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Anharmonic vibrational spectroscopy of Polycyclic Aromatic Hydrocarbons (PAHs)

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While powerful techniques exist to accurately account for anharmonicity in vibrational molecular spectroscopy, they are computationally very expensive and cannot be routinely employed for large species and/or at non-zero vibrational temperatures. Motivated by the study of Polycyclic Aromatic Hydrocarbon (PAH) emission in space, we developed a new code, which takes into account all modes and can describe all IR transitions including bands becoming active due to resonances as well as overtones, combination and difference bands. In this article, we describe the methodology that was implemented and discuss how the main difficulties were overcome, so as to keep the problem tractable. Benchmarking with high-level calculations was performed on a small molecule. We carried out specific convergence tests on two prototypical PAHs, pyrene ($C_{16}H_{10}$) and coronene ($C_{24}H_{12}$), aiming at optimising tunable parameters to achieve both acceptable accuracy and computational costs for this class of molecules. We then report the results obtained at 0 K for pyrene and coronene, comparing the calculated spectra with available experimental data. The theoretical band positions were found to be significantly improved compared to harmonic Density Functional Theory (DFT) calculations. The band intensities are in reasonable agreement with experiments, the main limitation being the accuracy of the underlying calculations of the quartic force field. This is a first step towards calculating moderately high-temperature spectra of PAHs and other similarly rigid molecules using Monte Carlo sampling.

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I. INTRODUCTION

Polycyclic Aromatic Hydrocarbons (PAHs) are a family of organic compounds composed of aromatic rings containing carbon atoms, and whose peripheral bonds are saturated by hydrogen atoms. PAHs (or closely related species) are thought to be the carriers of the so-called Aromatic Infrared Bands (AIBs) at ~ 3.3 , 6.7, 7.7, 8.6, 11.3 μm .^{1,2} These features are among the strongest emission features observed in the interstellar medium.³ Therefore, the infrared (IR) spectroscopy of PAHs is of paramount importance in astrophysics. PAHs absorb starlight in the visible and ultraviolet domain via electronic transitions and re-emit most of this energy in the mid-IR in a vibrational de-excitation cascade. In such model PAHs should emit most of the flux in the AIBs from highly vibrationally excited states corresponding to a thermal temperature of the order of ~ 1000 K.⁴ Each AIB results from the superposition of the emission of molecules at different temperatures, each of which being the superposition of a large number of individual vibrational transitions from a statistical distribution of excited states. Due to anharmonicity, each of these states is expected to have a shifted position with re-

spect to the corresponding 1-0 fundamental transition. Temperature-dependent effects of anharmonicity in modelling astronomical PAH spectra received only relatively sparse attention.^{5,6} These works followed measurements on a few neutral gas-phase PAHs, which reported the overall band shift and broadening of the most intense fundamental bands with temperature up to ~ 900 K.⁷ Similar data became also available from theory.⁸⁻¹⁰ In most AIB modelling studies though, theoretical calculations are performed at the harmonic approximation level. The harmonic frequencies and first derivatives of the dipole moment are easily obtained using commonly available quantum chemistry computer codes that implement many different levels of theory for the electronic states of molecules such as the Hartree-Fock, Density Functional Theory (DFT) or Coupled-Cluster levels. Databases of such harmonic computed spectra are available^{11,12} and include PAHs with up to a few hundreds C atoms. Anharmonicity in the Hamiltonian is then typically taken into account by some overall frequency scaling factor chosen in order to obtain a better agreement with laboratory data measured in rare-gas matrices at low temperature. Note that laboratory data are themselves affected by matrix effects which can be modeled as anharmonicity effects.¹³ Scaling procedures can be refined in order to take into account the nature of the vibrational modes.^{14,15} In any case, this completely neglects the anharmonic shifts of hot bands with respect to the fun-

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damental, just correcting (empirically) for the effect of anharmonicity on the absolute position of the 1-0 transitions. In order to perform simulations that can be compared with the AIB spectrum, temperature effects are then only simply included by adding some ad-hoc band redshift and broadening, see e. g. Ref. 16 and the discussion of synthetic PAH spectra by Boersma *et al.*¹⁶ and Pauzat.¹⁷

In order to progress in the analysis of the astrophysical spectra, one has clearly to go beyond this simplified treatment. Efforts should be dedicated to describe connected states at non-zero temperatures. For this, the approximation of the electric dipole moment by its Taylor expansion truncated to the linear terms is not sufficient. Indeed this approximation only describes vibrational transitions connecting harmonic states which differ by only one quantum in only one (IR-active) vibrational mode. Any other vibrational transitions, including e.g. overtones, combination and difference bands, are thus completely neglected at this level of treatment. The reason for using such seemingly crude approximations is that performing actual anharmonic calculations for anything but the smallest PAHs quickly becomes prohibitively expensive from a computational point of view. Complete variational calculations of rovibrational levels with *ab-initio* potential energy and electric dipole surfaces are only feasible for very small molecules with up to four or five atoms.^{18–20} However, the field is making constant progress and variational vibrational calculations can now be achieved for systems with up to 11 atoms.²¹ Vibrational Self Consistent Field (VSCF) methods, based on the *ansatz* that the vibrational wave-function can be represented by a product of functions depending on a single coordinate,^{22,23} can treat systems as large as a couple of hundreds of atoms, including PAHs.²⁴ However, they neglect all correlation between different degrees of freedom and are generally not accurate enough for our purposes. Techniques like the Vibrational Mean Field Configuration Interaction (VMFCI)^{25,26} improve on this by contracting several modes into one to describe correlated wave-functions, and can scale up to molecules as big as naphthalene ($C_{10}H_8$), provided that the modes involved in important resonances form disjoint subsets of limited sizes. For much larger molecules, perturbation theory up to second order (VPT2) or beyond,²⁷ is the method of choice to treat inter-modes correlation starting from VSCF solutions, or both intra-mode anharmonicity and inter-modes correlation starting from the harmonic approximation. However, this treatment breaks down in the case of near degeneracy, be it accidental or due to symmetry. While analytical VPT2 equations have existed for long, symmetric and spherical tops deserve special care due to symmetry degeneracies of modes²⁸, furthermore, accidental near-degeneracies of energy levels have to be detected and treated in a specific manner. Since in PAHs vibrational modes of the same type cluster at given frequencies, near-degeneracies are unavoidable for these molecules, in particular in the spectral region of C-

H stretches and C-H bends. Resonances must therefore be taken out of the perturbation treatment, and explicitly accounted for, without omitting any interacting state nor resonant term of the vibrational potential, using for example the generalized second-order vibrational perturbation theory (GVPT2).^{28–31}

This GVPT2 approach requires to first define some appropriate thresholds as a function of the desired target accuracy to determine which terms are to be considered “resonances” and which can safely be treated using perturbation theory. Then, one defines polyads of interacting harmonic vibrational states connected by resonant terms and solves the corresponding variational problems, while perturbative corrections for non-resonant terms can be added either before or after. In principle, for vanishingly small thresholds all terms are considered resonances and one goes back to the limit of a full variational calculation in a basis of harmonic vibrational states. In general, in addition to thresholds for discriminating between resonant and non-resonant terms, one also needs to implement some truncation scheme to keep the size of the polyads and of the associated variational problems small enough to be tractable without degrading the accuracy of the results. This kind of calculation has been implemented and successfully applied to compute the vibrational spectrum of PAHs involving transitions from the ground vibrational state and possibly some of the lowest lying ones in energy.^{32–34} Solving for all vibrational states up to higher energies quickly becomes prohibitive, since the number of states involved explodes, behaving as a multi-factorial with energy. However, if one is interested in reproducing the overall envelope of the spectral profile resulting from the superposition of hot bands as a function of vibrational excitation (or temperature), a possible approach is to give up a complete solution in favor of a Monte Carlo sampling of vibrational states.^{8,10,35–37} We developed and implemented a code suitable for this approach,³⁸ and in this paper we test its behaviour benchmarking its calculations of the 0 K spectrum against both experimental measurements for some small-medium PAHs and very high-level VMFCI calculations for a smaller species, namely ethylene oxide. In Section II we describe the Van Vleck perturbation approach and we give computational details. Then in Section III we proceed to compare the outcomes of our calculations with reference calculations, which we use to define a strategy for choosing tuning parameters that provide both acceptable accuracies and computational costs. We also compare with experimental results, assessing the accuracy that can be expected from these calculations for fundamental bands, bands becoming weakly active due to resonances, and also overtone/combination/difference bands stemming from transitions mediated by the second order terms in the expansion of the electric dipole moment. In Section IV we discuss our results, assessing the applicability of the method and its estimated computational costs for future calculations of spectra of PAHs and similarly rigid molecules considering all significantly

populated levels at moderately high T using Monte Carlo sampling.

II. COMPUTATIONAL FORMALISM

A. Van Vleck perturbation theory

In this section, we give a brief review of the Van Vleck approach to perturbation theory applied to molecular vibrations.^{39–42} The anharmonic vibrational Hamiltonian H is written as

$$H = H^{(0)} + \lambda H^{(1)} + \lambda^2 H^{(2)}, \quad (1)$$

where λ is the perturbation parameter. The zero-order Hamiltonian is given as the harmonic normal-mode Hamiltonian

$$H^{(0)} = H^{\text{harm}} = \sum_i \frac{\hbar\omega_i}{2} (p_i^2 + q_i^2), \quad (2)$$

and the first- and second-order Hamiltonian are given by the cubic and quartic expansion of the vibrational potential

$$H^{(1)} = \frac{1}{3!} \sum_{ijk} \left(\frac{\partial^3 V}{\partial q_i \partial q_j \partial q_k} \right)_0 q_i q_j q_k, \quad (3)$$

$$H^{(2)} = \frac{1}{4!} \sum_{ijkl} \left(\frac{\partial^4 V}{\partial q_i \partial q_j \partial q_k \partial q_l} \right)_0 q_i q_j q_k q_l. \quad (4)$$

Higher-order terms in the potential are neglected. The Van Vleck procedure relies on an infinitesimal contact transformation which will lead to an effective Hamiltonian. The contact transformation is defined by $T = \exp(S)$ where the operator S is anti-hermitian $S^\dagger = -S$. This operator can be expanded up to the second-order in the perturbation parameter

$$S = \lambda S^{(1)} + \lambda^2 S^{(2)} + \dots \quad (5)$$

The transformed Hamiltonian $\tilde{H} = T^\dagger H T$ up to the second-order in the perturbation parameter is then written as

$$\tilde{H} = \tilde{H}^{(0)} + \lambda \tilde{H}^{(1)} + \lambda^2 \tilde{H}^{(2)} + \dots, \quad (6)$$

where

$$\tilde{H}^{(0)} = H^{(0)}, \quad (7)$$

$$\tilde{H}^{(1)} = H^{(1)} - [S^{(1)}, H^{(0)}], \quad (8)$$

$$\begin{aligned} \tilde{H}^{(2)} = & H^{(2)} - [S^{(2)}, H^{(0)}] - [S^{(1)}, H^{(1)}] \\ & + \frac{1}{2} [S^{(1)}, [S^{(1)}, H^{(0)}]]. \end{aligned} \quad (9)$$

Each term of the contact transformation $S^{(i)}$ is chosen such that it cancels out all non-diagonal terms of $\tilde{H}^{(i)}$. This is in general possible as long as no resonance occurs. For example, in the case of the anharmonic vibrational Hamiltonian of Eqs. (1), (2), (3) and (4), one can find in Ref. 43 the following expression for $S^{(1)}$

$$\begin{aligned} S^{(1)} = & \sum_{ijk} \frac{i\omega_k}{3! \Delta_{ijk}} \left(\frac{\partial^3 V}{\partial q_i \partial q_j \partial q_k} \right)_0 \left(2\omega_i \omega_j p_i p_j p_k \right. \\ & \left. + (\omega_i^2 + \omega_j^2 - \omega_k^2) (q_i q_j p_k + q_i p_k q_j + p_k q_i q_j) \right), \end{aligned} \quad (10)$$

where Δ_{ijk} is defined by

$$\begin{aligned} \Delta_{ijk} = & (\omega_i - \omega_j - \omega_k) (\omega_i + \omega_j - \omega_k) \\ & \times (\omega_i - \omega_j + \omega_k) (\omega_i + \omega_j + \omega_k) \end{aligned} \quad (11)$$

This expression can then be used in Eq. (9) and $S^{(2)}$ can be determined. In fact, at this point, it is not necessary to obtain an explicit expression for $S^{(2)}$. One just have to assume that the second order transformation exists and is not divergent, which is the case if no resonances occur. Then, the eigenvalues of the effective Hamiltonian are given by the usual Dunham expansion

$$\begin{aligned} E_{\mathbf{n}} = & \chi_0 + \sum_i \omega_i (n_i + 1/2) \\ & + \sum_{i \leq j} \chi_{ij} (n_i + 1/2) (n_j + 1/2) \end{aligned} \quad (12)$$

where the expression for the anharmonic coefficients χ_0 , and χ_{ij} can be found easily in the literature.^{27,30,43}

In the case of exact or near resonances, the expression for the first order contact transformation Eq. (10) cannot be used as it is. Indeed, Fermi resonances defined by the occurrences of triplets i, j and k such that $\pm\omega_i \pm \omega_j \pm \omega_k \simeq 0$ would result in a diverging first-order contact transformation. This implies that the contact transformation should be determined to infinite order. In the case of quasi-degenerate perturbation theory these resonant terms are simply excluded from the definition of the contact transformation. As a result the first-order transformed Hamiltonian $\tilde{H}^{(1)}$ is not diagonal. Similarly the occurrence of Darling-Dennison⁴⁴ (DD) resonances arises by diverging terms in the second-order contact transformation $S^{(2)}$ and occurring when $\pm\omega_i \pm \omega_j \pm \omega_k \pm \omega_l \simeq 0$. These terms should also be excluded, therefore preventing for a full cancellation of the non-diagonal terms in the second-order transformed Hamiltonian $\tilde{H}^{(2)}$. Note that DD resonances can occur even if there are no corresponding quartic derivatives since quartic coupling terms appear due to the application of the first order contact transformation. After using the contact transformation, we obtain an effective and non-diagonal Hamiltonian $\tilde{H}$ whose eigenvalues are identical to the Hamiltonian H at the second order of the perturbation. This effective Hamiltonian can be diago-

nalized using a finite variational basis.

