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The Gaia spectrophotometric standard stars survey: II. Instrumental effects of six ground-based observing campaigns*

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The Gaia SpectroPhotometric Standard Stars (SPSS) survey started in 2006, was awarded almost 450 observing nights and accumulated almost 100 000 raw data frames with both photometric and spectroscopic observations. Such large observational effort requires careful, homogeneous, and automatic data reduction and quality control procedures. In this paper, we quantitatively evaluate instrumental effects that might have a significant (i.e., $\geq 1\%$) impact on the Gaia SPSS flux calibration. The measurements involve six different instruments, monitored over the eight years of observations dedicated to the Gaia flux standards campaigns: DOLORES@TNG in La Palma, EFOSC2@NTT and ROSS@REM in La Silla, CAFOS@2.2 m in Calar Alto, BFOSC@Cassini in Loiano, and LaRuca@1.5 m in San Pedro Mártir. We examine and quantitatively evaluate the following effects: CCD linearity and shutter times, calibration frames stability, lamp flexures, second order contamination, light polarization, and fringing. We present methods to correct for the relevant effects which can be applied to a wide range of observational projects at similar instruments.

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1 Introduction

Gaia¹ is a cornerstone mission of the ESA (European Space Agency) Space Program, launched in 2013 December 19. The Gaia satellite is performing an all-sky survey to obtain parallaxes and proper motions to μas precision for about 10^9 objects down to a limiting magnitude of $V \approx 20$ mag, and astrophysical parameters (T_{eff} , $\log g$, $E(B - V)$, metallicity, etc.) plus 2–30 km s⁻¹ precision – depending on spectral type – radial velocities for several millions of stars with $V < 16$ mag. Such an observational effort will have an impact on all branches of astronomy and astrophysics, from solar system objects to distant QSOs (Quasi Stellar Objects,

or Quasars), from the Galaxy to fundamental physics. The exquisite quality of Gaia data will allow a detailed reconstruction of the 6D spatial structure and velocity field of the Milky Way galaxy within ≈ 10 kpc from the Sun, providing answers to long-standing questions about the origin and evolution of our Galaxy, from a quantitative census of its stellar populations, to a detailed characterization of its substructures, to the distribution of dark matter. Gaia will also determine direct geometric distances to many kinds of distance standard candles, setting the cosmological distance scale on extremely firm bases. The Gaia scanning law will cover the whole sky repeatedly (≈ 70 times on average, over the planned 5-years Gaia lifetime, see Carrasco et al. 2007; Lindegren et al. 2008), therefore, a large number of transient alerts will be a natural byproduct (for example, many SNe will be discovered, see Altavilla et al. 2012)². For a review

* Based on data obtained with BFOSC@Cassini in Loiano, Italy; EFOSC2@NTT in La Silla, Chile; DOLORES@TNG in La Palma, Spain; CAFOS@2.2 m in Calar Alto, Spain; LaRuca@1.5 m in San Pedro Mártir, Mexico (see acknowledgements for more details).

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¹ <http://www.cosmos.esa.int/web/gaia/home>

² Indeed, Gaia started already discovering SNe, as testified by the ATel (<http://www.astronomersteletgram.org/>) telegrams published so far.

of Gaia science cases, see Perryman et al. (2001), Mignard (2005), and Prusti (2011). The challenges in the data processing of this huge data volume require a team of hundreds of scientists and engineers (organized in the DPAC, or Data Analysis and Processing Consortium, see Mignard et al. 2008).

Gaia is usually described as a self-calibrating mission, but it also needs external data to fix the zero-point of the magnitude/flux system (Pancino et al. 2012) and radial velocities (Soubiran et al. 2013), and to train the classification and parametrization algorithms (Bailer-Jones et al. 2013). Gaia photometry will come from the astrometric array of CCDs in the Gaia G-band, a filter-less band whose profile is defined by the reflectivity of the mirrors and by the sensitivity of the detectors, and from the blue (BP) and red (RP) spectro-photometers which will produce dispersed images with $20 < \lambda/\Delta\lambda < 100$ over the spectral ranges 330–680 nm and 640–1050 nm, respectively (see also Jordi et al. 2010; Prusti 2012). The derived spectro-photometry is essential not only to properly classify stars or estimate interstellar extinction, but also to compute the chromaticity correction to centroids of stars, a fundamental piece of information to achieve the planned astrometric performance (Busonero et al. 2006; Lindegren et al. 2008).

The final conversion of internally-calibrated G instrumental magnitudes and BP/RP instrumental fluxes into physical units requires an external absolute flux scale, that our team is in charge of providing (Pancino et al. 2012). The ideal Gaia spectrophotometric standard stars (SPSS) grid should comprise of the order of 100 SPSS in the range $9 \leq V \leq 15$ mag, properly distributed in the sky to be observed by Gaia as many times as possible. The mission requirement is to calibrate Gaia data with an accuracy of a few percent (1–3 %) with respect to Vega (Bohlin 2007). We use as *Pillars* the three pure hydrogen WDs (White Dwarfs) adopted by Bohlin, Colina & Finley (1995) as fundamental calibrators. Because the *Pillars* are not always visible, we are also building an intermediate grid of calibrators (Primary SPSS) that are observable all year round, with 2–4 m class telescopes.

A set of more than 100 flux tables with $\leq 1\%$ internal consistency and $\leq 1\text{--}3\%$ absolute calibration with respect to Vega are of interest for many scientific applications (see Bohlin, Gordon & Tremblay 2014, for a review of spectrophotometric methods and catalogues), including dark energy surveys based on type Ia supernovae (Sullivan et al. 2011). Our observing campaign to build a grid of SPSS, for the flux calibration of Gaia data, is among the largest of this kind to date (Pancino et al. 2012). Observations started at the end of 2006, on six different telescopes and instruments, and will presumably end in 2015. Almost 100 000 raw frames were collected and we faced the challenge of analyzing these data as uniformly, automatically, and carefully as possible. To ensure that the maximum quality could be obtained from SPSS observations, a set of careful observations (Pancino et al. 2008, 2009, 2011) and data reduc-

tion (Cocozza et al. 2013; Marinoni et al. 2012) protocols was implemented, following an initial assessment (Altavilla et al. 2011, 2014; Marinoni 2011; Marinoni et al. 2013) of all the instrumental effects that can have an impact on the flux calibration precision and accuracy. Our final goal was to make sure that all the instrumental effects that could affect the quality of the final SPSS flux tables were under control, and corrected with residuals below 1%.

This paper presents the methods and results of such instrumental effects study, and is organized as follows: Section 2 presents the shutter characterization of the employed CCD cameras; Sect. 3 presents the CCD linearity studies; Sect. 4 presents the methods and results of the calibration frames monitoring campaign; Sect. 5 presents a study of lamp flexures with DOLORES at the TNG (Telescopio Nazionale Galileo); Sect. 6 presents a study of the effect of polarization on the accuracy of flux calibrations; Sect. 7 presents a method to correct low resolution spectra for second-order contamination effects; and finally, Sect. 8 presents our summary and conclusions

2 Shutter characterization

To obtain accurate and precise photometry in our survey, we need to characterize two shutter quantities: the *shutter delay* and the *minimum exposure time*.

The shutter delay time, also called shutter dead time or shutter offset, is the difference between the requested and the effective t_{exp} (exposure time). If the shutter closes

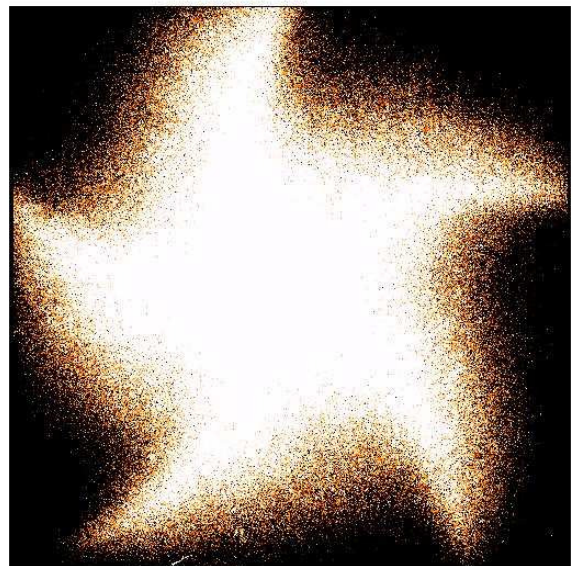


Fig. 1 An example of non-uniform CCD illumination caused by a too short t_{exp} , for an iris-type shutter, obtained with LaRuca @1.5 m in San Pedro Mártir, Mexico. The figure shows the ratio between a flat field obtained with a relatively long exposure time (20 s) and a very short one (0.1 s): ADU (Analogue-to-Digital Unit) variations of about $\pm 5\%$ are observed from the center to the corners of the image.

slower than it opens, the exposure time will be longer than requested, and the shutter delay might be positive, otherwise it is negative. Knowledge of the shutter delay time is necessary for high-precision photometry, where the effective t_{exp} is needed.

