
- Cappellari, M., Emsellem, E., Krajnović, D., et al. 2011, *MNRAS*, **413**, 813
- Chen, H.-W., Lanzetta, K. M., Webb, J. K., & Barcons, X. 2001, *ApJ*, **559**, 654
- Conroy, C., van Dokkum, P. G., & Kravtsov, A. 2015, *ApJ*, **803**, 77
- Conroy, C., Wechsler, R. H., & Kravtsov, A. V. 2006, *ApJ*, **647**, 201
- Cowie, L. L., Songaila, A., Hu, E. M., & Cohen, J. G. 1996, *AJ*, **112**, 839
- Croton, D. J., Springel, V., White, S. D. M., et al. 2006, *MNRAS*, **365**, 11
- Daddi, E., Dickinson, M., Morrison, G., et al. 2007, *ApJ*, **670**, 156
- Dalal, N., White, M., Bond, J. R., & Shirokov, A. 2008, *ApJ*, **687**, 12
- Damen, M., Labbé, I., Franx, M., et al. 2009, *ApJ*, **690**, 937
- Davies, F. B., & Furlanetto, S. R. 2016, *MNRAS*, **460**, 1328
- De Lucia, G., Weinmann, S., Poggianti, B. M., Aragón-Salamanca, A., & Zaritsky, D. 2012, *MNRAS*, **423**, 1277
- Dekel, A., & Silk, J. 1986, *ApJ*, **303**, 39
- Diemer, B., & Kravtsov, A. V. 2015, *ApJ*, **799**, 108
- Dressler, A. 1980, *ApJ*, **236**, 351
- Dressler, A., & Abramson, L. 2015, in *IAU Symp. 311, Galaxy Masses as Constraints of Formation Models*, ed. M. Cappellari & S. Courteau (Cambridge: Cambridge Univ. Press), 140
- Dressler, A., & Gunn, J. E. 1983, *ApJ*, **270**, 7
- Dressler, A., Kelson, D. D., Abramson, L. E., et al. 2016, *ApJ*, submitted (arXiv:1607.02143)
- Dressler, A., Oemler, A., Jr., Poggianti, B. M., et al. 2013, *ApJ*, **770**, 62
- Elahi, P. J., Knebe, A., Pearce, F. R., et al. 2016, *MNRAS*, **458**, 1096
- Fabricant, D., Franx, M., & van Dokkum, P. 2000, *ApJ*, **539**, 577
- Fang, J. J., Faber, S. M., Koo, D. C., & Dekel, A. 2013, *ApJ*, **776**, 63
- Feldmann, R., Hopkins, P. F., Quataert, E., Faucher-Giguère, C.-A., & Kereš, D. 2016, *MNRAS*, **458**, L14
- Feldmann, R., & Mayer, L. 2015, *MNRAS*, **446**, 1939
- Franx, M., van Dokkum, P. G., Schreiber, N. M. F., et al. 2008, *ApJ*, **688**, 770
- Furlanetto, S. R., Hernquist, L., & Zaldarriaga, M. 2004a, *MNRAS*, **354**, 695
- Furlanetto, S. R., Zaldarriaga, M., & Hernquist, L. 2004b, *ApJ*, **613**, 16
- Gadotti, D. A. 2009, *MNRAS*, **393**, 1531
- Gavazzi, G., Bonfanti, C., Sanvito, G., Boselli, A., & Scodreggio, M. 2002, *ApJ*, **576**, 135
- Geach, J. E., Smail, I., Moran, S. M., et al. 2011, *ApJL*, **730**, L19
- Genzel, R., Tacconi, L. J., Gracia-Carpio, J., et al. 2010, *MNRAS*, **407**, 2091
- Genzel, R., Tacconi, L. J., Lutz, D., et al. 2015, *ApJ*, **800**, 20
- Gladders, M. D., Oemler, A., Dressler, A., et al. 2013, *ApJ*, **770**, 64
- González, V., Bouwens, R., Illingworth, G., et al. 2014, *ApJ*, **781**, 34