After being determined, the infinitesimal contact transformation can be applied to any operators such as the dipole operator μ resulting in a transformed operator accounting for both mechanical and electric anharmonicity up to the second order in the perturbation parameter. Similarly as for the Hamiltonian, the dipole operator is expanded as

$$\mu = \mu^{(0)} + \lambda\mu^{(1)} + \lambda^2\mu^{(2)} + \dots, \quad (13)$$

where $\mu^{(0)}$ is the permanent dipole, and where $\mu^{(1)}$ and $\mu^{(2)}$ are respectively the linear and quadratic terms of the expansion of the dipole as a function of the normal mode coordinates about the equilibrium position, they are written as

$$\mu^{(1)} = \sum_i \mu_i^{(1)} q_i = \sum_i \left(\frac{\partial \mu}{\partial q_i} \right)_0 q_i, \quad (14)$$

$$\mu^{(2)} = \frac{1}{2} \sum_{ij} \mu_{ij}^{(2)} q_i q_j = \frac{1}{2} \sum_{ij} \left(\frac{\partial^2 \mu}{\partial q_i \partial q_j} \right)_0 q_i q_j. \quad (15)$$

Applying the perturbative transformation, the expression for the transformed dipole operator up to the second-order in the perturbation parameter is given by

$$\tilde{\mu} = \mu^{(0)} + \lambda\mu^{(1)} + \lambda^2 \left(\mu^{(2)} - [S^{(1)}, \mu^{(1)}] \right) + \dots, \quad (16)$$

where one should remember that resonant terms are omitted from the definition of $S^{(1)}$, as explained before. The intensities of the transitions are then calculated by computing the matrix elements of the dipole operator between two eigenstates of the effective Hamiltonian $\tilde{H}$. The intensity of a vibrational transition, neglecting rotational degrees of freedom, is then written as

$$I_{ab} = \frac{2\pi}{3\hbar c \epsilon_0} \omega_{ab} \langle \psi_a | \tilde{\mu} | \psi_b \rangle^2, \quad (17)$$

where $\omega_{ab} = \omega_b - \omega_a$ are the transitions between eigenfrequencies and where $|\psi_a\rangle$ are the eigenstates of $\tilde{H}$. I_{ab} is comparable to the experimental band intensity integrated over its rotational substructure.

B. Polyad construction and variational problem

In this section we give the computational details of our implementation of the Van Vleck perturbation theory described in Sec. II A. The computational protocol described in this section was implemented in our code: AnharmoniCaOs (the Cagliari-Orsay model for anharmonic molecular spectra in 2nd order perturbation theory).

The first step is to identify resonances. For this purpose, starting from an initial harmonic state $\mathbf{n} = (n_1, \dots, n_N)$, we first identify all harmonic states $\mathbf{n}' \neq \mathbf{n}$ directly connected to $\mathbf{n}$ through the cubic couplings.

For these states, the ratio of the coupling term $V_{\mathbf{n}\mathbf{n}'}^{(1)} = \langle \mathbf{n} | H^{(1)} | \mathbf{n}' \rangle$ and the harmonic energy difference between these states is used to define a resonance. A Fermi resonance occurs then if

$$\left| \frac{V_{\mathbf{n}\mathbf{n}'}^{(1)}}{E_{\mathbf{n}}^{(0)} - E_{\mathbf{n}'}^{(0)}} \right| \geq r_3 \quad (18)$$

where r_3 is a small threshold parameter. A list of Fermi resonances is thus built. From this list the first-order contact transformation $S^{(1)}$ can be defined by excluding all terms corresponding to a resonance as in Eq. (18). Applying the first-order contact transformation $S^{(1)}$ onto the Hamiltonian H we obtain the second-order part of the transformed Hamiltonian

$$H^{(2)} - [S^{(1)}, H^{(1)}] + \frac{1}{2} [S^{(1)}, [S^{(1)}, H^{(0)}]]. \quad (19)$$

This operator is then used to identify harmonic states $\mathbf{n}'' \neq \mathbf{n}$ directly connected to the initial state $\mathbf{n}$ through the quartic couplings which include the effect of the first-order transformation. Darling-Dennison (DD) resonances are then defined by

$$\left| \frac{V_{\mathbf{n}\mathbf{n}''}^{(2)}}{E_{\mathbf{n}}^{(0)} - E_{\mathbf{n}''}^{(0)}} \right| \geq r_4, \quad (20)$$

where $V_{\mathbf{n}\mathbf{n}''}^{(2)}$ is non-diagonal matrix element of operator (19). The list of DD resonances having been built, the second-order contact transformation can be defined implicitly. This transformation is used to define the transformed Hamiltonian $\tilde{H}$ (Eq. 6).

After this preliminary work, polyads are iteratively constructed as resumed in Fig. 1. A list of starting states is built by including the initial harmonic state $\mathbf{n}$ and the final harmonic states $\hat{\mathbf{n}}$ reached by a dipolar transition through the transformed dipole operator $\tilde{\mu}$. Each starting state will constitute the start of a polyad. Running through the resonances the list of states in each polyad is increased iteratively. When two polyads share a common state, they are merged. Symmetry can be taken advantage of by considering only states of a given symmetry type at a time. This way, one could think of separating explicitly the problem into a number of subproblems. However, in practice, given the way in which polyads are recursively built by following resonances, resonant terms can only be non-zero between harmonic states of the same symmetry. Hence symmetry separation is automatically enforced and there would not be any significant performance gain by implementing the explicit separation, which would make the code more complicated. It is only necessary to include a cutoff parameter to remove the numerical noise, introduced by cubic and quartic, non-symmetry-adapted Hamiltonian terms, in case the separate code used to obtain the potential does not explicitly use symmetry itself and thus produces very small non-zero spurious terms.

This procedure, depending on how resonances combine, can lead to polyads of very large, in principle even infinite, size. To keep the problem tractable, a truncation mechanism is therefore necessary. To obtain finite size polyads we used a “cost model”. The very initial harmonic states, i. e. the initial one and the ones connected to it by permitted transitions, are assigned a “budget” initialized to 1. Additional states newly added to a polyad “inherit” the budget of the state they are connected to, minus a “cost” which is inversely proportional to the “strength” of the resonance (the ratio between the non-diagonal element and the difference of the diagonal elements of the connected harmonic states) divided by a tuning parameter h . Therefore, a very strong resonance will add many states to a polyad, whereas a weak one will add few, possibly only one. Larger values of h will produce larger polyads, reducing truncation errors at the price of an increased computational cost.

To each distinct polyad corresponds a variational problem, for which an effective Hamiltonian is defined. What we want to obtain is an accurate description of the anharmonic states containing a non-negligible component of the initial harmonic starting states (the very first harmonic state and the ones connected to it by allowed transitions). Depending on how strong resonances are, and on chosen tuning parameter values, the eigensolutions we need may be a small fraction of the size of a given polyad. If the polyad is relatively small, complete solution using a standard direct method (e.g. Divide and Conquer⁴⁵ or Relatively Robust Representations⁴⁶) is more efficient. If, instead, the polyad is large (e. g. tens of thousands of harmonic states or more) and a small fraction of the eigensolutions is needed (e. g. less than $\sim 10\%$), then an iterative method like Jacobi-Davidson,^{47,48} explicitly tuned to select the eigenvalues with the largest component in the space of the initial starting states, may become more competitive. After the eigensolutions are found, the line intensities between anharmonic states are computed using Eq. (17). The final spectrum is then built from individual transitions weighted by the square of the component of the starting state it contains. In this way, when the same polyad is obtained more than once from different starting states, each transition will eventually converge to its exact value for complete coverage of the starting states.

C. Electronic structure calculations

The Van-Vleck method described in the previous sections relies on the use of a quartic force field corresponding to harmonic frequencies, cubic and quartic derivatives of the potential energy surface as well as the first and second derivatives of the dipole moment. These parameters are easily obtained from quantum chemical calculations. All electronic structure calculations were performed using the Gaussian09 suite of programs.⁴⁹ Geometry optimizations and frequency calculations were performed

using density functional theory (DFT) with the hybrid functional B97-1,⁵⁰ and with the TZ2P⁵¹ and 6-31G*^{52,53} basis sets. We chose this exchange-correlation functional over some more commonly used ones as e.g. B3-LYP^{54,55} because the former was found to provide more accurate band position and intensities in particular for aromatic molecules.^{32,33,56} As in Ref. 56 we used a (150,770) grid for the Kohn-Sham (KS) integration and a (75,194) grid for the Coupled Perturbed Kohn-Sham (CPKS) steps.

For each molecule, geometry optimization was first performed, then harmonic frequencies and normal modes were computed. Cubic and quartic normal mode derivatives of the potential around the equilibrium position were obtained using numerical differentiation of the analytical Hessian matrix, using a displacement step of $0.01 \text{ \AA} \cdot \text{amu}^{1/2}$. The calculation of all quartic derivatives requires $36N^2$ Hessian calculations, where N is the number of atoms. For large systems like pyrene or coronene such a calculation could only be performed with semi-empirical force-field and not with *ab-initio* or DFT methods. Note that only the semi-diagonal involving at most two different indices are needed for perturbation theory, the other terms just contribute to DD resonances. Therefore, we only computed the quartic derivatives with two identical indices $\left(\frac{\partial^4 V}{\partial q_i \partial q_i \partial q_j \partial q_k} \right)_0$. Similarly, second derivatives of the dipole moment were obtained by numerical differentiation of the analytical dipole first derivatives. Invariance of derivatives on the order of differentiation was used to check the numerical stability.

III. RESULTS

We now compare the results we obtained with AnharmoniCaOs with some reference calculations and experimental results. Comparison with reference high-level vibrational calculations performed with the VM-FCI method using exactly the same quartic force field will enable us to benchmark the effect of the different user-defined thresholds. In this way we can separate the effect of using the van Vleck method, and different levels of resonance detection and truncation of the size of variational calculations, from the effect of using a quartic force field and, in turn, from the level of theory used to obtain it. On the other hand, comparison with available experimental data for PAHs enables us to gauge the accuracy of this method *including also* the effect of truncating the nuclear potential to the quartic expansion and of the level of theory used to compute it. In addition, we perform some convergence tests on pyrene and coronene, to study the accuracy vs. computational cost as a function of the tunable parameters of AnharmoniCaOs specifically for the family of PAHs.

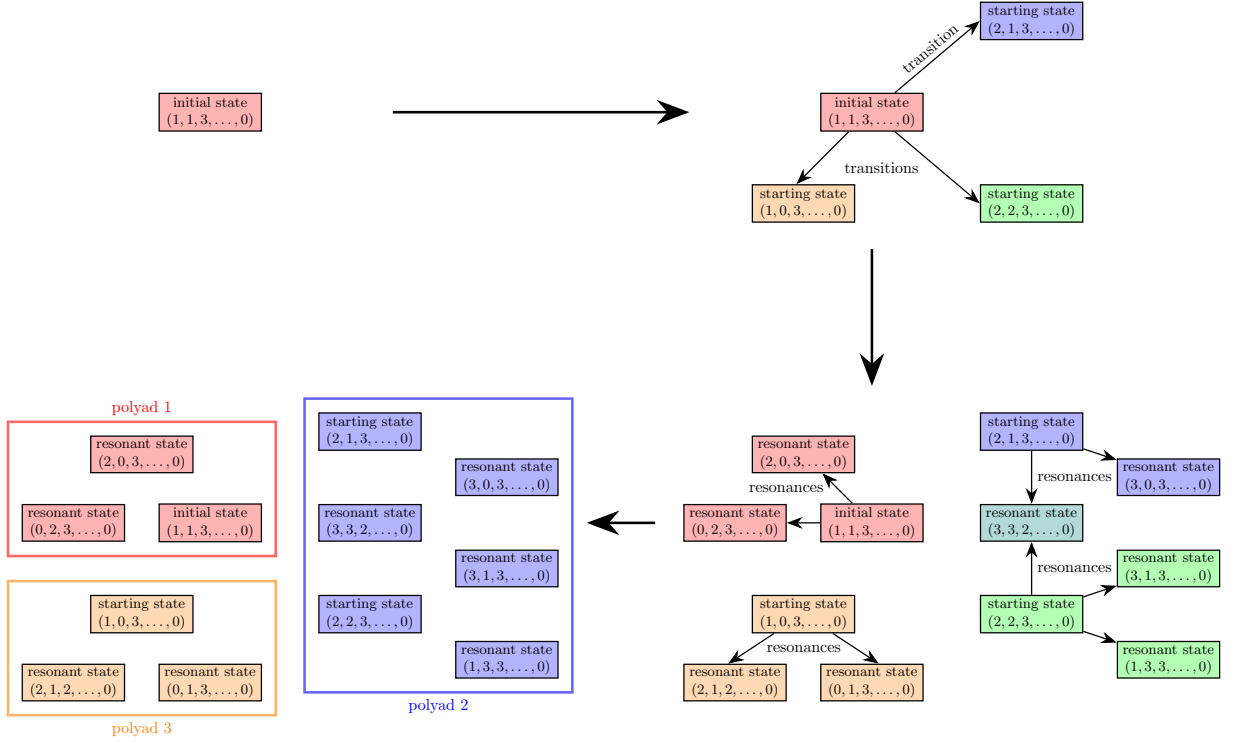


FIG. 1. Schematic algorithm of polyad definition

A. Comparison with VMFCI calculations for Ethylene oxide

A new VMFCI calculation of all vibrational energy levels of ethylene oxide up to about 3600 cm^{-1} has been performed (see supplementary material). The results compare favorably with previously published ones,^{21,57–59} using the same quartic force field for the nuclear potential. We used the very same force field with AnharmoniCaOs, and compared the results with the VMFCI ones for different values of the tuning parameters r_3 , r_4 , and h . Table I shows how the accuracy of AnharmoniCaOs changes for decreasing values of $r = r_3 = r_4$, while keeping fixed the parameter $h = 4$. Decreasing values of r cause more and more terms to be taken out of the perturbative treatment, and instead considered as resonances, in setting up and solving polyads in a variational way. For comparison, we also list in the same Table I the VSCF results, which we obtained as an intermediate step of our reference VMFCI calculation. Clearly, the VSCF discrepancies are much larger than those of our GVPT2 results, for all sets of tuning parameters used. Conversely, Table II compares the results of AnharmoniCaOs for fixed $r = r_3 = r_4 = 0.05$ and varying h . Interestingly enough, the accuracy of AnharmoniCaOs steadily improves for larger and larger h (at the price of a considerable increase in computational cost), while this is not the case for r . Indeed, we see in Table I that the best accuracy, for $h = 4$, is obtained with $r \approx 0.3$, and it gets *worse*, not better, if r is reduced while keeping h fixed. This is due to the way in which polyads are built, and their

truncation mechanism: no matter how weak it is, a resonance *always* adds at least one basis state to the starting ones. The idea behind this is that it is pointless to add a term to the list of resonances and then just neglect it because it is too weak. Conversely, if one decreases too much r without in parallel increasing h one gets an unbalanced truncation, with some weakly coupled basis states included (namely the ones connected to starting states via very weak resonances) while others, comparatively more strongly coupled, are not. Indeed, Table II shows that with $r = 0.05$ one does obtain more accurate results than with $r = 0.3$ *provided* that one then uses a large enough value of h to allow for a balanced truncation of polyads. All in all, it appears that $r \approx 0.3$ and $h \approx 4$ provide a good compromise in terms of accuracy versus computational cost.