The *minimum acceptable* t_{exp} is set by the finite time the shutter takes to travel from fully closed to fully open (and vice versa). It is not strictly related to the shutter delay. A shutter can have a very small shutter delay (for example because it opens and closes with very similar delays), but it might take some time to completely open. A different shutter can have a significant shutter delay (for example because it starts closing with a significant delay), but it might have a negligible minimum acceptable time because once it starts opening (or closing) it completes the operation in a very short time. Regardless of the shutter delay, a shutter that takes a long time to fully open (or close) might introduce significant³ illumination non-uniformities across the CCD. To avoid illumination non-uniformities, t_{exp} must be longer than a safe minimum (see Fig. 1 for an example).

2.1 Observations and data reductions

To measure shutter effects, series of imaging flat fields with different exposure times were obtained at each telescope (see Table 1). Each series consisted of triplets of flats with increasing exposure time, from very short to very long (thus useful for measuring CCD linearity as well, see Section 3). After each triplet, a monitoring triplet with constant exposure time was also taken, to monitor the lamp intensity constancy⁴, following the so-called *bracketed repeat exposure method* (see Sect. 3 for more details).

All flat field images were processed with IRAF⁵ by correcting them for overscan (where applicable) and subtracting a master bias. For all analyzed instruments the dark currents turned out to be negligible (see Sect. 4) and no correction for dark was applied. The median of images belonging to the same triplet was computed to remove cosmic ray hits and to reduce the noise.

To correct for lamp drifts, each median image resulting from a monitoring triplet was divided by a reference monitoring triplet (generally the first one of the same sequence), to derive a correction factor. The correction factor was used to correct each triplet – both monitoring and data triplets – and to report them to an ideal ADU level, free of lamp drift variations. The results of this correction are illustrated

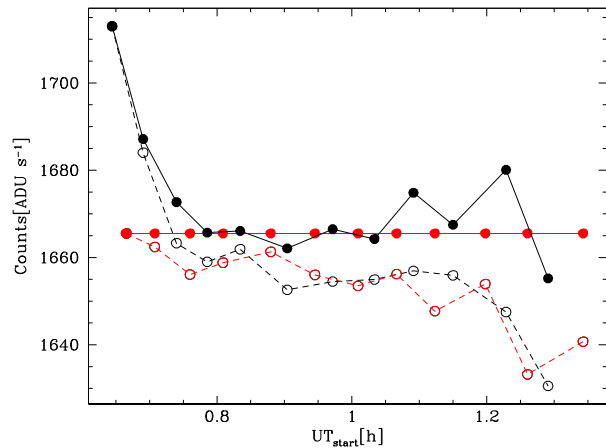


Fig. 2 Example of lamp drift correction, based on data obtained with LaRuca@1.5 m in August 2008 (see Table 1). Median ADU per second of each triplet are plotted as a function of time at triplet start. Red symbols represent monitoring triplets, while black symbols represent the actual series triplets. Empty symbols refer to uncorrected triplets, filled symbols to triplets corrected for lamp drift (see Sect. 2.1). After correction, monitoring triplets are aligned, while data triplets present some residual scatter (less than 1%), presumably caused by residual lamp instability after correction. A non-linear behaviour at low exposure times (left side of the plot) is clearly visible.

in Fig. 2. The same type of data sequences, extending to longer t_{exp} , can be used for the CCD linearity characterization described in Sect. 3, where the lamp drift correction is performed in the same way.

2.2 Shutter delay

We used two different methods to measure the shutter delay. In the first method, a linear extrapolation of the ADU (Analogue-to-Digital Units) versus t_{exp} relation crosses the time axis at a value that is generally different from zero, and that corresponds to the shutter delay (Fig. 3, top panel). In the second method, the count rate (ADU per second) are plotted versus ADU and should ideally remain constant; in practice, a deviation occurs at low t_{exp} . By iteratively adjusting the exposure time to $t + \delta t$, the shutter delay corresponds to the δt that minimizes the residuals from a flat count rate, at the low t_{exp} end (see also Fig. 3, bottom panel). We applied both methods to the data in Table 1⁶. We obtained consistent values and thus we averaged the results from the two methods to obtain a final shutter delay⁷. Results are reported in Table 2.

³ Owing to our requirements for the flux calibration of Gaia data (van Leeuwen & Richards 2012), we cannot accept illumination non-uniformities above 1% approximately.

⁴ The most common cause of lamp instability is thermal drift, causing the intensity to increase for a given exposure time as the lamp gets warmer.

⁵ IRAF is the Image Reduction and Analysis Facility, a general purpose software system for the reduction and analysis of astronomical data. IRAF is written and supported by the IRAF programming group at the National Optical Astronomy Observatories (NOAO) in Tucson, Arizona. NOAO is operated by the Association of Universities for Research in Astronomy (AURA), Inc. under cooperative agreement with the National Science Foundation.

⁶ We note that the BFOSC shutter was changed in 2010 February, with a much faster one providing a negligible shutter delay, but vignetting with t_{exp} lower than 5 s (R. Gualandi, 2010, priv. commun.). Similarly, the old diaphragm shutter of CAFOS@2.2 m was replaced by a more efficient 2-blades shutter during Summer 2008 (Santos Pedraz 2010, priv. commun.).

⁷ In practice, we applied the linearity correction (see Sect. 3) to the shutter data and the shutter delay correction to the linearity data, before re-computing both the linearity fit and the shutter delay. As expected, such corrections had a negligible impact on the resulting values.

Table 1 Flat field sequences used for the shutter delay, minimum acceptable exposure time, and linearity measurements. The columns contain (see text for more information): Column 1 the used instrument@telescope combination; Col. 2 the sequence of exposure times used; Col. 3 the exposure time of the monitoring exposures, when available; Col. 4 the used filter or grism; Col. 5 the used slit width in arcseconds, when applicable; Col. 6 the observation date; Col. 7 the experiment for which the measurements were used, with ‘a’ indicating the shutter delay measurement, ‘b’ the minimum acceptable exposure time measurement, and ‘c’ the linearity measurement. Of type a and c observations, the ones taken on 2009 January 29 lack monitoring triplets, because the data were downloaded from the TNG archive.

Instrument	t_{exp} (s)	$t_{\text{exp}}^{\text{m}}$ (s)	Filter/Grism	Slit (")	Date (yyyy month dd)	Notes
BFOSC@Cassini	0.5, 1, 2, 3, 4, 7, 10, 20, 30, 40, 50, 55, 60, 65	5	B	–	2009 May 27	a,b
	6, 26	10	gr7	2	2010 Aug 30	c
EFOSC2@NTT	10, 60, 120	40	gr3	2	2008 Nov 26	c
	3, 15, 30	10	gr3	2	2008 Nov 27	c
DOLORES@TNG	0.1, 0.2, 0.5, 1, 2, 3, 6, 15, 20	10	B	–	2008 Nov 28	a,b
	5, 20, 55	25	VHR-V	2	2008 Jan 17	c
CAFOS@2.2 m	1, 2, 4, 5, 6, 7, 8, 9	–	B	–	2008 Jan 29	a,b
	1, 5, 10, 20, 25, 30	15	B200	2	2007 Apr 01	c
LaRuca@1.5 m	0.1, 0.5, 1, 2, 3, 4, 5, 6, 20, 30, 40	–	R	–	2008 Apr 17	b
	0.1, 0.2, 0.5, 1, 2, 3, 4, 5, 6, 15, 20, 30	–	R	–	2008 Sep 09	b
	0.1, 0.2, 0.4, 0.6, 0.8, 1, 4, 6, 8, 8.7	2	V	–	2010 Sep 23	a,b,c
	0.1, 0.2, 0.4, 0.6, 0.8, 1, 2, 3, 4, 10, 20	–	White	–	2008 Aug 20	b
	1, 2, 4, 8, 15, 18, 20, 25, 30, 35, 36, 37	10	B	–	2008 Aug 23	a,c
	2, 4, 7, 14, 28, 40, 56, 62, 68, 74, 80, 87, 89	20	R	–	2010 Jul 17	a,c
	0.2, 0.4, 0.6, 0.8, 1, 2, 3, 4, 5, 6, 10, 15, 20, 30	–	R	–	2010 Jul 17	b

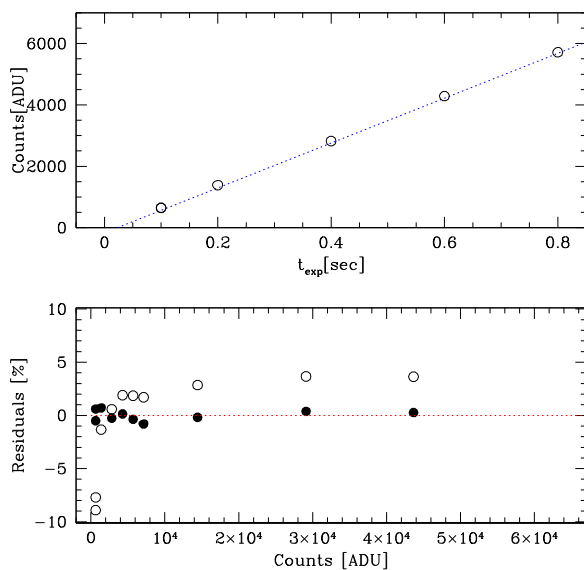


Fig. 3 Example of shutter delay determination, for Calar Alto observations obtained on 2010 September 23. *Top panel*: a zoom into the low t_{exp} region of the classical ADU versus t_{exp} plot; the dotted line represents a linear fit of all data and it does not intersect the X axis at zero, but at -0.023 s. *Bottom panel*: residuals of the count rate (ADU s^{-1}) from their average value, as a function of ADU; empty dots are the measured values, while filled dots are corrected for a shutter delay of -0.011 s. The average of the two determinations, -0.017 s, is reported in Table 2.