- Guglielmo, V., Poggianti, B. M., Moretti, A., et al. 2015, *MNRAS*, **450**, 2749
- Gunn, J. E., & Gott, J. R., III 1972, *ApJ*, **176**, 1
- Guo, Q., White, S., Li, C., & Boylan-Kolchin, M. 2010, *MNRAS*, **404**, 1111
- Harker, J. J., Schiavon, R. P., Weiner, B. J., & Faber, S. M. 2006, *ApJL*, **647**, L103
- Hearin, A. P., & Watson, D. F. 2013, *MNRAS*, **435**, 1313
- Holmberg, E. 1965, *ArA*, **3**, 387
- Hopkins, P. F., Kereš, D., Oñorbe, J., et al. 2014, *MNRAS*, **445**, 581
- Hopkins, P. F., Torrey, P., Faucher-Giguère, C.-A., Quataert, E., & Murray, N. 2016, *MNRAS*, **458**, 816
- Huertas-Company, M., Pérez-González, P. G., Mei, S., et al. 2015, *ApJ*, **809**, 95
- Johnson, S. D., Chen, H.-W., & Mulchaey, J. S. 2015, *MNRAS*, **449**, 3263
- Jones, T., Wang, X., Schmidt, K. B., et al. 2015, *AJ*, **149**, 107
- Kawinwanichakij, L., Papovich, C., Quadri, R. F., et al. 2014, *ApJ*, **792**, 103
- Kelson, D. D. 1998, PhD thesis, Univ. California
- Kelson, D. D. 2014, *ApJ*, submitted (arXiv:1406.5191)
- Kennicutt, R. C., Jr. 1998, *ApJ*, **498**, 541
- Kereš, D., Katz, N., Davé, R., Fardal, M., & Weinberg, D. H. 2009, *MNRAS*, **396**, 2332
- Kereš, D., Katz, N., Weinberg, D. H., & Davé, R. 2005, *MNRAS*, **363**, 2
- Knebe, A., Pearce, F. R., Thomas, P. A., et al. 2015, *MNRAS*, **451**, 4029
- Lacey, C., & Cole, S. 1993, *MNRAS*, **262**, 627
- Lang, P., Wuyts, S., Somerville, R., et al. 2014, *ApJ*, **788**, 11
- Larson, R. B., Tinsley, B. M., & Caldwell, C. N. 1980, *ApJ*, **237**, 692
- Leitner, S. N. 2012, *ApJ*, **745**, 149
- Lemson, G., & Kauffmann, G. 1999, *MNRAS*, **302**, 111
- Lilly, S. J., & Carollo, C. M. 2016, *ApJ*, submitted (arXiv:1604.06459)
- Lu, Y., Mo, H. J., Lu, Z., Katz, N., & Weinberg, M. D. 2014, *MNRAS*, **443**, 1252
- Madau, P., & Dickinson, M. 2014, *ARA&A*, **52**, 415
- Magdis, G. E., Bureau, M., Stott, J. P., et al. 2016, *MNRAS*, **456**, 4533
- Marchesini, D., Muzzin, A., Stefanon, M., et al. 2014, *ApJ*, **794**, 65
- Martig, M., Bournaud, F., Teyssier, R., & Dekel, A. 2009, *ApJ*, **707**, 250
- McDermid, R. M., Alatalo, K., Blitz, L., et al. 2015, *MNRAS*, **448**, 3484
- Mirocha, J., Furlanetto, S. R., & Sun, G. 2016, *MNRAS*, submitted (arXiv:1607.00386)
- Moore, B., Lake, G., Quinn, T., & Stadel, J. 1999, *MNRAS*, **304**, 465
- Morishita, T., Abramson, L. E., Treu, T., et al. 2016, *ApJ*, submitted (arXiv:1607.00384)
- Moustakas, J., Coil, A. L., Aird, J., et al. 2013, *ApJ*, **767**, 50
- Muzzin, A., Marchesini, D., Stefanon, M., et al. 2013, *ApJ*, **777**, 18
- Neistein, E., van den Bosch, F. C., & Dekel, A. 2006, *MNRAS*, **372**, 933
- Nelson, E. J., van Dokkum, P. G., Förster Schreiber, N. M., et al. 2016, *ApJ*, **828**, 27