B. Convergence tests on pyrene

The quartic force field and first and second derivatives of the electric dipole moment were obtained using DFT with the B-971⁵⁰ exchange-correlation functional and the TZ2P⁵¹ Gaussian basis set, as described in Sect. II C. We then used the AnharmoniCaOs code with the resulting quartic force field and second order Taylor expansion of the electric dipole, using the harmonic ground vibrational state as a starting state. We started with $r = 0.3$ and $h = 4$, based on the results of Sect. III A, as a reasonable compromise between speed and accuracy, bearing in

mode	VMFCI	VSCF	$r=1.0$	$r=0.5$	$r=0.3$	$r=0.2$	$r=0.1$	$r=0.05$	$r=0.01$
1	2920 ^a	3015	2955	2955	2918	2918	2918	2920	2930
2	1496	1525	1502	1502	1502	1503	1502	1503	1503
3	1271	1286	1271	1271	1271	1271	1271	1272	1272
4	1122	1160	1128	1128	1128	1128	1129	1127	1132
5	879	891	878	878	878	878	879	878	879
6	3029	3108	3043	3043	3043	3043	3060	3065	3097
7	1148	1172	1154	1154	1154	1154	1156	1156	1166
8	1018	1050	1025	1025	1025	1025	1027	1028	1037
9	2910	3032	2920	2920	2920	2921	2921	2920	2949
10	1468	1487	1474	1474	1474	1474	1475	1475	1482
11	1124	1158	1131	1131	1131	1131	1132	1132	1133
12	822	842	822	822	822	822	822	822	826
13	3041	3125	3058	3058	3058	3058	3055	3054	3085
14	1146	1169	1151	1151	1151	1151	1152	1152	1152
15	793	837	802	802	802	802	802	803	803
root mean square error		55.3	12.1	12.1	8.0	8.0	10.8	11.8	25.2

^a VMFCI assignments of step n are in terms of eigenstates of step $n - 1$ and not in terms of the initial harmonic oscillator (HO) basis functions. So, we use Thomas *et al.*²¹ assignments to relate the VMFCI frequencies to those assigned in this work on the basis of the dominant HO basis function. In fact, VMFCI assignments coincides with those of Ref. 21 except for ν_1 see supplementary material.

TABLE I. Fundamental frequencies for ethylene oxyde as a function of the threshold r and for $h = 4$.

mode	VMFCI	$h=4$	$h=8$	$h=16$	$h=18$
1	2920 ^a	2920	2920	2920	2918
2	1496	1503	1503	1503	1503
3	1271	1272	1272	1271	1271
4	1122	1127	1127	1127	1127
5	879	878	878	878	878
6	3029	3065	3065	3053	3043
7	1148	1156	1156	1155	1155
8	1018	1028	1028	1027	1026
9	2910	2920	2919	2918	2917
10	1468	1475	1474	1474	1474
11	1124	1132	1132	1132	1132
12	822	822	822	822	822
13	3041	3054	3054	3051	3050
14	1146	1152	1152	1152	1152
15	793	803	803	803	802
root mean square error		11.8	11.8	8.9	7.0

^a Same remark as in Tab. I.

TABLE II. Fundamental frequencies for ethylene oxyde as a function of the parameter h and for $r = 0.05$.

mind that our final goal will eventually be to perform (at least) tens or hundreds of thousands of such individual calculations to obtain a good enough Monte Carlo sampling of the energy dependence of vibrational bands. Still, in the case of pyrene, we conducted exploratory calculations where these parameters were pushed further, to check convergence with respect to the number of resonances and the number of states being included in the polyads to be explicitly diagonalised. For the smallest value of the threshold ($r = 0.05$) we obtained very large

polyads. The size of the largest polyad, keeping fixed $h = 8$, ramped up from ~ 650 with $r = 0.3$, to ~ 2700 with $r = 0.2$, to ~ 15000 with $r = 0.1$, and to ~ 36000 with $r = 0.05$ harmonic states. In particular, the polyads containing the states involved in the fundamental transitions in the C–H region include states spanning an energy range increasing from ~ 3015 – 5120 cm^{-1} for $r = 0.3$, to ~ 2540 – 5120 cm^{-1} for $r = 0.2$, ~ 500 – 9800 cm^{-1} for $r = 0.1$, and ~ 450 – 10000 cm^{-1} for $r = 0.05$.

Some spectra for various values of r are shown in Fig. 2 for the region of C–H stretches and in Figs. 3, 4, and 5 for three other representative spectral ranges. We studied how the positions of specific bands change with different r values. A large set of bands as they result from calculations with different r values is given in Table V in Supplementary Material. In some cases, most notably in the C–H stretch spectral region, with decreasing r values some bands split in several ones due to resonances. This happens, for example, for the bands which are computed at 3034, 3037, and 3042 at the $r = 0.3$ level, which split into a multitude of bands at smaller r values. Some of these results are summarised in Fig. 6, which shows the ratio between anharmonic and harmonic frequencies of unambiguously identified fundamentals, for the different r values. In general, this exploration shows that $r = 0.3$ (and $h = 8$) already provides an acceptable level of accuracy for most bands, when compared with harmonic calculations. In almost all cases in which the same fundamental band can be unambiguously identified for all r values, the calculation for $r = 0.3$ already provides $\geq 90\%$ of the anharmonic correction to frequencies. Going from $r = 0.3$ to $r = 0.05$ only changes the positions of fairly strong bands (i. e. with peak intensities larger

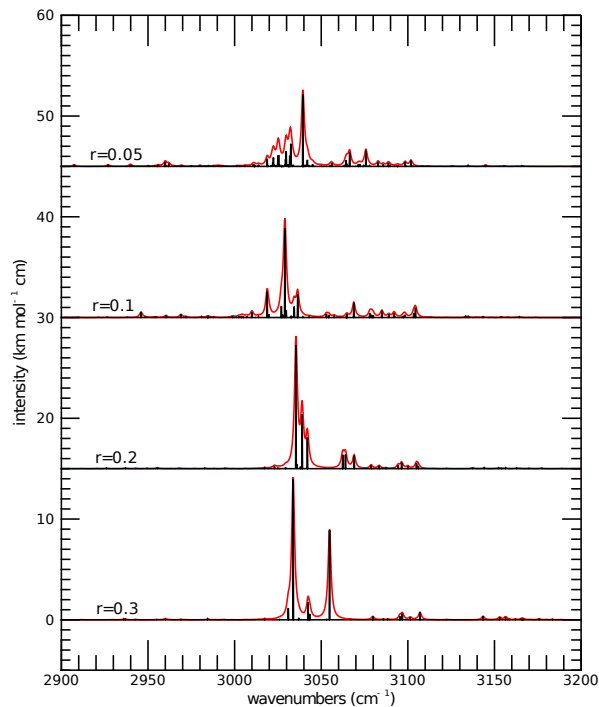


FIG. 2. Anharmonic spectrum of pyrene in the C-H stretch region for various values of the threshold r , keeping fixed $h = 8$. Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width. Spectra computed with different values of r are shifted for clarity by multiples of $15 \text{ km mol}^{-1} \text{ cm}$ with respect to the previous one.

than $\sim 1 \text{ km mol}^{-1} \text{ cm}$ when convolved with a 1 cm^{-1} wide Lorentzian) by less than $\sim 3\text{-}5 \text{ cm}^{-1}$, usually (but not always) slightly redwards. Only in very few particular cases, e. g. around 1430 cm^{-1} , when some resonance is taken into account due to decreasing r , a significant band may split in two close ones, but the overall spectral structure does not change much anyway except for the C-H stretch region. Very weak bands (i. e. bands with peak intensities $< 1 \text{ km mol}^{-1} \text{ cm}$ when convolved with a 1 cm^{-1} wide Lorentzian) are more sensitive. However, in most cases this is because some of them are “peripheric” states in big polyads, “borrowing” just very little intensity from fundamentals to which they are very indirectly connected by resonances. These states are indeed expected to be less well described, due to basis truncation errors.

Bands in the C-H stretches region, in contrast to other strong ones, appear very sensitive to thresholds and more difficult to get to converge. This is clearly due to the crowding of a much larger number of resonating states than the ones involved in transitions in other spectral re-

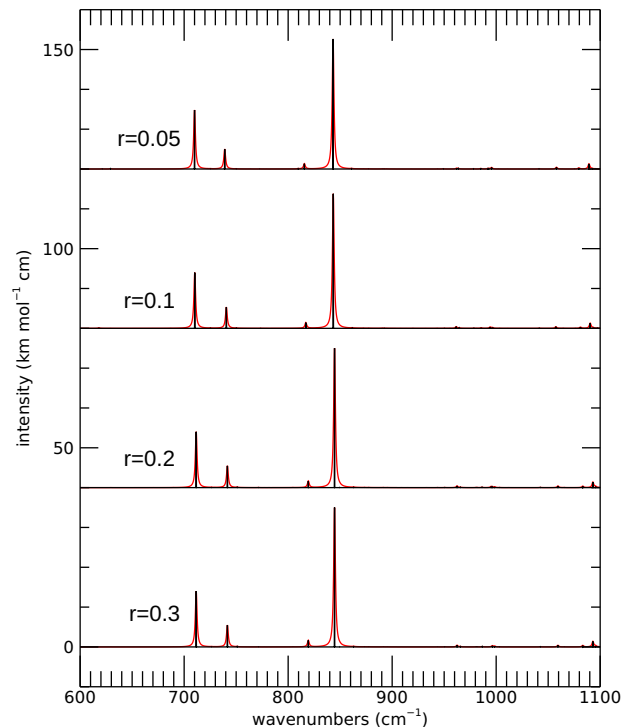


FIG. 3. Anharmonic spectrum of pyrene from 600 to 1100 cm^{-1} computed with various values of the threshold r , keeping fixed $h = 8$. Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width. Spectra computed with different values of r are shifted for clarity by multiples of $40 \text{ km mol}^{-1} \text{ cm}$.

gions. When more and more resonances are included, and larger polyads considered, the fundamental bands split in a few fairly strong and a multitude of weak bands. In some cases, this makes it actually impossible to “identify” a given fundamental in this spectral region with any specific anharmonic transition in the $r = 0.05$ calculation. Upon examining the spectra in Fig. 2, a side effect of this is also apparent: the total, integrated band intensity somewhat decreases when more and more resonances are treated explicitly. This happens due to the multitude of very weak Fermi resonances. Fermi resonances cause some transitions, which would be IR-inactive in the double harmonic approximation, to “borrow” some intensity from a fundamental transition. When such a Fermi resonance is treated explicitly, the total intensity is conserved, i. e. the “borrowed” intensity that appears from the band becoming active simultaneously disappears from the resonating fundamental. In contrast, when the same Fermi resonance is treated via a pertur-

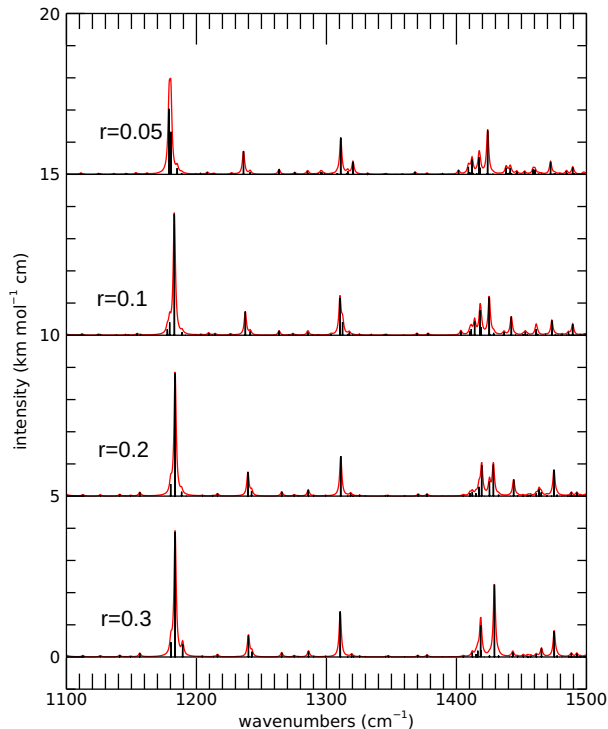


FIG. 4. Anharmonic spectrum of pyrene from 1100 to 1500 cm^{-1} computed with various values of the threshold r , keeping fixed $h = 8$. Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width. Spectra computed with different values of r are shifted for clarity by multiples of $5 \text{ km mol}^{-1} \text{ cm}$.

bative treatment, the “borrowed” intensity from a fundamental by a combination/difference band appears as a second order contribution, whereas the intensity decrease from the resonating fundamental appears as a third order one. Consequently, the latter is *not* included in the second order perturbative treatment. Some caution is therefore in order when considering band intensities obtained by these kinds of calculations, when a large number of weak resonances can add up and reach a total “borrowed” intensity which amounts to a sizeable fraction of the fundamental transition intensities. We actually implemented in our code an optional intensity correction which enforces the conservation of total intensity when computing the contribution of mechanical anharmonicity to the van Vleck transformed second derivatives of the electric dipole moment. This correction is not really well-balanced, since it approximately includes only one selected 3rd order correction to the transformed dipole moment operator, and not for example the one,

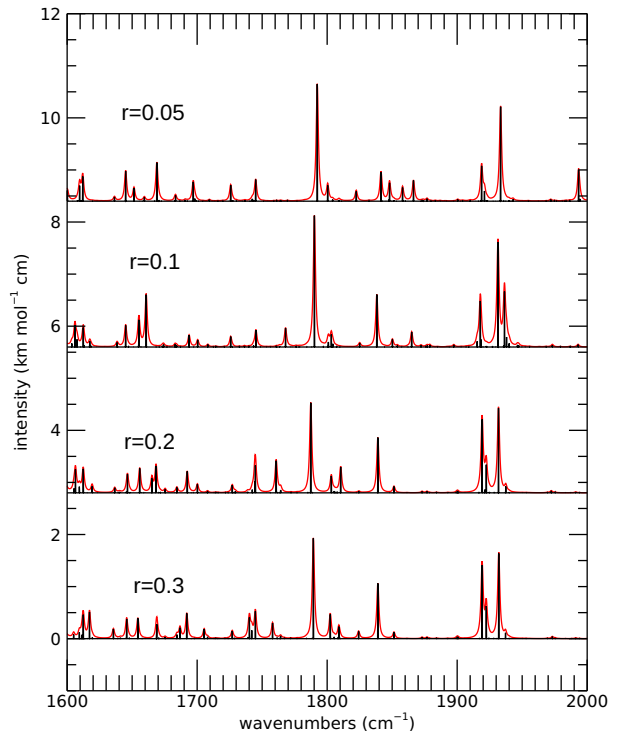


FIG. 5. Anharmonic spectrum of pyrene from 1600 to 2000 cm^{-1} computed with various values of the threshold r , keeping fixed $h = 8$. Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width. Spectra computed with different values of r are shifted for clarity by multiples of $2.8 \text{ km mol}^{-1} \text{ cm}$.

at the same order, transferring intensity in the opposite direction from an IR-active combination/difference band (produced by a sizeable second derivative of the dipole moment itself) to a fundamental resonating with it. This new functionality should be more fully tested under conditions in which other sources of error are negligible, hence it is disabled by default in AnharmoniCaOs, we just use it in the next section to get an estimate of how large its effect might be.