2.3 Minimum acceptable exposure time

The effect of illumination variations was estimated by dividing each of the flat fields in a series (Table 1) by the longest non-saturated flat field in the same series. The resulting ratio images (an example is shown in Fig. 1) should have oscillations or variations which are caused by noise only, and no systematic large scale patterns with amplitudes higher than roughly 1%, as required by the Gaia mission calibration goals. The minimum acceptable exposure times, $t_{\text{exp}}^{\text{min}}$, obtained with this criterion are listed in Table 2. While the shutter delay is generally much lower than one second, $t_{\text{exp}}^{\text{min}}$ can be higher than that. We stress the fact that if flat fields acquired with too short t_{exp} are used to correct all frames of one night, they will affect also scientific images taken with relatively long t_{exp} , and should thus be avoided.

3 CCD linearity

Linearity is a measure of how consistently the CCD responds to different light intensity over its dynamic range. CCDs can exhibit non-linearities, typically at either or both low and high signal levels (Djorgovski & Dickinson 1989; Walker 1993). High quality CCDs show significant response deviations ($\geq 1\%$) from linearity only close to saturation, i.e., at high signal levels when the potential well depth is almost full. Of course the CCD response strongly deviates at saturation, when the potential well depth is full and additional incoming photons do not increase the photoelectrons in a given pixel. When observations are restricted to the linear portion of the dynamic range, the CCD performs as a

Table 2 Resulting shutter delays and minimum exposure times for the relevant flat field series in Table 1. The columns contain (see text for more information): Column 1 the used instrument@telescope combination; Col. 2 the shutter delay; and Col. 3 the minimum acceptable exposure time.

Instrument	δt (s)	$t_{\text{exp}}^{\text{min}}$ (s)
BFOSC@Cassini ^a	-0.300 ± 0.050	≈ 5
BFOSC@Cassini ^b	negligible	≈ 5
EFOSC2@NTT	$+0.008 \pm 0.001$	< 0.1 (if any)
DOLORES@TNG	-0.011 ± 0.002	< 1 (if any)
CAFOS@2.2 m ^c	(no data)	3
CAFOS@2.2 m ^d	-0.017 ± 0.070	< 0.5
LaRoca@1.5 m	$+0.028 \pm 0.004$	5

^a Valid before 2010 February (see text).

^b Valid after 2010 February (see text).

^c Valid before Summer 2008 (see text).

^d Valid after Summer 2008 (see text).

detector suitable for accurate spectrophotometric measurements.

A few different methods have been presented in the literature to measure and correct for CCD linearity losses, as the *bracketed method* described by Gilliland et al. (1993), the *bracketed repeat-exposure method*, and the *ratio method* described by Baldry et al. (1999). Another method related to the *ratio method* is described by Leach et al. (1980) and Leach (1987): the *variance method*. In our Gaia SPSS observing campaigns, we tested the CCD linearity with two methods, described below.

The *classical method* is based on a series of imaging flat fields with increasing t_{exp} . Usually, at very high count levels the 1:1 relationship between ADU and t_{exp} breaks, and an increment of t_{exp} does not correspond anymore to a fixed increase in ADU (an example can be found in Fig. 4). We used the relevant data listed in Table 1, reduced and corrected for lamp drift as described in Sect. 2.1. We considered that a deviation of $\geq 1\%$ from linearity was not acceptable given our requirements, and thus we report in Table 3 the ADU level at which such deviations occur.

The second method we tested is described by Stello et al. (2006) and is a variation on the literature methods described above. It requires the acquisition of at least two spectroscopic flat fields, one reaching the saturation level and the other covering a fainter intensity range. The choice of slit width and grism (or grating) is not crucial, but it is important to obtain flat fields with a wide – and preferably monotonic increasing – intensity range along one direction (either a CCD column or line). The frames were collapsed along the dispersion direction, and corrected for lamp drifts as explained in Sect. 2.1, but taking into account that lamp drifts in spectra can show also colour variations. Afterwards, we applied the method as described in the original paper by Stello et al. (2006), by computing the intensity ratio of exposure pairs, normalized by their respective t_{exp} ,

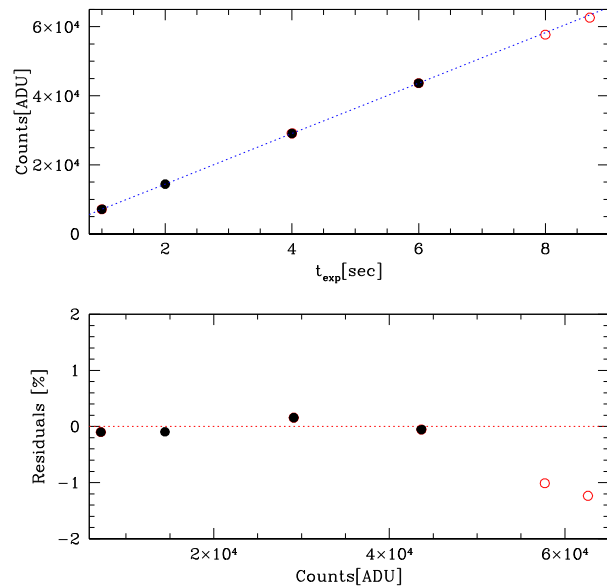


Fig. 4 Example of linearity test for Calar Alto observations obtained on 2010 September 23, with the SITE#1d_15 CCD. *Top panel*: a zoom into the high ADUs region of the classical ADU vs. t_{exp} plot; the dotted line represents a linear fit of all data. *Bottom panel*: residuals of the count rate (ADU s^{-1}) from their average value, as a function of counts; the dotted line represents a constant count rate. In both panels, points deviating by more than 1% are represented by empty circles and coloured in red.

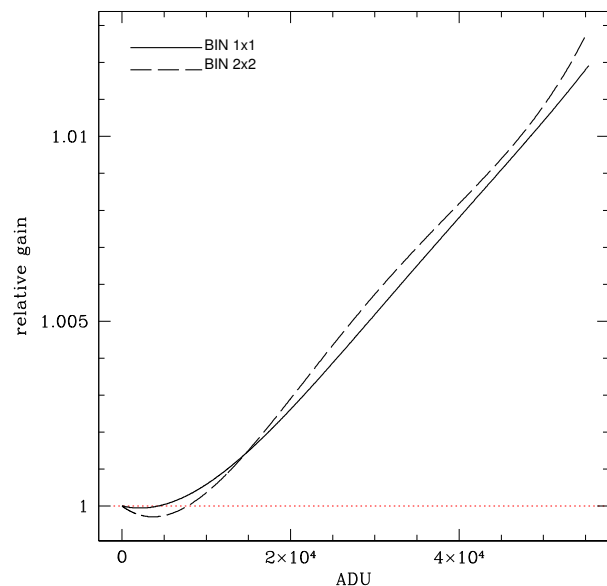


Fig. 5 Final gain curves obtained with the Stello et al. (2006) method for NTT, using data obtained on the 26 and 2008 November 27, with 1x1 and 2x2 binning, respectively, with the ESO grism #3. The dashed line is for the 2x2 binning while the solid line is for the 1x1 binning. Very small deviations from linearity start to appear already around 5000 ADU, while deviations above 1% kick in at $\approx 47\,500$ ADU for the 2x2 binning and at 49 000 ADU for the 1x1 binning.

Table 3 Resulting maximum acceptable ADU levels obtained from the flat series in Table 1. The columns contain (see text for more information): Column 1 the used instrument@telescope combination; Col. 2 the CCD used; Col. 3 the validity period of each CCD; Col. 4 the maximum acceptable ADU level obtained with the classical method; Col. 5 the maximum acceptable ADU level obtained with the Stello et al. (2006) method when applicable; and Col. 6 the final adopted value for each CCD: when there were no sufficient data, the safe value was assumed conservatively on the basis of that obtained for similar CCDs; in those cases, the values are indicated in parentheses.