- Noeske, K. G., Weiner, B. J., Faber, S. M., et al. 2007, *ApJL*, **660**, L43
- Oemler, A., Jr., Dressler, A., Gladders, M. G., et al. 2013, *ApJ*, **770**, 63
- Omand, C. M. B., Balogh, M. L., & Poggianti, B. M. 2014, *MNRAS*, **440**, 843
- Oppenheimer, B. D., Davé, R., Kereš, D., et al. 2010, *MNRAS*, **406**, 2325
- Ostriker, E. C., & Shetty, R. 2011, *ApJ*, **731**, 41
- Pacifici, C., Kassin, S. A., Weiner, B., Charlot, S., & Gardner, J. P. 2013, *ApJL*, **762**, L15
- Peng, Y., Maiolino, R., & Cochrane, R. 2015, *Nature*, **521**, 192
- Peng, Y.-j., Lilly, S. J., Kovač, K., et al. 2010, *ApJ*, **721**, 193
- Poggianti, B. M., Calvi, R., Bindoni, D., et al. 2013a, *ApJ*, **762**, 77
- Poggianti, B. M., Moretti, A., Calvi, R., et al. 2013b, *ApJ*, **777**, 125
- Poggianti, B. M., Fasano, G., Omizzolo, A., et al. 2016, *AJ*, **151**, 78
- Popping, G., Caputi, K. I., Trager, S. C., et al. 2015, *MNRAS*, **454**, 2258
- Quintero, A. D., Hogg, D. W., Blanton, M. R., et al. 2004, *ApJ*, **602**, 190
- Reddy, N. A., & Steidel, C. C. 2009, *ApJ*, **692**, 778
- Rodighiero, G., Daddi, E., Baronchelli, I., et al. 2011, *ApJL*, **739**, L40
- Saito, S., Leauthaud, A., Hearin, A. P., et al. 2016, *MNRAS*, **460**, 1457
- Salim, S., Rich, R. M., Charlot, S., et al. 2007, *ApJS*, **173**, 267
- Salmon, B., Papovich, C., Finkelstein, S. L., et al. 2015, *ApJ*, **799**, 183
- Salpeter, E. E. 1955, *ApJ*, **121**, 161
- Sánchez, S. F., Kennicutt, R. C., Gil de Paz, A., et al. 2012, *A&A*, **538**, A8
- Schawinski, K., Urry, C. M., Simmons, B. D., et al. 2014, *MNRAS*, **440**, 889
- Schimminovich, D., Wyder, T. K., Martin, D. C., et al. 2007, *ApJS*, **173**, 315
- Schreiber, C., Elbaz, D., Pannella, M., et al. 2016, *A&A*, **589**, 35
- Schreiber, C., Pannella, M., Elbaz, D., et al. 2015, *A&A*, **575**, 74
- Song, M., Finkelstein, S. L., Ashby, M. L. N., et al. 2016, *ApJ*, **825**, 5
- Sparre, M., Hayward, C. C., Springel, V., et al. 2015, *MNRAS*, **447**, 3548
- Speagle, J. S., Steinhardt, C. L., Capak, P. L., & Silverman, J. D. 2014, *ApJS*, **214**, 15
- Spergel, D., Gehrels, N., Baltay, C., et al. 2015, arXiv:1503.03757
- Spergel, D., Gehrels, N., Breckinridge, J., et al. 2013, arXiv:1305.5422
- Springel, V., Di Matteo, T., & Hernquist, L. 2005, *ApJL*, **620**, L79
- Stark, D. P., Ellis, R. S., Bunker, A., et al. 2009, *ApJ*, **697**, 1493
- Stark, D. P., Ellis, R. S., Charlot, S., et al. 2016, *MNRAS*, submitted (arXiv:1606.01304)
- Stark, D. P., Schenker, M. A., Ellis, R., et al. 2013, *ApJ*, **763**, 129

- Steinhardt, C. L., & Speagle, J. S. 2014, [ApJ](#), **796**, 25