We carried out similar convergence tests for coronene. Its quartic force field and derivatives of the dipole moment were obtained using DFT with the B97-1⁵⁰ exchange-correlation functional. However, for coronene we used the 6-31G* Gaussian basis set to perform all the numerical derivatives required to obtain the quartic force field, since using the TZ2P⁵¹ basis-set would be computationally too expensive. For AnharmoniCaOs we used $r = 0.3$, $r = 0.2$, and $r = 0.1$, keeping fixed $h = 6$. Calculations with $r = 0.05$ and $h = 8$ were computation-

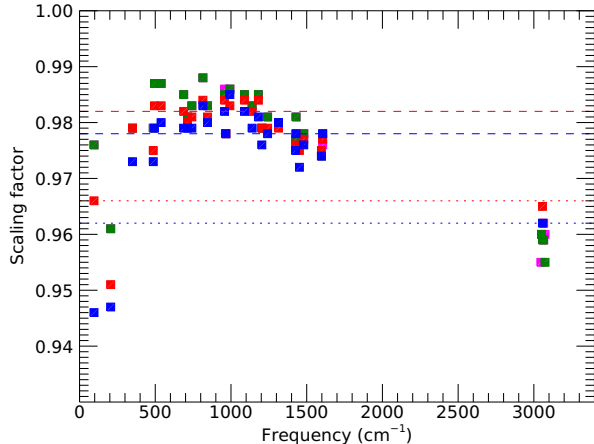


FIG. 6. Ratio between the computed anharmonic frequencies and the harmonic frequencies of unambiguously identified fundamentals of pyrene, as a function of the harmonic frequencies. Magenta squares are for $r = 0.3$, green ones for $r = 0.2$, red ones for $r = 0.1$, blue ones for $r = 0.05$. The overlaid lines represent the scaling factors for C-H stretches (dotted) and for all other bands (dashed). Red lines mark scaling factors from Ref. 56, namely 0.966 for C-H stretches and 0.982 for all the other bands. Blue lines mark scaling factors obtained from our best anharmonic calculation for pyrene, namely 0.962 for C-H stretches and 0.978 for all other bands except the two lowest modes, for which it is 0.946

ally too expensive. The results of this exploration for coronene were by and large the same as for pyrene, with the C-H stretch region being the only fairly sensitive one to thresholds and relatively difficult to converge.

C. Comparison with experimental results and previous theoretical calculations

Benchmarking our calculations on PAHs at 0 K would ideally require spectra recorded in gas-phase at very low temperatures. Such data are now becoming available for the CH stretch region.^{60,61} We therefore took the band positions reported for pyrene in ref. 61. Since this data is limited to the CH stretch range and does not report absolute intensities, we also used experimental data from Joblin *et al.*,⁶² both in Ne matrix at 4 K and in gas phase at relatively high temperatures (~ 300 -500 °C) due to the low vapour pressure of PAHs. Since the matrix effect on the band position is expected to be weak for neon¹³, we used the measured positions in Ne as the best available experimental data to be compared with our calculated positions. For most of the spectrum, the bands in Ne are very sharp and well resolved. We remark though that the band structure in the CH stretch range is less resolved in Ne matrices compared to gas-phase at low temperature, which can be observed from the values reported in Ta-

ble III. Again this indicates a strong role of anharmonicity in this range. Ne matrix spectra lack an absolute calibration of band intensities. In order to compare the experimental integrated band intensities with calculated band strengths, we therefore used the gas-phase spectra from Joblin *et al.*⁶² Error bars are not easy to derive but we proceeded as follows for the gas-phase data that are available to us.⁶³ A key point in deriving absolute intensities is to estimate the PAH column density. For gas-phase experiments in which all the PAH sample was evaporated at the temperature of interest, we were able to derive the column density of molecules in the gas-phase. If the evaporation is not total then this density is controlled by the vapour pressure, which is less precisely known. For coronene at 723 K we could use two spectra corresponding to total evaporation with different column densities. For pyrene at 523 K, we had only one such case and used a second case by applying an average global scaling on all bands. We deemed this useful to get an estimate of the relative errors from one experiment to the other but the absolute values might be less accurate in this case compared to coronene. In hot gas phase spectra it is often difficult to separate overlapping bands. After subtraction of a continuum from the spectra, we defined intervals to calculate band intensities. We then identified the corresponding bands in the theoretical spectra and summed them for comparison. The slight differences that can be found between the new gas-phase intensities and the ones previously reported⁶² is attributed to changes in the considered intervals and to the way the continuum level is subtracted.

1. Pyrene

The quartic force field and first and second derivatives of the electric dipole moment were obtained as described in Sections II C and III B. We then used the Anharmonic-CaOs code with the resulting quartic force field and second order Taylor expansion of the electric dipole, using the harmonic ground vibrational state as a starting state.

The results of our calculations can be compared with those obtained by Mackie *et al.*⁶⁴ and with the experimental data as described above. The comparison is shown in Table III for band positions and Table IV for band intensities. We remark that the calculated spectrum shown in Ref. 64 is nearly identical to the one we obtained with $r = 0.1$ and $h = 8$, shown in Fig. 2, which is not completely converged yet. We tentatively suppose that the smallest thresholds they tested for convergence, necessarily limited by computational constraints (the run including the largest polyads we tested required about 100 Gbytes of RAM) were smaller than the ones we arrived at. Since the computed spectrum changes abruptly every time the decrease of the threshold causes a significantly larger number of states to be included in the polyads, they probably did not reach threshold values low enough to obtain the top spectrum we show in Fig. 2.

TABLE III – continued from previous page

Gross pos. (μm)	Position (cm^{-1})			
	Ne Matrix 4 K		Theory 0 K	
	*Gas-phase (low T)			
	harm. scaled		anharmonic	
	(a)	(a)		(b)
5.9	1695		1697	$\nu_{35} + \nu_{68}$
6.0	1670, 1665		1669	$\nu_{36} + \nu_{69}$
	1648		1651, 1645	$\nu_{16} + \nu_{35}, \nu_{20} + \nu_{67}$
6.2	1617		1617, 1612, 1610	$\nu_{19} + \nu_{68}, \nu_{38} + \nu_{67}, \nu_{19} + \nu_{69}$
	1605	1606	1600, 1596	$\nu_{44}, \nu_{12} + \nu_{34}$
6.5		1597	1585, 1575	$\nu_{26}, \nu_{12} + \nu_{41} + \nu_{68}, \nu_{16} + \nu_{19}$
	1545		1537, 1544	$\nu_{39} + \nu_{66}, \nu_{20} + \nu_{68}$
	1517		1521, 1505	$\nu_{16} + \nu_{37}, \nu_{38} + \nu_{68}$
6.8	1496	1481	1473, 1489, 1485	$\nu_{45}, \nu_{14} + \nu_{21}, \nu_{21} + \nu_{66}$
			1460, 1459	$\nu_{15} + \nu_{39}, \nu_{51} + \nu_{64}$
7.0		1442	1442, 1438	$\nu_{36} + \nu_{70}, \nu_{27}$
	1438, 1436	1430	1424	ν_{28}
		1427	1418, 1417, 1412	$\nu_{34} + \nu_{63}, \nu_{46}, \nu_{10} + \nu_{53}$
7.7	1311	1314	1311, 1320, 1317	$\nu_{47}, \nu_{13} + \nu_{67}, \nu_{32} + \nu_{64}$
8.0	1244	1242	1241, 1236	$\nu_{40} + \nu_{68}, \nu_{29}$
8.4	1185, 1182	1181	1180	$\nu_{49}, \nu_{16} + \nu_{40},$
				$\nu_{39} + \nu_{41} + \nu_{53}, \nu_{39} + \nu_{41} + \nu_{53}$
9.1	1097, 1087	1089	1089, 1081, 1080	$\nu_{30}, \nu_{12} + \nu_{34}, \nu_{53} + \nu_{63}$
9.4	1063		1058	$\nu_{39} + \nu_{70}$
10.0	1004	993	996, 993	$\nu_{31}, \nu_{52} + \nu_{65}$
10.4	966	966	964, 963	$\nu_{66}, \nu_{41} + \nu_{69}$
11.8	844	845	843	ν_{67}
	822	815	816	ν_{32}
13.4	745	742	739	ν_{68}
14.0	713	713	710	ν_{69}
18.5	542	539	538	ν_{52}
20.0	499	495	493	ν_{34}
20.5	488	487	483	ν_{70}
29.0		350	347, 342	$\nu_{53}, \nu_{22} + \nu_{72}$
50		206	198	ν_{71}

To assess the overall accuracy of our calculations, we compare in Tables III and IV our results with the laboratory data that have been described above, as well as with the scaled harmonic calculation (see Supplementary material for all harmonic results and scaling factors). We also compared with the previously published anharmonic calculations by Mackie *et al.*⁶⁴ Since in this latter paper only relative intensities are given, we estimated absolute intensities for all possible bands by multiplying them by the absolute intensity we calculated for the band the authors chose for reference. This excludes C-H stretches

that the authors studied separately from the rest of the spectrum with an unknown scaling factor. Concerning band positions, we listed all states which carry a significant intensity and only one position was listed when these states are too close to be resolved. Assigning a given theoretical band to an experimental one can be tricky. We did our best on the basis of band proximity. This can be disputable especially in the case of the CH stretch region for which the calculations face some difficulties, as discussed above and in Sect. IV. For anharmonic calculations we label bands by the leading harmonic base state in

Gross pos. (μm)	Integration range (cm^{-1})	Intensities ($\text{km} \cdot \text{mol}^{-1}$)				
		Gas 570 K	Theory 0 K			
			harm. (a)	anharm.		(b)
				(a)	(b)	
				nc	c	
3.3	[2850-3250]	140(± 10)	95	99	74	
5.2	[1900-1950]	9.5 (± 0.5)		9.0	9.3	18.2
5.3	[1830-1880]	3.8 (± 0.1)		5.5	5.7	2.0
5.6	[1780-1830]	7.8 (± 0.7)		8.8	9.1	14.1
5.7	[1715-1760]	5.8 (± 0.2)		2.7	2.9	
5.9	[1677-1704]	—		1.9	2.0	
6.0	[1655-1677]	2.5 (± 0.2)		2.7	2.7	
6.1	[1630-1655]	—		2.8	2.9	
6.2	[1560-1620]	11.4 (± 1.2)	13.2	17.3	16.4	7.1
6.6	[1531-1552]	—		2.6	2.7	
6.8	[1467-1525]	—		3.7	3.5	3.0
7.0	[1405-1430]	11.4 (± 0.1)	11.2	8.8	7.8	6.1
7.7	[1300-1327]	2.4 (± 0.1)	4.9	5.2	4.5	6.1
8.0	[1224-1245]	1.9 (± 0.3)	2.3	2.6	2.4	2.0
8.4	[1165-1200]	10.5 (± 0.1)	14.0	13.7	12.1	14.1
9.1	[1074-1100]	5.4 (± 1.2)	7.0	5.4	4.7	8.1
10.0	[988-1003]	2.2 (± 0.6)		1.7	1.6	
11.8	[824-863]	100 (± 6)	112	101	93	101
13.4	[730-750]	20.8 (± 1.4)	17.5	15.2	14.0	24.2
14.0	[695-725]	46 (± 1)	44.8	45.3	42.9	41.4
18.5	[526-550]		2.5	2.3	2.1	
20.0	[487-501]			2.7	2.4	
20.5	[474-487]		2.1	1.9	1.7	
29.0	[335-354]			1.4	1.3	
50	[185-212]			9.7	9.0	

TABLE IV. List of the integrated intensities for pyrene. Theoretical data are (a) this work and (b) from Ref. 64. Since only relative intensities are reported in this latter study, we have scaled the values to the theoretical anharmonic intensity of the main CH out-of-plane bend mode. For our anharmonic results we list two columns, the “nc” one is standard GPVT2, “c” includes our correction borrowing some third order perturbation theory terms to enforce conservation of the intensity of transitions in resonances treated with perturbation theory. Experimental values have been derived from measurements in gas-phase at 570 K from two independent spectra taken from Refs. 62 and 63. The mean values are listed and the values in brackets provide the scatter around these values. Since in these spectra band overlapping is unavoidable, we selected integration ranges and integrated the corresponding theoretical bands. The listed ranges are the ones considered for the theoretical spectra and have been slightly shifted to cover the band envelopes in gas-phase spectra.

the expansion of the upper state of the transition. The harmonic modes are given in tables in the supplementary material. While in some cases anharmonic states are close enough to a single, well-defined harmonic state, in others they happen to be linear combinations of many harmonic states with coefficients of about the same magnitude, which can be as small as 10-20% for the leading harmonic state (which is nonetheless chosen to label the transition). Clear examples of both cases occur in the C-H stretch region. The ν_{42} harmonic state remains fairly well defined in the anharmonic calculation, with a large fraction of its intensity therefore in a single band at a computed position of 3040 cm^{-1} . The corresponding anharmonic state is a linear combination of the harmonic states ν_{42} (38%), $\nu_{45} + \nu_{48} + \nu_{53}$ (19%), $\nu_4 + \nu_{45}$ (4%), plus many other harmonic states contributing each less than 3%. Conversely, the ν_{43} and ν_{23} harmonic states

appear to distribute over a large number of anharmonic states via strong resonances. For example, the anharmonic state corresponding to the band at 3076 cm^{-1} , which we labelled ν_{23} , is a linear combination of the harmonic states ν_{23} (the “leading” one with 12%), $\nu_4 + \nu_{27}$ (11%), ν_{24} (9%), plus a large number of other harmonic states contributing less. The harmonic state ν_{43} , while contributing to many states, is not the leading one in any.

When compared to values in Ne matrices and excluding the CH stretch range, band positions appear to be accurate on average to better than 0.8%. The worst case is the band at $20.0 \mu\text{m}$ for which the shift reaches 1.2%. There is a very slight tendency to underestimate band positions, with an average ratio between theoretical and experimental positions of 0.998, which shows that systematic error is minimal. Of course Ne matrix positions will also be affected by some matrix shift effect, even if

it is expected to be small. Concerning integrated band intensities, we listed the integration limits we adopted to allow comparison with the gas phase spectra. We give here for pyrene the results both using our higher order correction enforcing conservation of band intensity and without it. The difference between the two theoretical sets of values are much smaller than the differences with experimental values, with the uncorrected ones appearing, if any, slightly better on average. For band intensities, the largest differences between calculated and experimental values appear in some of the weakest bands, the one at $5.7 \mu\text{m}$ being lower by a factor of two in calculations, while the one at $7.7 \mu\text{m}$ is larger by the same factor. The C-H stretch at $3.3 \mu\text{m}$, typically one of the most difficult bands to compute accurately with DFT, is lower in calculations by $\sim 40\%$, the remaining bands are within 25% . The fractional differences have both signs and average to about zero.