Instrument	CCD	Validity Period	ADU _{max} ^(Classic)	ADU _{max} ^(Stello)	ADU _{max} ^(Adopted)
BFOSC@Cassini	EEV 1300×1340B (old)	Before 2008 Jul 17	...	...	(50 000)
	EEV 1330×1340B (new)	After 2008 Jul 17	...	60 000	60 000
EFOSC2@NTT (1×1 bin)	CCD#40 LORAL/LESSER	All runs (from 2007)	...	49 000	49 000
EFOSC2@NTT (2×2 bin)		All runs (from 2007)	...	47 000	47 000
DOLORES@TNG	E2V 4240 (Marconi)	All runs (from 2007)	60 000	60 000	60 000
CAFOS@2.2 m	SITE#1d.15	All runs (from 2007)	55 000	53 000	53 000
LaRuca@1.5 m	SITE#1	Until 2009 Jul	60 000	...	60 000
	ESOPO	2009 Oct 20–22	...	...	(60 000)
	E2V 4240 (Marconi) #1	2009 Oct–2010 Dec.	60 000	...	60 000
	E2V 4240 (Marconi) #2	2011 Mar 9–17	...	...	(57 000)
	SITE#4	2011 May 2–11	...	...	(60 000)

Table 4 List of grisms used for the SPSS campaigns with the four spectrographs. The columns contain: Column 1 the instrument@telescope combination; Col. 2 the grism name; Col. 3 the minimum wavelength; Col. 4 the maximum wavelength; Col. 5 the approximate wavelength at which second order contamination starts (if any, see text); Col. 6 the approximate wavelength at which fringing starts (if any, see text); and Col. 7 the approximate relative amplitude of the fringing pattern, as measured on our spectra before correcting for fringing. All figures were derived from our SPSS spectra.

Instrument	Grism	$\lambda_{\min}$ (Å)	$\lambda_{\max}$ (Å)	$\lambda_{2\text{ndorder}}^{(\text{start})}$ (Å)	$\lambda_{\text{fringing}}^{(\text{start})}$ (Å)	Max. Fringing (±%)
BFOSC@Cassini	#3	3300	6420	—	—	—
	#5	4800	9800	—	≈ 7000	10–15
CAFOS@2.2 m	B200	3200	9000	—	—	—
	R200	6300	11000	> 12000	≈ 8500	≤ 5
EFOSC2@NTT	#5	5200	9350	—	≈ 7300	≤ 10
	#11	3380	7520	—	—	—
	#16	6015	10320	≈ 6000	≈ 7300	≤ 10
DOLORES@TNG	LR-B	3000	8430	≈ 6000	—	—
	LR-R	4470	10073	≈ 9500	≈ 8000	5–15

which is called *gain-ratio curve*. This ratio is the starting point of an iterative procedure of inversion to determine the gain at different intensity levels.

In conclusion, given our simple purpose of estimating a safe ADU limit for our observations, the classical method and the high-precision method (Stello et al. 2006) give substantially the same results. While the required spectroscopic observations are much faster to obtain, the Stello et al. (2006) method is highly sensitive to different choices for the fit of the gain-ratio curves (function, order, and rejection method), and to their extrapolation at zero counts. Therefore, the method is not as straightforward to apply as the classical method. To estimate the ADU level at which our data deviate more than 1% from linearity, we conservatively chose the lowest of the two estimates, in the few cases where the two methods provided different results. Results for the employed instruments can be found in Table 3.

4 Calibration frames

The quality and stability of the calibration frames was tested on all the acquired calibration data, producing recommendations for the calibration plans that were progressively refined as more and more data were available for the analysis. We found a surprising variety of behaviours among different instruments and observing sites. Considering that we acquired of the order of 100 000 frames in our campaigns (see Pancino et al. 2012, for the SPSS campaigns description), it was important to implement automated procedures to identify those calibration frames that needed a closer inspection, before validating the corresponding reduced data. Whenever possible, we acquired daily or nightly calibration frames, but in case of problems, we used the derived calibration plan presented in Table 5 to decide how to analyze the data (i.e., which of the archival or available frames were

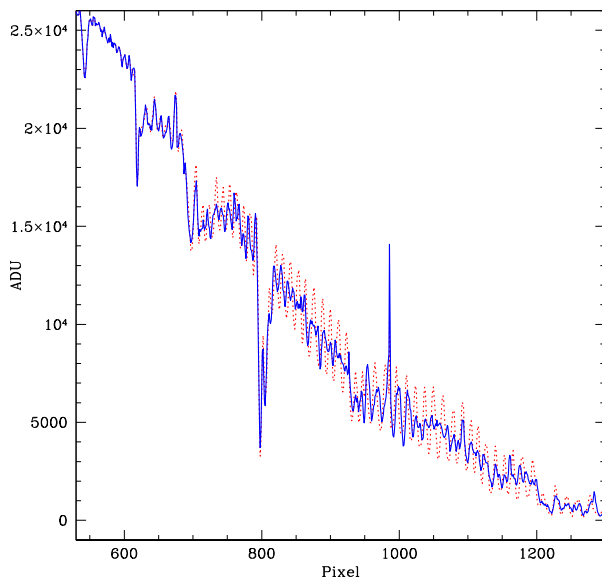


Fig. 6 Example of fringing correction (see text) on a spectrum of BD+284211, observed in Loiano on 2013 October 28, with gr5 and a 2'5 slit. The blue solid line is the spectrum after fringing correction, based on a day-time flat, while the red dotted spectrum is before the correction. Beatings are observed, caused by small differences in the fringing pattern of the spectrum and the flat used for the correction. For example, around pixel 900 the residuals are below 5%, while around pixel 1000 they remain higher, around 15%.

acceptable) and how to judge the quality of the applied reductions. In the following sections we briefly describe the monitoring strategy and the main conclusions on each instrument and telescope.

4.1 CCD cosmetics

We employed BPMs (Bad Pixel Masks), i.e., images of the same dimension of the CCD frames, where good pixels are identified by zero values and bad pixels by non-zero values. BPMs were created once or twice a year, using the ratio of two flat fields with different count levels – where pixels responding non-linearly stand clearly out – by means of the IRAF *ccdmask* task. They were applied with the IRAF task *fixpix*, which replaces bad pixels with estimates of their correct value. The estimates are built using a linear interpolation across the nearest non-bad pixels, either across lines or columns, or both, depending on the local characteristics of the BPM (see task description for more details), i.e., differently for bad lines, bad columns, or clusters of bad pixels.

4.2 Fringing

Fringing, caused by multiple reflections and interferences of red or infrared light in thin CCD substrates (Howell 2012;

Lesser 1990), was not relevant⁸ for the imaging filters used in our observations (mostly *B*, *V*, and *R*)⁹, but was present and relevant in all our spectra, because we needed to cover the whole Gaia wavelength range, from 330 to 1050 nm. All the instruments used for spectroscopy do suffer from fringing (see Table 4), well visible in our spectra (see Fig. 6). For NTT we took flat fields very close to our red grism observations, while for all other telescopes we only relied upon day-time lamp flats. However, while flat fielding does reduce the intensity of fringes, the best fringing reduction strategy for spectroscopy, i.e., observing triplets of spectra with the star in different positions along the slit, could not be applied at the chosen telescopes because it would have been too time consuming, observationally.

The solution we adopted was to apply specific spectra processing steps to reduce the impact of fringing, adapting the procedure used in the *mkfringe* task¹⁰ STIS (Space Telescope Imaging Spectrograph) for IRAF (Malumuth et al. 2003). Briefly, the most appropriate flat field available for each scientific spectrum was extracted from a region covering exactly the same pixels covered by the scientific spectrum, collapsed to a 1D spectrum along the dispersion direction, and normalized to one. The region that does not contain a significant fringing pattern was then flattened exactly to 1, to avoid adding extra noise to the scientific spectra. Finally, this specially prepared 1D spectroscopic flat was aligned to the scientific spectra and scaled until the residuals of the fringing pattern were minimized, using the IRAF *telluric* task, and applied to scientific spectra to correct them for fringing.

We found that even using spectroscopic flat fields obtained in the day-time (thus with different incident light patterns), the procedure greatly helps in reducing the amplitude of fringes, up to a factor of 2–3 in relative intensity (see Fig. 6). Sometimes, the closeness in position of the day-time flats fringe patterns to the SPSS fringe pattern produced beatings in the residuals of the fringing correction, with regions where the pattern was almost completely erased and regions where a residual fringing pattern remained. This has to be taken into account for applications of the method to scientific cases that could be sensitive to this kind of residual effects.

Additionally, in our spectroscopic campaign we often have observations of the same SPSS from different telescope and observing nights, and thus spectroscopic fringing residuals can be further reduced by building a median of the different spectra. The very red end of spectra ($\lambda \gtrsim 9500$ Å), where fringing is strong and S/N is low, will also be replaced with model or semi-empirical template spectra, as done for example by Bohlin (2007).

⁸ As in the above cases, fringing is considered to be relevant when the fringing pattern in images or spectra has an amplitude larger than 1% of the signal.

⁹ The procedure for correcting fringing in images involves the construction of a “super-flat” and is well described elsewhere (Gullixson 1992; Newberry 1991; Tyson 1990).