- Steinhardt, C. L., Speagle, J. S., Capak, P., et al. 2014, [ApJL](#), **791**, L25
- Stringer, M. J., Shankar, F., Novak, G. S., et al. 2014, [MNRAS](#), **441**, 1570
- Tacchella, S., Dekel, A., Carollo, C. M., et al. 2016, [MNRAS](#), **457**, 2790
- Taylor, E. N., Hopkins, A. M., Baldry, I. K., et al. 2015, [MNRAS](#), **446**, 2144
- Thom, C., Tumlinson, J., Werk, J. K., et al. 2012, [ApJL](#), **758**, L41
- Thomas, D., Maraston, C., Bender, R., & Mendes de Oliveira, C. 2005, [ApJ](#), **621**, 673
- Tinker, J. L. 2016, *MNRAS*, submitted (arXiv:1607.06099)
- Tinsley, B. M. 1968, [ApJ](#), **151**, 547
- Tinsley, B. M. 1972, *A&A*, **20**, 383
- Tomeczak, A. R., Quadri, R. F., Tran, K.-V. H., et al. 2014, [ApJ](#), **783**, 85
- Treu, T., Schmidt, K. B., Brammer, G. B., et al. 2015, [ApJ](#), **812**, 114
- Valentinuzzi, T., Fritz, J., Poggianti, B. M., et al. 2010, [ApJ](#), **712**, 226
- Voit, G. M., Bryan, G. L., O'Shea, B. W., & Donahue, M. 2015, [ApJL](#), **808**, L30
- Vulcani, B., Poggianti, B. M., Fritz, J., et al. 2015a, [ApJ](#), **798**, 52
- Vulcani, B., Treu, T., Schmidt, K. B., et al. 2015b, [ApJ](#), **814**, 161
- Watkins, L. L., Evans, N. W., & An, J. H. 2010, [MNRAS](#), **406**, 264
- Wechsler, R. H., Bullock, J. S., Primack, J. R., Kravtsov, A. V., & Dekel, A. 2002, [ApJ](#), **568**, 52
- Wechsler, R. H., Zentner, A. R., Bullock, J. S., Kravtsov, A. V., & Allgood, B. 2006, [ApJ](#), **652**, 71
- Weinmann, S. M., van den Bosch, F. C., Yang, X., & Mo, H. J. 2006, [MNRAS](#), **366**, 2
- Weisz, D. R., Dolphin, A. E., Skillman, E. D., et al. 2014, [ApJ](#), **789**, 148
- Wellons, S., Torrey, P., Ma, C.-P., et al. 2015, [MNRAS](#), **449**, 361
- Wetzel, A. R., Cohn, J. D., White, M., Holz, D. E., & Warren, M. S. 2007, [ApJ](#), **656**, 139
- Whitaker, K. E., Franx, M., Bezanson, R., et al. 2015, [ApJL](#), **811**, L12
- Whitaker, K. E., Franx, M., Leja, J., et al. 2014, [ApJ](#), **795**, 104
- Williams, R. J., Quadri, R. F., Franx, M., et al. 2010, [ApJ](#), **713**, 738
- Williams, R. J., Quadri, R. F., Franx, M., van Dokkum, P., & Labbé, I. 2009, [ApJ](#), **691**, 1879
- Wise, J. H., Demchenko, V. G., Halicek, M. T., et al. 2014, [MNRAS](#), **442**, 2560
- Wuyts, S., Förster Schreiber, N. M., van der Wel, A., et al. 2011, [ApJ](#), **742**, 96
- York, D. G., Adelman, J., Anderson, J. E., Jr., et al. 2000, [AJ](#), **120**, 1579
- Young, L. M., Bureau, M., Davis, T. A., et al. 2011, [MNRAS](#), **414**, 940
- Zabludoff, A. I., Zaritsky, D., Lin, H., et al. 1996, [ApJ](#), **466**, 104
- Zolotov, A., Dekel, A., Mandelker, N., et al. 2015, [MNRAS](#), **450**, 2327