Our results are generally consistent with those published by Mackie *et al.*⁶⁴ for pyrene, even if a detailed one-to-one comparison depends on the possibility of unambiguously matching them based on what is given in Table 1 of their supplementary material. For this reason, not all theoretical band positions reported by Mackie *et al.*⁶⁴ are listed in last column of Table III, but only those whose identification could be matched with a corresponding band in our calculations. In particular, in the C-H stretch region we only included the three bands with the largest component of the three C-H stretch IR-active fundamentals. Differences are non-negligible nonetheless, as could be expected from using two different implementation of Van-Vleck theory (see Sect. IV) and could also

originate from small differences in the initial quartic force field. We remark that the best agreement with the low-temperature gas-phase data in the CH stretch range is found with our not completely converged calculation with $r = 0.1$, extremely similar to the one shown in Ref. 64. This hints that the accuracy we can achieve with this kind of calculation is, in this case, limited by the accuracy of the underlying DFT calculations yielding the quartic force field. The apparent slightly better accuracy of the $r = 0.1$ calculation with respect to the nominally better one with $r = 0.05$ is likely due to a partial accidental cancellation of errors. The complete, anharmonic spectrum of pyrene, computed with $r = 0.05$ and $h = 8$, is shown in the supplementary material, and available in tabulated form from the online database by Mallocci, Joblin, and Mulas.⁶⁵

2. Coronene

As for pyrene, the quartic force field and first and second derivatives of the electric dipole moment were obtained as described in Sections II C and III B. We then used the results obtained by our AnharmoniCaOs code with $r = 0.1$ and $h = 6$, leading to the largest polyads we could computationally handle, to obtain positions and intensities of the permitted transitions. We did not use our correction to enforce intensity conservation, since we saw in the case of pyrene that its effect is small compared to the errors arising from the underlying DFT calculations.

TABLE V: List of the main band positions for coronene. Theoretical data are from this work. The experimental band positions are those of the main bands in the Ne matrix spectrum recorded at 4 K taken from Ref. 62.

Gross pos. (μm)	Position (cm^{-1})			
	Ne Matrix 4 K	Theory 0 K		
		harm. scaled	anharmonic	
3.3	3070	3066	3121	$\nu_{30} + 2\nu_{61}, \nu_{30} + \nu_{61} + \nu_{62}, \nu_{30} + 2\nu_{62}$
			3120	$\nu_2 + \nu_{51}, \nu_2 + \nu_{52}$
			3112	$\nu_{30} + \nu_{71}, \nu_{30} + \nu_{71}$
			3071	ν_{45}, ν_{46}
			3056	$\nu_{10} + \nu_{51}, \nu_{10} + \nu_{52}, \nu_{21} + \nu_{71}, \nu_{21} + \nu_{72}$
	"3035"	3048	3049	$\nu_{21} + 2\nu_{61}, \nu_{21} + \nu_{61} + \nu_{62}, \nu_{21} + 2\nu_{62}$
			3050	$\nu_{49} + \nu_{75}, \nu_{49} + \nu_{76}, \nu_{50} + \nu_{75}, \nu_{50} + \nu_{76}$
			3038	$\nu_{10} + \nu_{51} \nu_{10} + \nu_{52}$
			3025	ν_{47}, ν_{48}
			1918, 1920	$\nu_{35} + \nu_{91}, \nu_{35} + \nu_{92}, \nu_{36} + \nu_{91}, \nu_{36} + \nu_{92}$
5.3	1926, 1913, 1898	1889, 1892, 1895	$\nu_7 + \nu_{35}, \nu_7 + \nu_{36}$	
		1886	$\nu_6 + \nu_{14}$	
5.6	1809, 1801, 1786	1783	$\nu_7 + \nu_{37}, \nu_7 + \nu_{38}$	

Continued on next page

TABLE V – continued from previous page

Gross pos. (μm)	Position (cm^{-1})			
	Ne Matrix 4 K		Theory 0 K	
		harm. scaled		anharmonic
5.9	1721		1708	$\nu_{14} + \nu_{37}, \nu_{14} + \nu_{38}$
	1697		1692	$\nu_{17} + \nu_{91}, \nu_{17} + \nu_{92}$
			1681	$\nu_{37} + \nu_{93}, \nu_{37} + \nu_{94}, \nu_{38} + \nu_{94}, \nu_{38} + \nu_{94}$
6.2	1621	1614	1619	ν_{49}, ν_{50}
7.6	1317	1309	1315	ν_{55}, ν_{56}
			1310	$\nu_{14} + \nu_{41}, \nu_{14} + \nu_{42}$
8.8	1139	1134	1143	ν_{59}, ν_{60}
11.7	857	857	862	ν_{14}
13.0	772	767	767	ν_{63}, ν_{64}
18.2	549	548	553	ν_{15}
26.5		376	379	ν_{65}, ν_{66}
81		121	126	ν_{16}

In Table V we compare our computed results for the band positions of coronene with laboratory data that were described above and with the scaled harmonic calculation, using the same approach as that described for pyrene. The band positions appear to be accurate to better than 0.7% except for the weak band at $\sim 5.3 \mu\text{m}$ for which the normalized shift between theory and experiment is 1.1%. As for pyrene, there is a balance between negative and positive values for this shift, since the average ratio between theoretical and Ne matrix positions is ~ 0.998 .

Theoretical band intensities appear somewhat less accurate than the ones for pyrene. The extreme cases are the bands at $5.3 \mu\text{m}$ and $13.0 \mu\text{m}$, which are both lower by a factor ~ 3 in the calculations compared to the experiment. The average difference is about 35%, with errors on both the positive and negative sides, with a prevalence of theoretical underestimation of experimental values. However, the accuracy is rather good for the most intense bands at 3.3 and $11.7 \mu\text{m}$. This is consistent with what we found in the case of pyrene. The slightly worse agreement with experimental data relative to the case of pyrene, is likely due to the smaller basis set used to obtain the quartic force field for coronene. The complete, anharmonic spectrum of coronene is shown in the supplementary material, and available in tabulated form from the online database by Mallocci, Joblin, and Mulas.⁶⁵

IV. DISCUSSION AND CONCLUSIONS

We performed anharmonic calculations of the vibrational spectra of neutral pyrene and coronene at 0 K using our AnharmoniCaOs code, and compared them to

the best available experimental data. The results are overall fairly satisfactory, yielding a significant improvement over the conventional double harmonic DFT calculations normally used for molecules of this size. As we already remarked at the beginning of Sect. III C, the most meaningful comparison should be with high resolution, low temperature gas phase measurements, which are only available in the C-H stretch region for a small set of PAHs.^{60,61} For the present work, we mostly relied on Ne matrix spectra for band positions and high temperature gas-phase spectra for band intensities.^{62,63} Our conclusions might slightly change if/when full high resolution, low temperature gas phase spectra of such species become available.

The accuracy of our anharmonic calculations is clearly related to several different independent limiting factors, respectively:

1. The representation of the true adiabatic potential energy surface as a quartic force field, and the dipole moment as a Taylor expansion truncated to second order, both in terms of cartesian normal coordinates ;
2. The accuracy of the determination of the quartic force field and quadratic dipole parameters via quantum chemistry calculations;
3. The accuracy of the GVPT2 method itself;
4. The accuracy of our AnharmoniCaOs implementation (as a function of its parameters).

Point 1 is an acceptable approximation when oscillations have a relatively small amplitude. Therefore, it is suitable for semirigid molecules at not too high vibrational

Gross pos. (μm)	Exp. Range (cm^{-1})	Intensities ($\text{km} \cdot \text{mol}^{-1}$)		
		Gas 570 K	Theory 0 K	
			harm.	anharm.
3.3	[2995-3110]	190 (1)	134.7	177
5.3	[1870-1906]	23 (3)		31.5
5.6	[1770-1825]	18 (4)		5.9
5.9	[1670-1720]	20 (3)		13.9
6.2	[1600-1629]	19 (2)	15.2	24.9
7.6	[1300-1330]	37 (2)	42.3	22.7
8.8	[1128-1158]	22 (4)	17.3	14.4
11.7	[845-880]	138 (s)	167.9	147
13.0	[745-785]	32 (1)	13.6	13.2
18.2	[530-565]	44 (2)	42.7	23.0
26.5	[360-395]		7.7	5.9
81	[110-140]		6.9	4.6

TABLE VI. List of the integrated intensities for coronene. Theoretical data are from this work. Experimental values have been derived from measurements in gas-phase at 570 K from two independent spectra taken from Refs. 62 and 63. The mean values are listed and the values in brackets provide the scatter around these values. The strongest band at $11.7 \mu\text{m}$ (annotation (s)) was saturated in one of the two experimental spectra so only a single value could be reported. Since in these spectra band overlapping is unavoidable, we selected integration ranges and integrated the corresponding theoretical bands. The listed ranges are the ones considered for the theoretical spectra and have been slightly shifted to cover the band envelopes in gas-phase spectra.

energies. So, it should be appropriate for PAHs, unless they have side groups that give rise to internal rotations, that are poorly represented by a truncated Taylor expansion in normal coordinates.⁶⁶ We notice that it would be desirable in some cases, to opt for non-cartesian coordinates so that the associated quartic force field would have a better asymptotic behaviour, and the potential would be reasonably described even along large amplitude motions.^{67,68} Such coordinates, however, are not directly applicable in our case: the Van Vleck approach used here for GVPT2 requires a description of the Hamiltonian in terms of normal coordinates, since it exploits their commutation properties to obtain a formally simple representation of the infinitesimal contact transformations and of the transformed Hamiltonian and dipole moment operators.

Point 2 is basically a matter of trade-off between computational cost and accuracy. Band positions can be fairly accurate for PAHs, as shown also by our results, even using moderate size basis sets (such as the triple zeta ones we used) with DFT using hybrid exchange-correlation functionals.⁶⁹ This can be further improved, approaching spectroscopic accuracy, by using more accurate (and computationally way more expensive) methods with larger basis sets, either only for the harmonic frequencies or also for some of the higher derivatives of the potential, and for the first and/or second derivatives of the dipole moment (see Ref. 70 and references therein). Absolute intensities are more difficult to get with the same accuracy as for positions using DFT calculations, but they also become fairly precise when using high (computationally expensive) levels of theory.⁷¹

As to point 3, GVPT2 has been used successfully for many years (see e.g. Ref. 29), and it has been reviewed

by Bloino *et al.* in Ref. 70. In general, GVPT2 should be appropriate as long as point 1 is valid.

Finally point 4 has been assessed in our comparisons with higher level vibrational calculations, in Sect. III A. Our code has the limitation that rotational degrees of freedom, and their interaction terms with vibrational ones (i.e. terms due to the Coriolis and centrifugal pseudo forces) are not included in our approximated Hamiltonian. This is not expected to be an important limit for relatively large, semi-rigid, asymmetric top molecules such as e. g. pyrene. However, it is always possible that a small number of individual states undergo accidental Coriolis resonances. In contrast, strong Coriolis interactions are bound to occur for vibrational states which are degenerate due to symmetry reasons, and this will be more likely for symmetric species like coronene. However, we remark that the agreement between theoretical and experimental data for coronene does not seem significantly worse than that of pyrene, hinting that Coriolis coupling is not crucial.

Summing up, we find that for the cases of pyrene and coronene presented here the main limitation appears to be point 2, since we clearly reached a point at which increasing the accuracy of our GPVT2 calculation did not improve the agreement with experimental data. Indeed, it looks like there is a “sweet spot” in the accuracy of the GPVT2 calculations where some accidental cancellation of errors with the underlying DFT-based quartic force field produces the best results, even if this may not be general.

The calculations presented here are similar to those performed by Mackie *et al.* in Ref. 64. The latter authors reported an agreement of calculated band positions with experimental data (in rare gas matrices) of $0.4 \pm 0.6\%$, to

be compared with our value of better than 0.8%. Since both studies include different sets of laboratory data, it is difficult to conclude which one provides the best results. Still, it can be noticed that some individual bands positions differ by more than numerical noise, when comparing both theoretical studies (cf. Table III). The implementation of GVTP2 in Ref. 64 is different from ours. In particular, a different strategy is used to define which cubic and quartic terms of the potential should be included in the treatment of resonances rather than in the perturbation expansion, and subsequently the construction (and truncation) of polyads differ. In spite of these differences, the overall agreement can be seen as a validation of both approaches. With the parameters used, the results reported here appear slightly more accurate for the region of combination bands, whereas the results in Ref. 64 appear somewhat better in the C-H stretches region. We remark once more that we actually did obtain very nearly the same results as Mackie *et al.*⁶⁴ for the C-H stretches of pyrene with our not completely converged anharmonic calculation, as discussed in Sect. III C 1 and recalled in the previous paragraph, where we hinted that this can be due to an accidental cancellation of errors. The availability of two independent codes performing similar calculations is of course important, as this can be used to test both. Indeed, the close similarity of the results, at least with some choice of the tunable parameters, validates both codes and the correctness of their results. In addition we remark that AnharmoniCaOs is freely distributed under an open source license via SourceForge, guaranteeing it will remain available and easy to find and download for the foreseeable future.

Scaled harmonic calculations of fundamentals appear to be not much less accurate than our anharmonic ones. However, they depend on the use of empirically calibrated parameters. Moreover, harmonic calculations are limited only to fundamental transitions and completely neglect combination, difference and overtone bands, some of which may be of sizeable intensity especially when strong Fermi resonances occur, as shown also in Refs. 34, 64 and 61. Indeed, Fig. 2, and Table V in Supplementary Material, clearly show that some bands, most notably (but not only) in the C-H stretch region, split from their simplistic harmonic structure into several, sometimes a multitude of, close ones, due to resonances. Such a complex structure was confirmed for the C-H stretch region by the low temperature, gas-phase measurements of Refs. 34, 64 and 61. Only anharmonic calculations can account for this. We also note that as a by-product we can also use this kind of anharmonic calculation to estimate purely theoretical frequency scaling factors. Fig. 6 shows the empirical scaling factors for the level of theory B97-1/TZ2P,^{56,60,61} as well as the ones we can derive from our calculations on pyrene, which are consistent with them.

The main limiting factor to this kind of anharmonic calculations, at least for species of up to a few tens of atoms in size, is related to the computational cost of

computing the cubic and quartic force field constants, as well as the second derivatives of the electric dipole moment, which must be obtained via numerical differentiation since currently there is no readily available code that can compute them analytically. However, some general schemes have been proposed⁷² to obtain such higher order derivatives in the framework of the Density Functional Theory, even if their computational implementation is still a work in progress and is still private. When such codes will become available and well-tested, anharmonic calculations, at least of the ground vibrational state and of the states connected to it by permitted transitions, are likely to become commonplace. The AnharmoniCaOs code, besides being usable to process quartic force fields and first and second order electric dipole moment derivatives to obtain spectra from the vibrational ground state, can be used for vibrationally excited states as well, and is currently being tested to obtain moderately high temperature spectra of the same PAHs studied here.