¹⁰ <http://stdas.stsci.edu/cgi-bin/gethelp.cgi?mkfringe>

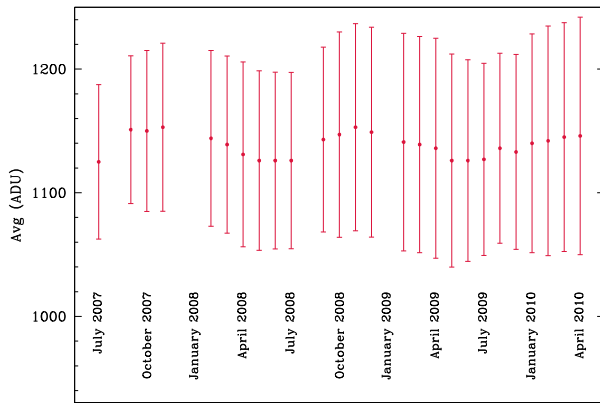


Fig. 7 Example of the long-term behaviour of the average counts (with their sigma) of REM monthly master dark frames with $t_{\text{exp}} = 60$ s, covering roughly 3 years. A seasonal trend is clearly visible.

4.3 Dark frames

Dark correction is the subtraction of the electron counts which accumulate in each pixel due to thermal noise. The reduction of dark currents is the main reason why all astronomical CCDs are cooled to liquid nitrogen temperatures. We took dark exposures as long as the longest scientific exposure, at least once per year. As expected, all used instruments showed at most a few ADU per pixel per hour, except for the REM (Rapid Eye Movement) robotic telescope in La Silla, which is cooled by a Peltier system. The REM staff creates monthly master dark frames, containing the bias structure as well (see below), thus we applied the closest master dark frame in time, taken with the correct t_{exp} . A seasonal effect on the dark frames average level can be clearly seen in Fig. 7, while the 2D shape of the frames was extremely stable and the ADU increased linearly with t_{exp} .

4.4 Overscan and bias frames

The bias current is an offset, preset electronically, to ensure that the Analogue-to-Digital Converter (ADC) always receives a positive value and operates in a linear regime, as much as possible. The offset for each exposure given by the bias level has to be subtracted before further reduction, and may be modelled as $A + B(x_i, y_i)$, where B is the pixel-to-pixel variation or 2D structure, which is time-invariant unless particular problems occur, while A is the overall bias level, that can change slightly during the night or even seasonally, owing to temperature changes.

Generally, A can be measured using a strip of pixels, called “overscan”, acquired by continuing to readout the CCD beyond its real physical extent: the result is an oversized array with a strip of signal-free pixels. In our survey, the overscan was available only for DOLORES@TNG and LaRuca@1.5 m in San Pedro Mártir, thus in all other cases

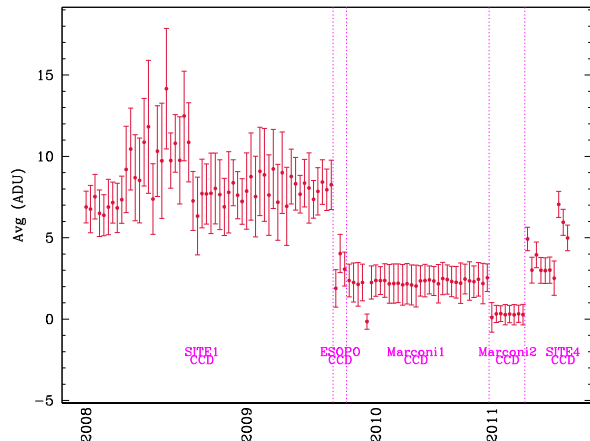


Fig. 8 Example of long-term behaviour of the average counts (with their sigma) of San Pedro Mártir daily master bias frames, covering more than three years. The CCD changes are marked and correspond to clear jumps in ADU and also in variations of the ADU stability with time.

we used the 2D bias frames to implicitly correct for A as well. This produces some additional uncertainty if A varies during the night¹¹, because the bias calibration frames were generally acquired in the afternoon (and also in the morning for the less stable instruments, like LaRuca). The different stability of the four CCDs used with LaRuca is illustrated in Fig. 8, where the bias level – after overscan subtraction¹² – of all the master bias frames is plotted.

4.5 Flat fields

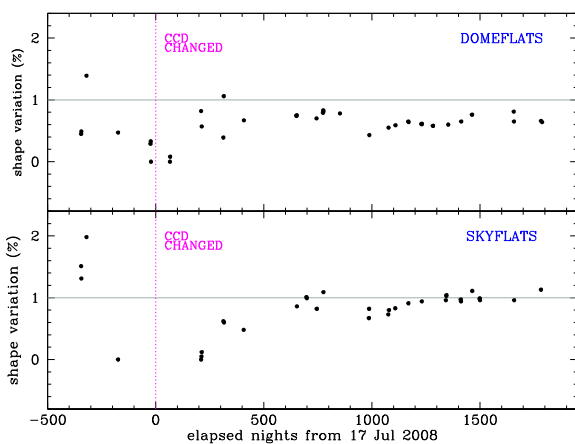
After bias subtraction and – if needed – dark correction, data values are directly related to the number of photons detected in each CCD pixel. But different pixels can be characterized by a different sensitivity: The flat fielding correction accounts for the non-uniform CCD response to incident light that can show variations on all scales, from the whole CCD to the single pixel. In imaging, flat field frames are acquired using the twilight sky or a screen uniformly illuminated by a lamp, with each filter; they are used to correct for sensitivity variations on all scales, acting as a sort of *sensitivity map*. Massey & Hanson (2013) recommend the use of dome or screen flat fields to correct for the small-scale sensitivity variations and of sky flat fields for large-scale variations. In spectroscopy, only the small scale variations are corrected with flat fields obtained through the spectrograph, i.e., with a slit and a grism (or grating). Of course, flat field frames with low signal or saturated were always rejected (see Sect. 3).

¹¹ However, we found that in both the photometry and long-slit spectroscopy cases, the standard sky subtraction procedures we adopted were sufficiently accurate to get rid of any residual bias offset, at least within the required 1% level.

¹² It can happen that after overscan correction the average level of a bias frame remains above zero by a few ADUs, because of the way the CCD is read out.

Table 5 Calibration plan recommendations following our calibration frames monitoring (see text). For each type of calibration frame, the minimum acceptable stability timescale is reported.

Instrument	Dark	Photometric Bias	Spectroscopic Bias	Photometric Dome Flat	Photometric Sky Flat	Spectroscopic Lamp flat	Spectroscopic Sky flat
BFOSC@Cassini	—	4 days	4 days	300 days	varies ^c	1 day	1 week
EFOSC2@NTT	—	5 days	1 day	250 days	250 days	1 day	1 week
DOLORES@TNG	—	1 day	1 day	—	250 days	5 days	1 week
CAFOS@2.2 m	—	1 week	1 week	varies ^b	1 day	1 day	1 week
LaRoca@1.5 m	—	varies ^a	—	—	5 days	—	—
ROSS@REM ^d	1 month	—	—	—	—	—	—

^a Varies from 1 to 7 days, depending on the CCD used.^b Was 1 day until Summer 2008, and 1 week afterwards.^c Was 4 days until 2010 February, and 100 days afterwards.^d REM flat fields were not applied, to avoid reinforcing the strong light-concentration effect from which ROSS suffers.**Fig. 9** Example of long-term behaviour of the normalized shape variation ($\Delta S/S$, see text) of BFOSC nightly dome (*top panel*) and sky (*bottom panel*) flat fields in the *B* filter. The shape variations are rarely above 1%, globally, showing an impressive stability over a period covering more than 2000 nights. The CCD change in Summer 2008 is marked by a vertical dotted line.

For imaging flat fields, we were interested in monitoring the large scale shape variations with time. We thus took all the masterframes obtained for each instrument, normalized them by their mode, and smoothed them with a boxcar to remove small scale variations. We then computed a stability function K_{stab} , using the pixel-to-pixel difference of the counts in each masterflat (F_{pix}) with respect to a reference masterflat ($F_{\text{pix}}^{\text{ref}}$):

$$K_{\text{stab}} = \sum_{\text{pix}=1}^n (F_{\text{pix}} - F_{\text{pix}}^{\text{ref}})^2. \quad (1)$$

We can then define a *normalized shape variation* as

$$\frac{\Delta S}{S} = \sqrt{\frac{K_{\text{stab}}}{K_{\text{stab}}(1\%)}}, \quad (2)$$

where $K_{\text{stab}}(1\%)$ is computed in the case where the pixel-to-pixel variation is equal to 1% everywhere. Using this indi-

cator, we have a significant variation (i.e., a global variation above 1%) when $\Delta S/S > 1$.

Using the K_{stab} global indicator, we explored all the master flat fields obtained during our SPSS survey with automatic procedures (see Fig. 9 for an example), thus identifying the flat fields that needed to be visually inspected for anomalies, before applying them to the scientific images.

4.5.1 Illumination variations

For imaging, dome and sky flat fields are sufficient to correct for the CCD response variations (Massey 1997; Massey & Hanson 2013). Instrumental effects, like internal light reflections or light concentration, need instead ad hoc procedures (Koch et al. 2004; Manfroid, Selman & Jones 2001). Among the imagers used in the SPSS campaigns, only ROSS@REM showed a significant (up to 5–10%, depending on the amount of incident light) cushion-shaped illumination variation, that was worsened by the application of the flat-field correction, similarly to what happens with light concentration. Because ROSS@REM was only used for relative photometry, we decided not to apply any flat-field correction to REM images and to keep the SPSS as much as possible in the same *X* and *Y* position in the CCD during each time series. Nevertheless, relative photometry with ROSS@REM proved to be extremely difficult and only rarely met our requirement of $\sim$ millimag accuracy.