SUPPLEMENTARY MATERIAL

See supplementary material for:

- A description of the different ethylene oxide force fields that have been used for benchmarking purposes.
- A table with the comparison of the VMFCI reference calculation with HI-RRBPM calculation D of Ref. 21 for the force field adimensionned with DFT quadratic force constants;
- A table with the comparison of the VMFCI reference calculation with a previous VMFCI calculation using a different contraction-truncation scheme for the force field adimensionned with coupled-cluster quadratic force constants;
- A table with the list of harmonic vibrational normal modes of pyrene;
- A table with the list of harmonic vibrational normal modes of coronene;
- A table listing the (significantly IR-active) anharmonic states of pyrene as a function of r values, for fixed $h = 8$;
- Figures of the complete anharmonic spectra of pyrene and coronene.

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Anharmonic vibrational spectroscopy of Polycyclic Aromatic Hydrocarbons (PAHs)

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This supplementary material contains:

- A description of the different ethylene oxide force fields that have been used for benchmarking purposes
- Table I: Comparison of the VMFCI reference calculation with HI-RRBPM calculation D of Ref. 1 for the force field adimensionned with DFT quadratic force constants
- Table II: Comparison of the VMFCI reference calculation with a previous VMFCI calculation using a different contraction-truncation scheme for the force field adimensionned with coupled-cluster quadratic force constants
- Table III: List of harmonic vibrational normal modes of pyrene.
- Table IV: List of harmonic vibrational normal modes of coronene
- Table V: List of (significantly IR-active) anharmonic states of pyrene as a function of r values, for fixed $h = 8$
- Figures 1–27: Anharmonic spectra of pyrene and coronene

I. VMFCI CALCULATIONS

A. Different ethylene oxide force fields

The ethylene oxide quartic force field of Bégue *et al.*² has now become a classical benchmark for methods solving the anharmonic vibrational Schrödinger equation. However, there are three different versions of this force field:

1. The version provided as supplementary material of Ref. 2, which was actually only employed for the largest P-VMWCI calculation. It lacks one cubic and seven quartic force constants that were removed from the original DFT quartic force field. This was necessary to alleviate the computational effort required to perform the $\gamma = 3$ P-VMWCI calculation of Ref. 2.
2. The version used in the main text of present work and in the VMFCI calculations of Refs. 1 and 2. It is the combination of the CCSD(T)/cc-pVTZ quadratic constants and the full set of B3LYP/6-31+G(d,p) cubic and quartic force constants in mass-weighted, Cartesian, normal coordinates. In terms of adimensional coordinates, the

force field obtained amounts to cubic and quartic force constants rescaled by using CCSD(T)/cc-pVTZ quadratic constants. It is this force field that has been used for benchmarking AnharmoniCaOs.

3. The version used in recent investigations.^{1,3,4} It differs from the previous one by the fact that the cubic and quartic force constants were adimensionned by using the B3LYP/6-31+G(d,p) quadratic constants. This results in a noticeably different force field: for instance, the zero point energy with this force field differs by 2 cm^{-1} from that obtained with the previous force field, (compare the two tables below).

B. Comparison of the new VMFCI reference calculation with HI-RRBPM results

Two fairly large numerical calculations of the 200 lowest eigenvalues of the vibrational Hamiltonian associated to force field 3 have appeared recently. One is based on the Adaptive Vibrational Configuration Interaction (A-VCI) algorithm,⁴ and the other on the Hierarchical Intertwined Reduced-Rank Block Power Method (HI-RRBPM).¹ We postpone a discussion of the remarkable A-VCI results to a forthcoming publication. Here, we just validate a new Vibrational Mean Field Configuration Interaction (VMFCI) contraction-truncation scheme

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against the latest HI-RRBPM calculation (calculation “D” of Ref. 1).

The new contraction-truncation scheme is the following:

- The calculation starts with the same modal basis set of 100 harmonic function on each of the 15 modes, then a VSCF calculation is converged to machine precision in seven VMFCI iterations. The common zero point energy (ZPE) of the 15 modes is then 12539.269841 cm^{-1} .
- Then, we make 6 contractions: the four C–H stretching modes $\nu_1 - \nu_6 - \nu_9 - \nu_{13}$, the two H–C–H scissoring modes $\nu_2 - \nu_{10}$, the three ring breathing and deformation modes $\nu_3 - \nu_5 - \nu_{12}$, the two wagging modes $\nu_4 - \nu_{11}$, the two twisting modes $\nu_7 - \nu_{14}$, the two rocking modes $\nu_8 - \nu_{15}$, with truncation thresholds on the sum of the energies of 34072 cm^{-1} for the stretchings, 17000 cm^{-1} for the ring deformations and 30000 cm^{-1} for the four other contractions. This partition is iterated once and the 6 ZPE’s decrease down to $12501.358 \pm 0.001 \text{ cm}^{-1}$. All this is abbreviated by VSCFCI($\{\nu_1, \nu_6, \nu_9, \nu_{13}\}$, 34072; $\{\nu_2, \nu_{10}\}$, 30000; $\{\nu_3, \nu_5, \nu_{12}\}$, 17000; $\{\nu_4, \nu_{11}\}$, 30000; $\{\nu_7, \nu_{14}\}$, 30000; $\{\nu_8, \nu_{15}\}$, 30000).
- Next, we make a coarser partition of only two contractions and play on both individual energy and sum of energy thresholds in a single VMFCI step: VMFCI($\{\{\nu_1, \nu_6, \nu_9, \nu_{13}\}_{17523}, \{\nu_2, \nu_{10}\}_{12272}, \{\nu_8, \nu_{15}\}_{12077}\}$, 17914; $\{\{\nu_3, \nu_5, \nu_{12}\}_{11459}, \{\nu_4, \nu_{11}\}_{11156}, \{\nu_7, \nu_{14}\}_{9633}\}$, 15280), (the subscripts are the thresholds in cm^{-1} on the individual energies of the component of the new contractions).
- Finally, the two previous contraction are contracted in a VCI step with a truncation threshold of 6367.5 cm^{-1} on the second component and a threshold of 11956.6 cm^{-1} on the sum of the component energies. The size of this VCI is 430769 basis functions, it splits into 4 C_{2v} -symmetry blocks of dimensions 108147 for A1, 107421 for A2, 107740 for B1, 107461 for B2.

The whole calculation is denoted VMFCI(11956.6) in the table below and the wave numbers are compared to those of calculation D of Ref. 1.

TABLE I: Lowest 200 vibrational energy levels (in cm^{-1}) of ethylene oxide for force field 3

irreps.	VMFCI(11956.6)	Ref. 1	Assignment ¹
A ₁	12461.57	12461.65	ZPE
B ₂	792.80	793.10	ν_{15}
B ₁	821.98	822.00	ν_{12}
A ₁	878.32	878.33	ν_5
A ₂	1017.31	1017.49	ν_8
A ₁	1121.64	1121.94	ν_4
B ₁	1124.07	1124.37	ν_{11}
B ₂	1146.00	1146.19	ν_{14}
A ₂	1148.22	1148.40	ν_7
A ₁	1270.95	1270.94	ν_3

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irreps.	VMFCI(11956.6)	Ref. 1	Assignment ¹
B ₁	1467.59	1467.72	ν_{10}
A ₁	1495.49	1495.74	ν_2
A ₁	1588.03	1589.09	$2\nu_{15}$
A ₂	1611.50	1611.63	$\nu_{15} + \nu_{12}$
A ₁	1641.39	1641.21	$2\nu_{12}$
B ₂	1670.69	1670.78	$\nu_{15} + \nu_5$
B ₁	1695.36	1695.13	$\nu_5 + \nu_{12}$
A ₁	1755.07	1754.81	$2\nu_5$
B ₁	1806.39	1806.81	$\nu_{15} + \nu_8$
B ₂	1832.96	1832.89	$\nu_{12} + \nu_8$
A ₂	1889.73	1889.81	$\nu_5 + \nu_8$
A ₂	1908.54	1908.90	$\nu_{15} + \nu_{11}$
B ₂	1910.58	1910.80	$\nu_{15} + \nu_4$
B ₁	1928.71	1929.15	$\nu_{15} + \nu_7$
A ₁	1938.65	1938.92	$\nu_{15} + \nu_{14}$
B ₁	1940.37	1940.58	$\nu_{12} + \nu_4$
A ₁	1944.67	1944.71	$\nu_{12} + \nu_{11}$
A ₂	1962.19	1962.04	$\nu_{12} + \nu_{14}$
B ₂	1970.84	1970.60	$\nu_{12} + \nu_7$
A ₁	1998.09	1997.73	$\nu_5 + \nu_4$
B ₁	1999.74	1999.39	$\nu_5 + \nu_{11}$
B ₂	2021.35	2021.12	$\nu_5 + \nu_{14}$
A ₂	2024.24	2024.00	$\nu_5 + \nu_7$
A ₁	2032.16	2032.78	$2\nu_8$
B ₂	2060.62	2060.74	$\nu_{15} + \nu_3$
B ₁	2086.60	2086.38	$\nu_{12} + \nu_3$
A ₂	2128.92	2128.77	$\nu_8 + \nu_4$
B ₂	2135.29	2135.21	$\nu_8 + \nu_{11}$
A ₁	2140.80	2140.60	$\nu_5 + \nu_3$
B ₁	2152.76	2152.89	$\nu_8 + \nu_{14}$
A ₁	2165.51	2165.40	$\nu_8 + \nu_7$
A ₁	2236.23	2236.66	$2\nu_4$
B ₁	2242.75	2243.99	$\nu_4 + \nu_{11}$
A ₂	2246.58	2247.45	$\nu_{15} + \nu_{10}$
A ₁	2249.90	2250.81	$2\nu_{11}$
B ₂	2264.70	2266.13	$\nu_7 + \nu_{11}$
A ₂	2267.87	2268.95	$\nu_{14} + \nu_{11}$
B ₂	2268.07	2268.67	$\nu_{14} + \nu_4$
A ₂	2275.16	2275.77	$\nu_7 + \nu_4$
B ₂	2282.99	2284.58	$\nu_{15} + \nu_2$
A ₁	2288.29	2288.50	$\nu_{12} + \nu_{10}$
A ₂	2289.88	2290.67	$\nu_3 + \nu_8$
A ₁	2290.79	2291.41	$2\nu_{14}$
B ₁	2293.56	2294.02	$\nu_7 + \nu_{14}$
A ₁	2296.10	2296.40	$2\nu_7$
B ₁	2310.81	2311.26	$\nu_{12} + \nu_2$
B ₁	2345.42	2345.49	$\nu_5 + \nu_{10}$
A ₁	2370.35	2370.74	$\nu_5 + \nu_2$
B ₂	2381.33	2388.00	$3\nu_{15}$
A ₁	2389.76	2389.51	$\nu_3 + \nu_4$
B ₁	2391.37	2391.01	$\nu_3 + \nu_{11}$
B ₁	2402.50	2405.27	$2\nu_{15} + \nu_{12}$
B ₂	2415.20	2415.12	$\nu_3 + \nu_{14}$
A ₂	2415.99	2415.73	$\nu_3 + \nu_7$
B ₂	2428.33	2429.22	$\nu_{15} + 2\nu_{12}$
B ₁	2458.05	2458.83	$3\nu_{12}$
A ₁	2465.15	2468.36	$2\nu_{15} + \nu_5$

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irreps.	VMFCI(11956.6)	Ref. 1	Assignment ¹
B_2	2481.23	2482.05	$\nu_{10} + \nu_8$
A_2	2484.11	2484.74	$\nu_{15} + \nu_5 + \nu_{12}$
A_2	2506.69	2507.98	$\nu_2 + \nu_8$
A_1	2510.12	2513.46	$\nu_5 + 2\nu_{12}$
A_1	2537.36	2537.41	$2\nu_3$
B_2	2546.70	2547.55	$\nu_{15} + 2\nu_5$
B_1	2567.16	2567.33	$2\nu_5 + \nu_{12}$
B_1	2583.86	2584.32	$\nu_{10} + \nu_4$
A_1	2587.72	2588.24	$\nu_{10} + \nu_{11}$
A_2	2595.27	2597.09	$2\nu_{15} + \nu_8$
A_2	2599.54	2600.34	$\nu_{10} + \nu_{14}$
B_2	2601.60	2602.28	$\nu_{10} + \nu_7$
A_1	2612.35	2613.87	$\nu_2 + \nu_4$
B_1	2616.37	2617.02	$\nu_2 + \nu_{11}$
A_1	2616.89	2618.09	$\nu_{15} + \nu_{12} + \nu_8$
A_1	2629.75	2630.88	$3\nu_5$
A_2	2640.84	2641.75	$\nu_2 + \nu_7$
B_2	2641.13	2641.91	$\nu_2 + \nu_{14}$
A_2	2646.04	2646.24	$2\nu_{12} + \nu_8$
B_1	2676.07	2677.43	$\nu_{15} + \nu_5 + \nu_8$
B_1	2693.79	2696.02	$2\nu_{15} + \nu_{11}$
A_1	<i>2699.42</i>	<i>2701.69</i>	$2\nu_{15} + \nu_4$
B_2	<i>2700.39</i>	<i>2701.38</i>	$\nu_5 + \nu_{12} + \nu_8$
A_2	2711.51	2713.94	$2\nu_{15} + \nu_7$
B_2	2721.54	2723.68	$\nu_{15} + \nu_{12} + \nu_{11}$
A_2	2728.32	2730.67	$\nu_{15} + \nu_{12} + \nu_4$
B_2	2734.81	2736.72	$2\nu_{15} + \nu_{14}$
B_1	2736.84	2737.37	$\nu_{10} + \nu_3$
A_1	2743.92	2744.96	$\nu_{15} + \nu_{12} + \nu_7$
B_1	2751.45	2752.71	$\nu_{15} + \nu_{12} + \nu_{14}$
A_1	2757.10	2758.29	$2\nu_{12} + \nu_4$
A_2	2760.48	2761.35	$2\nu_5 + \nu_8$
B_1	2760.82	2761.37	$2\nu_{12} + \nu_{11}$
A_1	2762.81	2763.88	$\nu_2 + \nu_3$
B_2	2775.18	2775.59	$2\nu_{12} + \nu_{14}$
A_2	2784.06	2787.72	$\nu_{15} + \nu_5 + \nu_{11}$
B_2	2785.74	2788.41	$\nu_{15} + \nu_5 + \nu_4$
A_2	2789.89	2790.40	$2\nu_{12} + \nu_7$
B_1	2803.09	2809.98	$\nu_{15} + \nu_5 + \nu_7$
A_1	2811.45	2813.50	$\nu_5 + \nu_{12} + \nu_{11}$
A_1	2812.63	2817.12	$\nu_{15} + \nu_5 + \nu_{14}$
B_2	2815.48	2817.46	$\nu_{15} + 2\nu_8$
B_1	2816.02	2817.70	$\nu_5 + \nu_{12} + \nu_4$
A_2	2832.39	2834.71	$\nu_5 + \nu_{12} + \nu_{14}$
B_1	<i>2840.51</i>	<i>2842.50</i>	$\nu_{12} + 2\nu_8$
B_2	<i>2841.30</i>	<i>2842.15</i>	$\nu_5 + \nu_{12} + \nu_7$
A_1	2851.94	2853.47	$2\nu_{15} + \nu_3$
A_1	2872.12	2873.05	$2\nu_5 + \nu_4$
A_2	<i>2872.23</i>	<i>2874.54</i>	$\nu_{15} + \nu_{12} + \nu_3$
B_1	<i>2872.96</i>	<i>2873.88</i>	$2\nu_5 + \nu_{11}$
B_2	2894.46	2895.89	$2\nu_5 + \nu_{14}$
A_2	2897.94	2898.62	$2\nu_5 + \nu_7$
A_1	2898.22	2901.60	$\nu_5 + 2\nu_8$
A_1	2900.68	2902.18	$2\nu_{12} + \nu_3$
B_1	2907.72	2911.04	ν_9
A_1	2913.03	2917.53	$\nu_{15} + \nu_8 + \nu_{11}$ ^a
B_1	2915.00	2917.90	$\nu_{15} + \nu_8 + \nu_4$

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irreps.	VMFCI(11956.6)	Ref. 1	Assignment ¹
A_1	2918.79	2923.02	ν_1
B_2	2929.63	2932.84	$\nu_{15} + \nu_5 + \nu_3$
B_2	2934.76	2937.33	$\nu_{12} + \nu_8 + \nu_4$
A_2	2940.11	2942.52	$\nu_{15} + \nu_8 + \nu_{14}$
B_2	2946.04	2949.06	$\nu_{15} + \nu_8 + \nu_7$
A_2	2949.03	2951.56	$\nu_{12} + \nu_8 + \nu_{11}$
B_1	2952.56	2954.44	$\nu_5 + \nu_{12} + \nu_3$
A_1	2953.40	2958.32	$2\nu_{10}$
A_1	2965.35	2967.45	$\nu_{12} + \nu_8 + \nu_{14}$
B_1	2980.82	2981.94	$\nu_{12} + \nu_8 + \nu_7$
B_1	2990.72	2995.63	$\nu_2 + \nu_{10}$
A_1	<i>2995.54</i>	3002.34	$2\nu_2$
A_2	<i>2999.17</i>	<i>3000.44</i>	$\nu_5 + \nu_8 + \nu_4$
B_2	3003.31	3006.30	$\nu_5 + \nu_8 + \nu_{11}$
A_1	3009.07	3009.47	$2\nu_5 + \nu_3$
B_2	3013.56	3020.05	$\nu_{15} + 2\nu_{11}$ ^{c1}
A_2	<i>3021.24</i>	<i>3025.44</i>	$\nu_{15} + \nu_4 + \nu_{11}$
B_1	<i>3023.14</i>	<i>3024.79</i>	$\nu_5 + \nu_8 + \nu_{14}$
A_2	3026.36	3027.95	ν_6
B_2	3030.94	3033.86	$\nu_{15} + 2\nu_4$ ^{c1}
B_1	<i>3034.01</i>	<i>3038.36</i>	$2\nu_{15} + \nu_{10}$
A_1	<i>3034.03</i>	<i>3035.70</i>	$\nu_5 + \nu_8 + \nu_7$
B_2	3038.39	3040.21	ν_{13}
A_2	<i>3041.93</i>	3048.87	$3\nu_8$
A_1	3042.71	3048.70	$\nu_{15} + \nu_7 + \nu_{11}$
B_1	<i>3043.77</i>	<i>3048.42</i>	$\nu_{12} + 2\nu_{11}$ ^b
B_1	3050.73	3057.21	$\nu_{15} + \nu_{14} + \nu_{11}$
A_1	3053.21	3060.39	$\nu_{12} + \nu_4 + \nu_{11}$
A_1	3056.21	3061.69	$\nu_{15} + \nu_{14} + \nu_4$
B_1	3060.53	3066.08	$\nu_{15} + \nu_7 + \nu_4$
B_2	3063.15	3066.94	$\nu_{15} + \nu_{12} + \nu_{10}$
A_1	<i>3063.78</i>	3071.18	$2\nu_{15} + \nu_2$
B_1	<i>3065.83</i>	<i>3069.95</i>	$\nu_{12} + 2\nu_4$ ^b
A_2	<i>3069.09</i>	3074.73	$\nu_{12} + \nu_{14} + \nu_4$ ^d
B_2	<i>3069.29</i>	<i>3072.99</i>	$\nu_{15} + 2\nu_{14}$ ^{c2}
B_1	3074.08	3079.51	$\nu_{15} + \nu_3 + \nu_8$
A_2	<i>3079.24</i>	3085.65	$\nu_{15} + \nu_{12} + \nu_2$
B_2	<i>3080.70</i>	<i>3084.54</i>	$\nu_{12} + \nu_{14} + \nu_{11}$
A_2	3083.97	3088.11	$\nu_{15} + \nu_7 + \nu_{14}$
B_2	3085.07	3088.58	$\nu_{15} + 2\nu_7$ ^{c2}
B_2	3092.78	3098.61	$\nu_{12} + \nu_7 + \nu_4$
A_2	3099.06	3102.95	$\nu_{12} + \nu_7 + \nu_{11}$
B_2	3100.52	3103.16	$\nu_{12} + \nu_3 + \nu_8$
B_1	3100.73	3104.30	$\nu_{12} + 2\nu_{14}$
B_1	3106.74	3109.15	$2\nu_{12} + \nu_{10}$
A_1	<i>3109.44</i>	3115.58	$\nu_{12} + \nu_7 + \nu_{14}$ ^e
A_1	<i>3109.53</i>	3125.67	$\nu_5 + 2\nu_4$ ^e
B_1	3114.87	3119.83	$\nu_5 + \nu_4 + \nu_{11}$
B_1	3118.20	3121.56	$\nu_{12} + 2\nu_7$
A_1	<i>3121.77</i>	3112.80	$\nu_5 + 2\nu_{11}$ ^e
A_2	3122.74	3125.96	$\nu_{15} + \nu_5 + \nu_{10}$
A_1	3125.17	3127.61	$2\nu_{12} + \nu_2$
A_1	3131.40	3134.62	$2\nu_8 + \nu_4$
A_2	<i>3137.83</i>	3143.98	$\nu_5 + \nu_{14} + \nu_{11}$
B_2	<i>3138.43</i>	<i>3141.97</i>	$\nu_5 + \nu_7 + \nu_{11}$
B_2	<i>3140.71</i>	<i>3142.77</i>	$\nu_5 + \nu_{14} + \nu_4$
B_1	3141.72	3146.30	$2\nu_8 + \nu_{11}$

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irreps.	VMFCI(11956.6)	Ref. 1	Assignment ¹
<i>A</i> ₂	3148.76	3154.15	$\nu_5 + \nu_7 + \nu_4$
<i>A</i> ₂	3155.47	3160.67	$\nu_5 + \nu_3 + \nu_8$
<i>B</i> ₂	3156.27	3162.47	$\nu_{15} + \nu_5 + \nu_2$
<i>B</i> ₂	3156.67	3163.50	$2\nu_8 + \nu_{14}$
<i>A</i> ₁	3160.98	3163.70	$\nu_5 + \nu_{12} + \nu_{10}$
<i>A</i> ₁	3162.51	3164.77	$\nu_5 + 2\nu_{14}$
<i>B</i> ₁	3165.89	3169.44	$\nu_5 + \nu_7 + \nu_{14}$
<i>A</i> ₁	3168.96	3171.54	$\nu_5 + 2\nu_7$
<i>A</i> ₁	3170.54	???	$4\nu_{15}$
<i>A</i> ₂	3171.59	3175.51	$\nu_{15} + \nu_3 + \nu_{11}$
<i>B</i> ₂	3175.16	3179.51	$\nu_{15} + \nu_3 + \nu_4$
<i>A</i> ₂	3178.65	3182.09	$2\nu_8 + \nu_7$
<i>B</i> ₁	3181.45	3189.11	$\nu_5 + \nu_{12} + \nu_2$
<i>A</i> ₂	3190.24	???	$3\nu_{15} + \nu_{12}$
<i>B</i> ₁	3192.71	3198.95	$\nu_{15} + \nu_3 + \nu_7$
<i>A</i> ₁	3201.63	3205.73	$\nu_{12} + \nu_3 + \nu_{11}$
<i>B</i> ₁	3201.86	3206.84	$\nu_{12} + \nu_3 + \nu_4$
<i>A</i> ₁	3206.74	3209.19	$\nu_{15} + \nu_3 + \nu_{14}$
<i>A</i> ₁	3216.11	???	$2\nu_5 + 2\nu_{12}$
<i>B</i> ₁	3221.25	3225.03	$2\nu_5 + \nu_{10}$
<i>A</i> ₂	3224.44	3227.09	$\nu_{12} + \nu_3 + \nu_{14}$
<i>B</i> ₂	3231.59	3235.46	$\nu_{12} + \nu_3 + \nu_7$

Table I. Inverted or permuted wave numbers are italicized, the HI-RRBPM numbers which deviate by more than 5 cm⁻¹ are in bold. Some wave numbers are missing in the HI-RRBPM list and are marked by the symbol ???. This is probably because they would appear higher in the spectrum. They are overtones or combinations bands of low frequency modes which does not seem to be well described in the HI-RRBPM calculation arguably due to too low basis truncation criteria. As a matter of fact, most of the bold number correspond to bands involving ν_5 , ν_{12} or ν_{15} . Some footnotes in this tables are related to those in Tab. 2 of Ref. 1.

^a main weight on ν_1 in VMFCI calculation, as for the eigenstate corresponding to 2918.79 cm⁻¹.

^b both assignments to $\nu_{12} + 2\nu_{11}$ in the new calculation.

^{c1,c2} inversion confirmed in the new calculation.

^d still assigned to $\nu_{15} + \nu_7 + \nu_{14}$ in the new calculation.

^e permutation of 3 levels in the new calculation instead of just 2.

C. Comparison of the new VMFCI reference calculation with a previous VMFCI calculation using a different contraction-truncation scheme

Here, we display the 200 lowest eigenvalues of the new VMFCI calculation used as a reference for the AnharmoniCaOs calculations in the main text. It is compared

to a previous VMFCI calculation, referred to as VMFCI(11300), which uses force field 2 but was compared to force field 3 HI-RRBPM calculations in Ref 1.

The contraction-truncation schemes of the two calculations are different. The contraction-truncation scheme of the new VMFCI(11956.6) calculation is the same as that of the previous section except that the 2-contraction step VMFCI($\{\{\nu_1, \nu_6, \nu_9, \nu_{13}\}_{17523}, \{\nu_2, \nu_{10}\}_{12272}, \{\nu_8, \nu_{15}\}_{12077}\}$, 17914; $\{\{\nu_3, \nu_5, \nu_{12}\}_{11459}, \{\nu_4, \nu_{11}\}_{11156}, \{\nu_7, \nu_{14}\}_{9633}\}$, 15280), is iterated once, leading to a common ZPE for the two contractions of 12492.353 ± 0.001 cm⁻¹. This ZPE is to be compared with that of the 6-contraction VSCFCI step of 12502.883 ± 0.001 cm⁻¹ and with the converged VSCF ZPE of 12540.280269 cm⁻¹.

The size of the final VCI is 429499 splitted into 4 *C*_{2V}-symmetry blocks of dimensions 107852 for A1, 107096 for A2, 107415 for B1, 107136 for B2.

The contraction-truncation scheme of the old VMFCI(11300) calculation is essentially the same as that of Ref. 2, with a larger final threshold: 11300 cm⁻¹ instead of 10721 cm⁻¹. It can be abbreviated as: VSCF/VSCFCI($\{\nu_1, \nu_6, \nu_9, \nu_{13}\}$, 34072; $\{\nu_3, \nu_5, \nu_{12}\}$, 17000)/VSCFCI($\{\{\nu_1, \nu_6, \nu_9, \nu_{13}\}, \nu_{10}, \nu_{15}\}$, 19464)/VCI(11300)

The size of the final VCI is larger than for the new one with 456201 basis functions splitted into 4 *C*_{2V}-symmetry blocks of dimensions 114852 for A1, 113489 for A2, 114278 for B1, 113582 for B2. However, the results are clearly less accurate (see Table S2 below).

TABLE II: Lowest 200 vibrational energy levels (in cm⁻¹) of ethylene oxide for force field 2

irreps.	VMFCI(11956.6)	VMFCI(11300) ¹
<i>A</i> ₁	12463.57	12463.65
<i>B</i> ₂	793.03	793.37
<i>B</i> ₁	822.48	822.75
<i>A</i> ₁	878.76	879.04
<i>A</i> ₂	1017.97	1018.96
<i>A</i> ₁	1122.14	1122.96
<i>B</i> ₁	1124.37	1125.18
<i>B</i> ₂	1146.45	1147.09
<i>A</i> ₂	1148.42	1149.02
<i>A</i> ₁	1271.11	1271.50
<i>B</i> ₁	1467.98	1468.34
<i>A</i> ₁	1496.00	1496.65
<i>A</i> ₁	1588.62	1589.38
<i>A</i> ₂	1612.27	1612.92
<i>A</i> ₁	1642.48	1642.86
<i>B</i> ₂	1671.38	1671.94
<i>B</i> ₁	1696.43	1696.71
<i>A</i> ₁	1755.98	1756.21
<i>B</i> ₁	1807.45	1808.96
<i>B</i> ₂	1834.31	1835.47
<i>A</i> ₂	1890.96	1891.96

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irreps.	VMFCI(11956.6)	VMFCI(11300) ¹
A_2	1909.35	1910.19
B_2	1911.54	1912.38
B_1	1929.48	1930.21
A_1	1939.41	1940.01
B_1	1941.36	1941.82
A_1	1945.44	1945.80
A_2	1963.29	1963.62
B_2	1971.50	1971.83
A_1	1999.08	1999.34
B_1	2000.51	2000.76
B_2	2022.30	2022.57
A_2	2024.90	2025.19
A_1	2033.60	2036.08
B_2	2061.03	2061.51
B_1	2087.26	2087.49
A_2	2130.39	2132.05
B_2	2136.49	2138.11
A_1	2141.59	2141.80
B_1	2154.23	2155.72
A_1	2166.39	2167.73
A_1	2237.29	2237.65
B_1	2243.54	2243.80
A_2	2247.41	2247.97
A_1	2250.51	2250.86
B_2	2265.43	2266.28
A_2	2268.70	2269.27
B_2	2269.00	2269.43
A_2	2275.80	2276.17
B_2	2283.65	2284.40
A_1	2289.19	2289.57
A_2	2290.60	2291.52
A_1	2291.68	2292.21
B_1	2294.23	2294.73
A_1	2296.53	2297.09
B_1	2311.89	2312.71
B_1	2346.25	2346.58
A_1	2371.38	2372.12
B_2	2382.39	2382.86
A_1	2390.42	2390.76
B_1	2391.82	2392.13
B_1	2403.66	2404.62
B_2	2415.82	2416.22
A_2	2416.33	2416.72
B_2	2429.68	2430.50
B_1	2459.82	2460.42
A_1	2466.23	2467.01
B_2	2482.36	2483.52
A_2	2485.46	2486.14
A_2	2508.00	2510.34
A_1	2511.90	2512.47
A_1	2537.63	2538.02
B_2	2547.89	2548.58
B_1	2568.81	2569.28
B_1	2584.82	2585.75
A_1	2588.50	2589.46
A_2	2596.91	2599.54
A_2	2600.64	2601.69