For spectroscopy, illumination variations caused by imperfections on the slit borders, vignetting, and other similar effects, are usually corrected with the help of twilight sky spectroscopic flat fields (Massey & Hanson 2013; Schönebeck et al. 2014), applied after normal (i.e., lamp) flat fielding. After some initial testing, we decided to obtain spectroscopic flats fields once per run, as a compromise between the stability of the results and the time needed to obtain the flats (see Table 5).

5 Lamp flexures

Low resolution spectrographs are usually mounted at the telescope or at the telescope derotator, and they move during observations. At each different position, the varying projection of the gravitational force leads to mechanical distortions, which can be seen on the wavelength calibration frames, where they produce linear or non-linear shifts of the lamp emission lines (Munari & Lattanzi 1992). A correct flux table must associate the right flux to the right wavelength, thus any error in wavelength has an impact on the flux calibration of a spectrum¹³.

We performed a test to evaluate the DOLORES@TNG lamp flexures on 2008 January 31, when the instrument was equipped with three separate calibration lamps (He, Ne, Ar) that could not be switched on simultaneously, and the procedure to obtain high S/N calibration lamps for each single star during night-time observations was too time consuming. We used the LR-R grism, the 2'' slit and the Ar lamp. Triplets of wavelength calibration lamp spectra were acquired at different positions of the derotator, covering a complete derotator circle, from -260° to $+100^\circ$ in steps of 10° forward, and from $+95^\circ$ up to -255° in steps of 10° backwards. The median of the three acquired lamp frames at each position was used to evaluate line shifts.

The resulting shifts of emission lines, as a function of derotator angle, are characterized by a quasi-sinusoidal trend, as shown in Fig. 10 for a selection of the examined lines. The typical size of the global oscillation was almost three pixels, with an error on the lines peak position ranging from 0.01 pixels in the central portion of the CCD to 0.06 pixels at the CCD borders. The shift amounts to almost $8 \text{ \AA}$ in the space of wavelengths and thus daytime lamps are clearly not appropriate for the wavelength calibration of spectra in the Gaia spectrophotometric campaigns. For this reason, most astronomers also take wavelength calibration lamps close to the position of each observed spectrum. Not being able to measure emission line shifts for all instruments, we also adopted the same strategy. In particular, if the night lamp observations required too much time – as was initially the case with TNG – we decided to employ day-time lamps with very high S/N, and shift them to match the night lamps taken with lower S/N in order to save time. In all other cases we used night-time lamps.

6 Polarization

The reflectivity of curved and plane mirrors may be different for different directions of linear polarization, in a wavelength dependent way (see Breckinridge & Oppenheimer

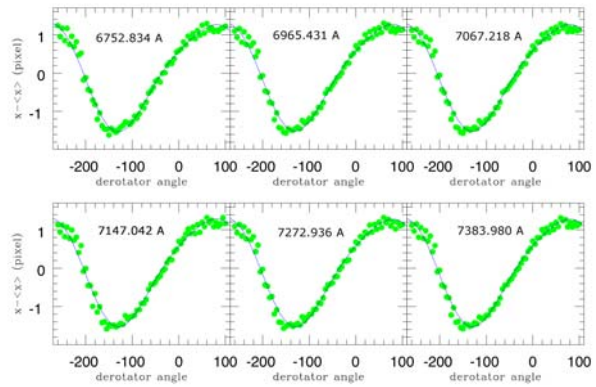


Fig. 10 Example of lamp flexure effect on emission lines position of a TNG Argon wavelength calibration lamp exposure with the LR-R grism (see text). Each panel represents the line peak shift in pixels (ordinate) of six individual argon lines as a function of DOLORES derotator angle position (abscissa). Each green dot corresponds to a peak measurement in a different lamp exposure. The typical shift of almost ± 1.5 pixels corresponds to almost $8 \text{ \AA}$. The thin black line is a polynomial fit to the data, to guide the eye.

2004, and references therein). An incoming beam of non-polarized light will be polarized according to this difference of efficiency. This might affect the sensitivity of the whole observing system, especially if the light beam encounters other polarization-sensitive optical elements after reflection. The degree of induced polarization depends on the *angle of incidence*: The maximum degree of (linear) polarization occurs for an incidence angle which depends on the material, called *polarization angle*. Refractive dispersing elements (grisms, in our case) may also – in principle – produce some degree of polarization, or may have different transmission efficiencies for different *intrinsic* polarization of the incoming light, thus introducing selective light losses. A combination of optical elements with different orientation could thus produce a significant variation in the efficiency, dependent on wavelength and/or pointing.

We can envisage three cases relevant to our SPSS campaigns:

1. The light beam encounters just one optical element potentially sensitive to polarization, for example the grism, as in the case of CAFOS@2.2 m in Calar Alto. If the incoming light is polarized, there may be a difference in the response depending on the actual orientation of the grism in the plane perpendicular to the direction of the light beam. If the incoming light is not polarized, there is no effect.
2. There is more than one element potentially sensitive to polarization, but their relative orientation in the planes perpendicular to the direction of the incoming beam are fixed, as is the case of the flat mirror feeding the grisms of BFOSC@Cassini. The net result is the same as above.
3. There is more than one element potentially sensitive to polarization, and their relative orientation in the planes perpendicular to the direction of the incoming beam

¹³ It is important to remark that – especially in those parts of the spectrum where the flux changes rapidly with wavelength – even small errors in the wavelength calibration can imply relatively large errors on the flux calibration. This problem becomes difficult to solve, i.e., by means of cross-correlation to align spectra, for featureless or nearly featureless stars, which are in general the best flux calibrators.

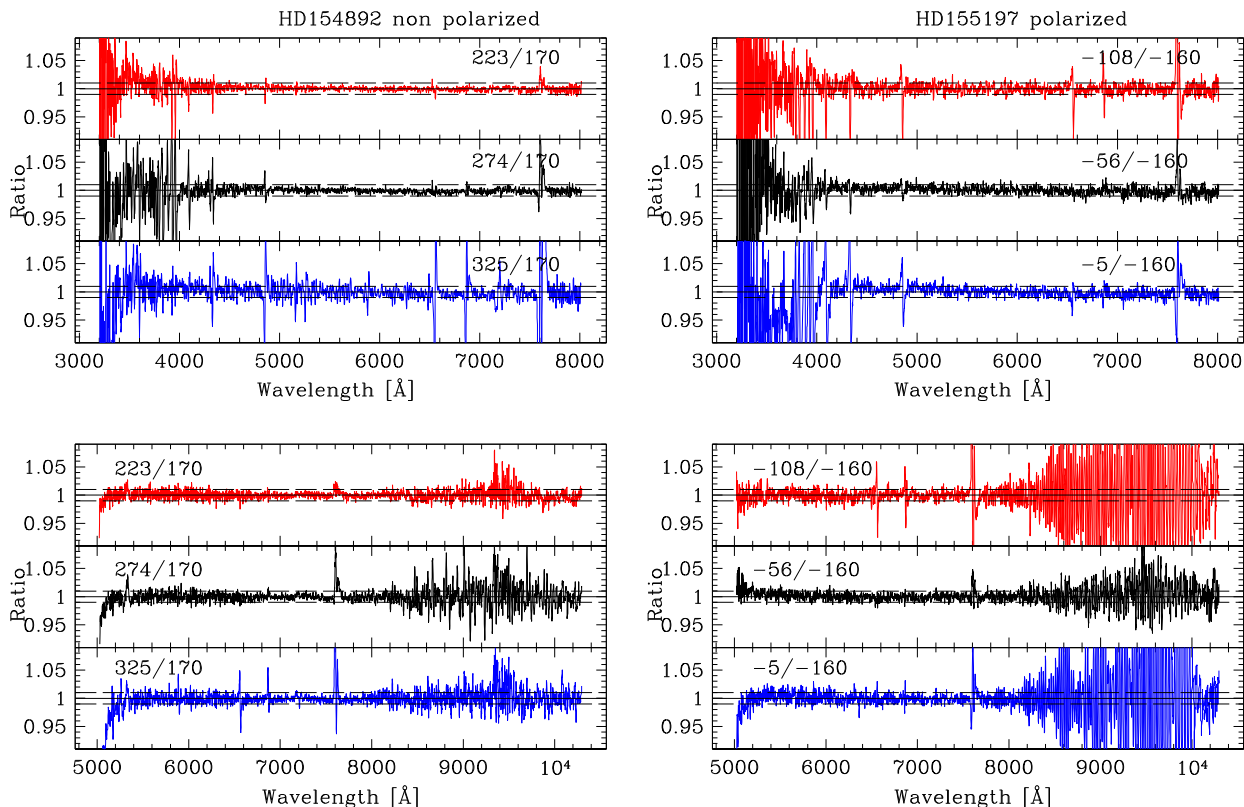


Fig. 11 Results of our test on the effect of polarization on spectrophotometry. The *left panels* report observations of the “non-polarized” star HD 154892, while the *right panels* of the “polarized” star HD 155197 (see text for more details). The three *top panels* for each star refer to the LR-B grism, while the *bottom panels* to the LR-R grism. Each sub-panel shows the ratio of spectra taken at different derotator angles with respect to a reference angle, which is 170° for the non-polarized star and -160° for the polarized star. All ratios are shown as a function of wavelength. Dashed lines in each sub-panel indicate $\pm 1\%$ ratio variations. Except for the noisy blue edge, and the red parts where fringing is important, it is clear that the effect is negligible for our purposes. The increase in the fringing residuals pattern for the polarized star is caused by the fact that the polarized star was accidentally placed in different positions along the slit for different polarization angles.

changes, depending on the telescope pointing. This is the case of the flat M3 mirror of the twin TNG and NTT telescopes¹⁴ that drives the main light path to the Nasmyth focus where it feeds the grism of the DOLORES and EFOSC2 spectrographs, mounted on the derotator device. In this case some polarization effect can be induced even on non-polarized incoming light, and the grism may respond slightly differently – in terms of efficiency – to differently polarized light, resulting in a light loss depending on the relative orientation of the two elements (see Giro et al. 2003).

6.1 The TNG polarization experiment

While potentially highly polarized sources should be avoided as candidate SPSS (as for instance DP white dwarfs or heavily extincted stars), there is no observational strategy that can completely avoid the problem. We thus tested the

above hypothesis on TNG, because it has the same NTT design, and they are the only two telescopes in our campaign where polarization could significantly affect spectrophotometric observations.

We observed two stars from Turnshek et al. (1990): HD 155197 which exhibits a polarization in the V band of $4.38 \pm 0.03\%$ – that we will call *polarized star* hereafter – and HD 154892, having a $0.05 \pm 0.03\%$ – that we will call *non-polarized star*. They were observed on a clear night (2012 June 1, kindly obtained by the TNG staff), with good seeing ($0.8''$ – $1.5''$), using a wide slit ($10''$) and the same grisms used for the SPSS campaign (LR-R and LR-B, see Table 4). Spectra were taken at different derotator angles of -160° , -108° , -56° , and -5° for the polarized star, and 170° , 223° , 274° , and 325° for the non-polarized star. The spectra were bias and flat-field corrected, extracted, and calibrated in wavelength with the help of narrow slit ($2''$) spectra and day-time calibration lamps, using standard recipes from the IRAF *apall* and *onedspec* packages, configured exactly as done for our SPSS reductions (Cocozza et al. 2013). We finally corrected all spectra for atmospheric extinction, using

¹⁴ The EFOSC2@NTT web site (<http://www.eso.org/sci/facilities/lasilla/instruments/efosc-3p6/>) reports: “Preliminary analysis of instrument polarization shows $< 0.1\%$ at field center and $\approx 0.4\%$ at the edge (07/07/2005).”

an extinction curve built from our own data, even if the air-mass difference among the used spectra was small: 0.10 for the polarized star and 0.05 for the non-polarized star.

As can be seen in Fig. 11, the overall shapes of the spectra taken at different derotator angles agree with one another within 1%, except at the blue edge (below 4000 Å), where the S/N is rather low, and above 8000 Å, as expected because of fringing and telluric absorption (which were not corrected for in this particular test). This confirms first of all that the night conditions were photometric. The additional absence of significant slopes or oscillations confirms that both the stellar intrinsic polarization and the polarization supposedly induced by the instrument configuration had no significant impact on our flux measurements (i.e., not more than 1%), even for different stellar polarization levels, wavelength regions, or derotator positions.

7 Second-order contamination

Low resolution spectrographs and large format CCDs allow for observations of a wide wavelength range, as needed for the Gaia SPSS grid. However, with both grisms and gratings the light from different orders can overlap, especially if no cross disperser or blocking filter is employed (Gutierrez-Moreno et al. 1994). Typically, light from the blue wavelengths of the second order can contaminate the red portion of the spectra, in the spectrographs used in our campaign. The characteristics of the contaminating light depend on both the spectral energy distribution of the observed source and the characteristics of the dispersing element. In particular, it is necessary to know the grism (Traub 1990) or grating (Szokoly et al. 2004) equation to apply a correction.

Two of our adopted instrument configurations, for EFOSC2@NTT and DOLORES@TNG, are known to possess significant second order contamination. For CAFOS@2.2 m in Calar Alto the only visible second order light started well after 10 000 Å, after the end of the first-order spectrum in the red¹⁵. For BFOSC@Cassini it was absent, both judging from the BFOSC manuals and web pages, and from the quality of our final spectra of CALSPEC¹⁶ stars. The chosen setups were unavoidable owing to our requirements to cover as much as possible the Gaia wavelength range (300–1100 nm) and to oversample the resolution of the BP and RP (Blue and Red Photometers) by a factor of up to 4–5, whenever possible.

We thus employed an adaptation of the method by Sánchez-Blázquez et al. (2006) to correct our spectra from second-order contamination. The original method requires

¹⁵ The manual mentions a second order contamination for the B200 grism starting after 6500 Å, but our data reach up to ≈8500 Å and no contaminating light was detected; for R200, the manual mentions a second-order contamination starting at 10 000 Å, but we could clearly see the second order spectrum, starting well after the end of the first order spectrum, with no overlap between the two orders.

¹⁶ <http://www.stsci.edu/hst/observatory/crds/calspec.html>

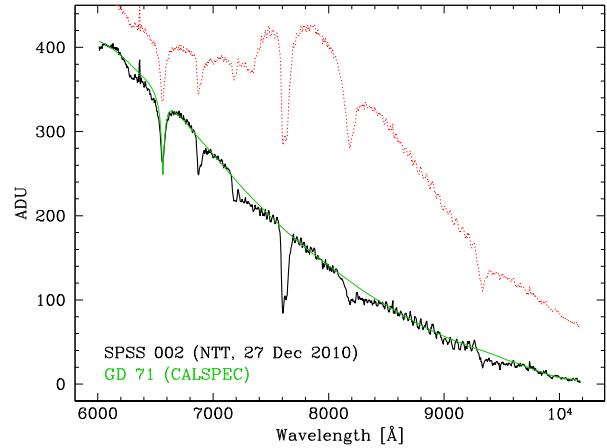


Fig. 12 Example of our second order contamination correction for a CALSPEC star (GD 71, one of the pillars) observed with NTT. The red (dotted) spectrum is our observed spectrum contaminated by second-order light coming from blue wavelengths. The black spectrum is the same spectrum after our second order correction procedure; it still contains telluric absorption bands and fringing. The green spectrum is the CALSPEC spectrum for the same star: The spectral shape is recovered to the 1% level.

observations of a blue and a red star, with well known calibrated flux in the literature. They are used to build a *response curve* for the contaminating second-order spectra, by solving the following system of equations:

$$S_b = C_1 T_b + C_2 T_{2b}, \quad (3)$$

$$S_r = C_1 T_r + C_2 T_{2r}, \quad (4)$$

where the ‘b’ subscript refers to the blue star and the ‘r’ one to the red star. S is the observed spectrum, including the second-order contamination, while T is the true spectrum obtained from the literature (thus free from second-order contamination), and T_2 is the second-order contaminating component. C_1 is the response curve for the first order light and C_2 for the second order light: once they are known, any spectrum observed with the same instrument can be corrected for second-order contamination as described by Sánchez-Blázquez et al. (2006). Since C_1 and C_2 were not computed each night in our SPSS campaign, we derived the corrected spectrum (S'_b for the blue star case) using

$$S'_b = S_b - C_2 S'_{2b}, \quad (5)$$

where S'_{2b} is the ratio S_b/C_1 , i.e., the second order contaminating spectrum, after resampling and shifting, according to the function mapping the first order into the second order wavelength range (see following sections for more details), and still not calibrated in flux.

The method by Sánchez-Blázquez et al. (2006) requires a single spectrum to cover the whole wavelength range, while in our case we always had two setups – with some overlap – observed in two slightly different moments of time. Our adaptation consisted in joining the blue and red

spectra of the observed stars, after reporting the blue spectrum to the airmass of the red one with a good extinction curve. A residual jump in the junction point does not have any effect on the quality of the final correction for second order.

While the original method by Sánchez-Blázquez et al. (2006) for deriving C_1 and C_2 can be applied directly, even in non-perfectly photometric conditions, our particular adaptation requires some care in those cases where extinction variations are large. Application of the method is safe in photometric and grey¹⁷ nights, where extinction variations are lower than $\approx 3\%$, roughly speaking.

7.1 NTT second order correction

The spectra of two well known flux standards, Feige 110 (a blue star, Landolt & Uomoto 2007) and LTT 1020 (a red star, Landolt 1992), were acquired with the blue and red EFOSC2 grisms #11 and #16 (see Table 4), used for the whole SPSS campaign. Each star was observed with a $10''$ slit to gather all the flux, and with a $2''$ slit to obtain a reliable wavelength calibration.

The spectra were extracted and wavelength calibrated as described in Sect. 6.1. We created an extinction curve from our own data, that compared well with the CTIO and ESO ones. We then joined the grism #11 and #16, reported to the same airmass, into a single spectrum, cutting them at $6000 \text{ \AA}$, i.e., well before the expected start of second order contamination. We used the absorption lines of the blue contaminating spectrum – well visible at red wavelengths – to determine the function that maps the blue contaminating wavelengths in the red wavelength domain:

$$\lambda_2 = -3815 + 3.271 \lambda_1 - 0.0001167 \lambda_1^2, \quad (6)$$

where λ_1 is the wavelength in $\text{\AA}$ of the first order, and λ_2 the corresponding wavelength of the second-order contaminating blue light. The residuals around this fit were of $2.9 \text{ \AA}$. The flux reference tables used for computing the response curves, i.e., T_b and T_r in the equations of the previous section, were obtained for LTT 1020 from CTIO¹⁸ (Hamuy et al. 1992, 1994) and for Feige 110 from the CALSPEC database (Bohlin, Dickinson & Calzetti 2001).

Once the two response curves C_1 and C_2 were derived, we computed the percentage of the blue light of a star that is expected to fall at each red wavelength, as

$$P = \frac{C'_2}{C_1}, \quad (7)$$

where C'_2 is C_2 reported into blue wavelengths using the relation above. The computation of P was repeated in two

¹⁷ Grey conditions are attained when the cloud coverage produced grey extinction variations, i.e., when extinction does not alter significantly the spectral shape (within 1%). This condition is almost always verified in the case of veils or thin clouds (Oke 1990; Pakštie & Solheim 2003) and can easily be checked a posteriori.

¹⁸ <ftp://ftp.eso.org/pub/stecf/standards/ctiostan>

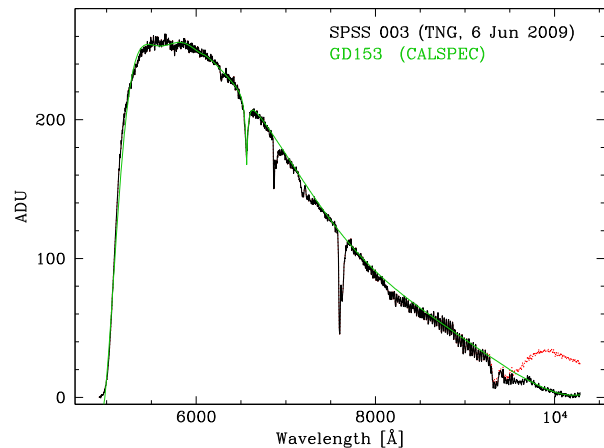


Fig. 13 Same as Fig. 12, but for one TNG spectrum of GD 153, taken in 2009 June. Apart from the telluric bands and the strong fringing that was not corrected for at this stage, the second order contamination – starting around $9000\text{--}9500 \text{ \AA}$ – is very well corrected even with our simple assumptions on the grism equation (see text).

different epochs (using LTT 1020 and Feige 110 on 2008 November 30, and LTT 9239 and Feige 110 on 2011 August 28) and appeared stable within 1%. The contamination starts around $6000 \text{ \AA}$, as shown in Fig. 12, and becomes severe around $6500 \text{ \AA}$.

7.2 TNG second order correction

The DOLORES red LR-R grism used in our SPSS campaign had a built-in order-blocking filter for wavelengths bluer than $\approx 5000 \text{ \AA}$, but unfortunately some second-order contamination from blue light was observed after $\approx 9500 \text{ \AA}$ ¹⁹. To obtain the percentual contamination, P , we used two well known standards: HZ 44 (a blue star, Landolt & Uomoto 2007) and G 146-76 (a red star, Høg et al. 2000). A third star, GD 153 (another blue star, Bohlin, Colina & Finley 1995) was observed to independently test the results. We also observed all LR-R spectra together with the Johnson I broadband filter in front of the spectrograph, to recover independently the uncontaminated spectral shape. They were observed with the LR-B and LR-R grisms (see Table 4) both with a $5''$ slit, the widest available at the time, and with the $2''$ slit for a better wavelength calibration. We also tested GD 153 and KF06T2 (Reach et al. 2005) on 2009 June 8 with similar results: the maximum difference between the two P curves was always lower than $\approx 0.3\%$. All spectra were extracted and wavelength calibrated as in the NTT case, and reported to the same airmass with a tabulated extinction curve²⁰. We joined the spectra arbitrarily at

¹⁹ We have also observed an effect that could be ascribed to second order contamination in the blue LR-B grism, starting at $6000 \text{ \AA}$ approximately. Because that region is always adequately covered by the red LR-R grism, we simply cut the contaminated section of the LR-B spectra.

²⁰ http://www.ing.iac.es/astronomy/observing/manuals/ps/tech_notes/tn031.pdf

5850 Å, i.e., well before the start of second-order contamination.

To determine a wavelength mapping relation, we obtained helium lamp spectra with $t_{\text{exp}} = 1$ h and a B filter to block first order light. Of the three measurable emission lines in LR-R, one was at 9553 Å, where no He line is expected. We thus could not derive a good fit of the function mapping the second order, as done for the NTT case. We used the simple formula $\lambda_1 = 2\lambda_2$ that derives from the simplest form of the grating equation

$$m\lambda = \sigma(\sin\beta \pm \sin\alpha). \quad (8)$$

We then derived the C_1 and C_2 curves and the percentage of contaminating flux using this basic assumption: the emission line appearing at 9553 Å corresponded well with the 4713 Å He emission line.

We tested the correction on TNG observations of various CALSPEC and literature standards (one example is reported in Fig. 13). We obtained good results, having differences with SPSS literature spectra of 1–2 % at most, for several stars. However, the presence of fringing and a deep telluric band in the contaminated region (after 9500 Å) causes high residuals which can be $\lesssim 10$ % in the worst cases. For many SPSS we had observations from other telescopes, but for the few SPSS observed only with TNG, a viable solution will be to replace the flux tables after 9500 Å with model spectra (see also Bohlin 2007).

8 Conclusions

To build a grid of ~ 100 SPSS with ≤ 1 % internal errors and ≤ 1 –3 % external errors (with respect to the Vega calibration by Bohlin 2007) we have carried out a systematic study of instrumental effects that could have an impact on the flux calibration of SPSS flux tables and integrated magnitudes, finding methods to beat each effect separately down, until their impact on fluxes (or ADUs) was below 1%. We collected from the literature the available methods and strategies to both evaluate and remove those instrumental effects, and carried out specific daytime and nighttime observations to quantify their effects. When necessary, we adapted the methods to our case, or developed our specific tests and data-reduction methods. Whenever feasible – given the difficulties to obtain telescope time for such programs – we compared different approaches to the same problem. In particular,

- we characterized the employed CCDs to evaluate their minimum acceptable exposure time, the shutter delay, and the linearity limit, using different methods, and deriving recommendations for our SPSS observing campaigns (see Tables 2 and 3);
- we carried out a calibration frames monitoring plan, that provided recommendations on the minimum frequency of each calibration type observations (see Table 5) for our observations;

- as part of our calibration frames monitoring, we devised a specific procedure to correct our spectra for fringing, after manipulation of spectroscopic flat fields (following the strategy by Malumuth et al. 2003) and using the IRAF *telluric* task;
- we have also quantitatively studied the effect of gravity-induced lamp flexures at DOLORES@TNG; we found shifts of about 3 pixels, with 0.01–0.06 errors, corresponding to almost 8 Å in the wavelength space, confirming that for this type of instruments day-time lamps are not sufficiently accurate. If night-time lamps required too much observing time, we acquired high S/N day-time lamps and shifted them onto the night-time ones to perform the wavelength calibration;
- we studied the effect of instrument-induced polarization with both NTT and TNG, the two instruments that should have the largest polarization in our SPSS campaigns. We observed stars with polarization levels up to 4 % at different instrument rotation angles, and concluded that even in the case of SPSS with those intrinsic polarization levels, the flux calibration remained stable within 1%;
- we tested and adapted to our case the method described by Sánchez-Blázquez et al. (2006) to correct our spectra for second order contamination; our method can be applied in photometric and grey nights with residuals generally within 1–2 % for NTT and TNG.

In summary, a few of the examined instrumental effects turned out to be negligible, while we devised specific methods to correct for the remaining ones, reducing their impact on the flux calibration of SPSS magnitudes and spectra to ≤ 1 %. The methods can be applied to a wide range of observational programs on similar instruments.

The actual reduction and analysis of the SPSS spectra will be presented in a subsequent paper, where we will further reduce the residuals of instrumental effect corrections by combining several independent observations for each SPSS. We will also replace the low S/N blue and red borders (roughly ≤ 3800 – 4000 Å and ≥ 9500 Å) with theoretical or semi-empirical templates.

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Mexico (under program IDs #27 and #28, with observations in 2008 August and 2010 July).

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