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irreps.	VMFCI(11956.6)	VMFCI(11300) ¹
B_2	2602.50	2603.53
A_1	2613.54	2615.39
B_1	2617.19	2618.85
A_1	2618.61	2621.28
A_1	2631.18	2631.55
A_2	2641.56	2643.15
B_2	2642.02	2643.48
A_2	2648.15	2650.27
B_1	2677.78	2680.29
B_1	2695.22	2697.64
A_1	2700.93	2703.48
B_2	2702.44	2704.29
A_2	2712.80	2715.21
B_2	2722.99	2725.03
A_2	2729.73	2731.66
B_2	2735.91	2737.56
B_1	2737.38	2737.96
A_1	2745.45	2747.23
B_1	2752.86	2754.35
A_1	2758.55	2760.02
B_1	2762.24	2763.28
A_2	2762.32	2763.83
A_1	2763.49	2764.81
B_2	2777.03	2778.08
A_2	2785.37	2786.85
B_2	2787.21	2788.85
A_2	2791.15	2792.07
B_1	2804.45	2805.89
A_1	2812.84	2813.86
B_1	2814.12	2815.20
A_1	2817.31	2818.25
B_2	2817.53	2822.93
A_2	2834.11	2834.86
B_2	2842.57	2843.35
B_1	2842.87	2847.79
A_1	2852.76	2854.33
A_2	<i>2873.20</i>	<i>2874.44</i>
A_1	<i>2873.63</i>	<i>2874.30</i>
B_1	2874.26	2874.81
B_2	2895.96	2896.62
A_2	2899.13	2899.84
A_1	2900.38	2905.43
A_1	2901.94	2902.67
B_1	2909.55	2910.86
A_1	2914.73	2916.69
B_1	2916.95	2922.14
A_1	2920.29	2924.52
B_2	2930.71	2931.82
B_2	2936.85	2941.73
A_2	2942.08	2946.78
B_2	2947.75	2951.93
A_2	2950.91	2954.90
B_1	2953.94	2954.55
A_1	2955.01	2956.38
A_1	2967.42	2971.25
B_1	2982.39	2985.81
B_1	2992.22	2993.97

continued on next page

irreps.	VMFCI(11956.6)	VMFCI(11300) ¹
<i>A</i> ₁	2997.03	3000.08
<i>A</i> ₂	3001.27	3005.26
<i>B</i> ₂	3005.13	3009.46
<i>A</i> ₁	3010.51	3011.16
<i>B</i> ₂	3015.27	3020.67
<i>A</i> ₂	3022.83	3027.65
<i>B</i> ₁	3025.19	3028.81
<i>A</i> ₂	3028.63	3029.75
<i>B</i> ₂	3032.42	3035.82
<i>B</i> ₁	3035.32	3038.05
<i>A</i> ₁	3035.53	3039.04
<i>B</i> ₂	3040.60	3041.99
<i>A</i> ₁	3043.97	3047.88
<i>A</i> ₂	<i>3044.37</i>	3055.15
<i>B</i> ₁	<i>3045.56</i>	<i>3049.21</i>
<i>B</i> ₁	3052.13	3055.92
<i>A</i> ₁	3054.74	3057.68
<i>A</i> ₁	3057.67	3061.06
<i>B</i> ₁	3061.83	3064.98
<i>B</i> ₂	3064.57	3066.76
<i>A</i> ₁	3065.11	3068.03
<i>B</i> ₁	3066.99	3069.49
<i>B</i> ₂	3070.40	3073.86
<i>A</i> ₂	3070.63	3073.86
<i>B</i> ₁	3075.31	3078.94
<i>A</i> ₂	3080.74	3084.02
<i>B</i> ₂	3082.16	3084.92
<i>A</i> ₂	3085.30	3087.97
<i>B</i> ₂	3086.36	3089.06
<i>B</i> ₂	3093.99	3096.23
<i>A</i> ₂	3100.17	3102.56
<i>B</i> ₂	3101.89	3104.44
<i>B</i> ₁	3102.46	3105.97
<i>B</i> ₁	3108.24	3109.55
<i>A</i> ₁	3110.78	3112.92 ^b
<i>A</i> ₁	3111.10	3116.35
<i>B</i> ₁	3116.21	3118.38
<i>B</i> ₁	3119.06	3122.34
<i>A</i> ₁	3122.94	3127.59
<i>A</i> ₂	3124.07	3126.44
<i>A</i> ₁	3126.88	3129.22
<i>A</i> ₁	<i>3134.00</i>	3146.79
<i>A</i> ₂	<i>3139.32</i>	<i>3143.09</i>
<i>B</i> ₂	<i>3139.67</i>	<i>3143.48</i>
<i>B</i> ₂	<i>3142.20</i>	<i>3146.04</i>
<i>B</i> ₁	<i>3143.98</i>	3156.68
<i>A</i> ₂	<i>3149.94</i>	<i>3153.20</i>
<i>A</i> ₂	3156.84	3159.94
<i>B</i> ₂	3157.70	3160.91
<i>B</i> ₂	<i>3159.06</i>	3170.99
<i>A</i> ₁	<i>3162.45</i>	<i>3163.64</i>
<i>A</i> ₁	<i>3164.00</i>	<i>3168.19</i>
<i>B</i> ₁	<i>3167.12</i>	<i>3170.35</i>
<i>A</i> ₁	3169.90	3173.84
<i>A</i> ₁	3172.17	3174.28
<i>A</i> ₂	3172.61	3175.82
<i>B</i> ₂	3176.29	3179.58

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irreps.	VMFCI(11956.6)	VMFCI(11300) ¹
<i>A</i> ₂	<i>3180.39</i>	3190.21
<i>B</i> ₁	<i>3183.15</i>	<i>3186.27</i>
<i>A</i> ₂	3191.97	3195.37
<i>B</i> ₁	3193.72	3197.44
<i>A</i> ₁	3202.65	3206.44
<i>B</i> ₁	3202.99	3206.78
<i>A</i> ₁	3207.67	3210.84
<i>A</i> ₁	<i>3217.88</i>	??? ^a
<i>B</i> ₁	3222.57	3223.58
<i>A</i> ₂	3225.77	3229.36
<i>B</i> ₂	3232.42	3235.64

Table II. Inverted or permuted wave numbers are italicized, the VMFCI(11300) numbers which deviate by more than 5 cm⁻¹ are in bold. Some wave numbers are missing in the VMFCI(11300) list and are marked by the symbol ???. This is because they appear higher in the spectrum. As a matter of fact, the quality of the VMFCI(11300) calculation deteriorate beyond 3110 cm⁻¹ where many bold and italicized numbers appear. Some footnotes in this tables are related to those in Tab.2 of Ref. 1.

^a assigned to $2\nu_5 + 2\nu_{12}$ in the new calculation, missing in the old one.

^b this wave number is assigned to $\nu_{12} + \nu_8 + \nu_{15}$ in the new calculation and to $\nu_{12} + \nu_7 + \nu_{14}$ in the old one.

II. HARMONIC VIBRATIONAL ANALYSES

The cubic and quartic force field constants, as well as the first and second electric dipole moment derivatives that our AnharmoniCaOs code uses were obtained via DFT calculations with the B97-1 exchange-correlation functional with the TZ2P and 6-31G* Gaussian basis sets, as described in the text. Since we used numerical differentiation along normal coordinates, we give here the results of the harmonic vibrational analyses for both Pyrene (Table III) and Coronene (Table IV).

TABLE III: List of harmonic vibrational normal modes of Pyrene. For each mode numbered according to spectroscopic conventions, we list the irreducible representation label ("sym." column), the wave number in cm^{-1} , and the intensity of the corresponding fundamental IR transition in $\text{km} \cdot \text{mol}^{-1}$.

mode	sym.	B97-1/6-31g*		B97-1/TZ2P	
		freq.	int.	freq.	int.
ν_1	A_g	3202.4	0.0	3182.2	0.0
ν_2	A_g	3192.4	0.0	3172.8	0.0
ν_3	A_g	3175.9	0.0	3157.5	0.0
ν_4	A_g	1678.5	0.0	1661.0	0.0
ν_5	A_g	1604.9	0.0	1587.8	0.0
ν_6	A_g	1435.6	0.0	1420.8	0.0
ν_7	A_g	1356.2	0.0	1347.3	0.0
ν_8	A_g	1270.5	0.0	1256.9	0.0
ν_9	A_g	1172.1	0.0	1164.9	0.0
ν_{10}	A_g	1094.5	0.0	1085.6	0.0
ν_{11}	A_g	809.8	0.0	811.3	0.0
ν_{12}	A_g	594.6	0.0	593.2	0.0
ν_{13}	A_g	409.6	0.0	409.6	0.0
ν_{14}	A_u	969.0	0.0	985.1	0.0
ν_{15}	A_u	896.2	0.0	908.2	0.0
ν_{16}	A_u	685.5	0.0	689.7	0.0
ν_{17}	A_u	400.9	0.0	399.9	0.0
ν_{18}	A_u	153.1	0.0	149.9	0.0
ν_{19}	B_{1g}	911.8	0.0	920.6	0.0
ν_{20}	B_{1g}	816.5	0.0	815.4	0.0
ν_{21}	B_{1g}	538.5	0.0	533.9	0.0
ν_{22}	B_{1g}	250.3	0.0	246.9	0.0
ν_{23}	B_{1u}	3202.0	71.116	3182.0	38.0
ν_{24}	B_{1u}	3176.1	0.368	3157.8	0.0018
ν_{25}	B_{1u}	3172.9	2.711	3154.4	1.2
ν_{26}	B_{1u}	1643.2	12.622	1626.5	13.2
ν_{27}	B_{1u}	1490.5	0.320	1478.9	1.0
ν_{28}	B_{1u}	1466.4	11.194	1456.5	8.0
ν_{29}	B_{1u}	1271.2	2.776	1264.3	2.3
ν_{30}	B_{1u}	1117.8	5.675	1109.2	7.0
ν_{31}	B_{1u}	1008.1	0.830	1011.5	1.8
ν_{32}	B_{1u}	830.2	3.477	830.3	4.6
ν_{33}	B_{1u}	700.6	0.151	700.8	0.029
ν_{34}	B_{1u}	502.6	2.417	503.7	3.0
ν_{35}	B_{2g}	980.0	0.0	990.7	0.0
ν_{36}	B_{2g}	960.8	0.0	977.0	0.0
ν_{37}	B_{2g}	837.2	0.0	859.2	0.0
ν_{38}	B_{2g}	784.5	0.0	783.6	0.0
ν_{39}	B_{2g}	583.0	0.0	586.5	0.0
ν_{40}	B_{2g}	513.0	0.0	511.7	0.0
ν_{41}	B_{2g}	264.2	0.0	259.7	0.0
ν_{42}	B_{2u}	3192.1	69.930	3172.7	42.8
ν_{43}	B_{2u}	3183.0	22.516	3165.0	12.7
ν_{44}	B_{2u}	1653.6	4.393	1635.8	2.3
ν_{45}	B_{2u}	1525.0	2.782	1508.5	3.2
ν_{46}	B_{2u}	1466.2	2.046	1453.4	3.2
ν_{47}	B_{2u}	1356.3	5.519	1338.3	4.9
ν_{48}	B_{2u}	1236.1	0.026	1226.2	0.020
ν_{49}	B_{2u}	1210.9	13.525	1202.2	14.0
ν_{50}	B_{2u}	1174.7	0.418	1160.4	0.28

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TABLE III – continued from previous page

mode	sym.	B97-1/6-31g*		B97-1/TZ2P	
		freq.	int.	freq.	int.
ν_{51}	B_{2u}	979.3	0.106	975.4	0.0025
ν_{52}	B_{2u}	549.2	2.649	549.2	2.6
ν_{53}	B_{2u}	355.8	1.327	356.5	1.7
ν_{54}	B_{3g}	3183.4	0.0	3165.5	0.0
ν_{55}	B_{3g}	3173.1	0.0	3154.6	0.0
ν_{56}	B_{3g}	1632.6	0.0	1618.3	0.0
ν_{57}	B_{3g}	1543.3	0.0	1528.9	0.0
ν_{58}	B_{3g}	1443.4	0.0	1433.0	0.0
ν_{59}	B_{3g}	1411.9	0.0	1393.9	0.0
ν_{60}	B_{3g}	1264.6	0.0	1260.7	0.0
ν_{61}	B_{3g}	1203.0	0.0	1194.9	0.0
ν_{62}	B_{3g}	1125.4	0.0	1122.9	0.0
ν_{63}	B_{3g}	744.4	0.0	745.2	0.0
ν_{64}	B_{3g}	503.5	0.0	504.3	0.0
ν_{65}	B_{3g}	459.2	0.0	459.7	0.0
ν_{66}	B_{3u}	968.6	1.753	983.8	1.4
ν_{67}	B_{3u}	859.3	109.082	860.0	111.9
ν_{68}	B_{3u}	759.5	11.473	755.1	17.5
ν_{69}	B_{3u}	722.5	25.863	725.6	44.8
ν_{70}	B_{3u}	497.2	1.318	496.3	2.1
ν_{71}	B_{3u}	216.2	6.968	210.2	10.6
ν_{72}	B_{3u}	100.2	0.409	97.3	0.64

TABLE IV: List of harmonic vibrational normal modes of coronene. For each mode numbered according to spectroscopic conventions, we list the irreducible representation label ("sym." column), the wave number in cm^{-1} , and the intensity of the corresponding fundamental IR transition in $\text{km} \cdot \text{mol}^{-1}$.

mode	sym.	B97-1/6-31g*		B97-1/TZ2P	
		freq.	int.	freq.	int.
ν_1	A_{1g}	3194.3	0.0	3174.4	0.0
ν_2	A_{1g}	1646.4	0.0	1629.3	0.0
ν_3	A_{1g}	1385.1	0.0	1366.1	0.0
ν_4	A_{1g}	1250.5	0.0	1240.7	0.0
ν_5	A_{1g}	1048.4	0.0	1043.5	0.0
ν_6	A_{1g}	483.9	0.0	483.0	0.0
ν_7	A_{1u}	946.8	0.0	964.2	0.0
ν_8	A_{1u}	526.4	0.0	530.4	0.0
ν_9	A_{2g}	3172.8	0.0	3154.1	0.0
ν_{10}	A_{2g}	1580.0	0.0	1565.2	0.0
ν_{11}	A_{2g}	1257.7	0.0	1254.9	0.0
ν_{12}	A_{2g}	931.2	0.0	941.2	0.0
ν_{13}	A_{2g}	637.5	0.0	640.1	0.0
ν_{14}	A_{2u}	873.5	158.034	873.2	167.908
ν_{15}	A_{2u}	562.7	24.821	558.4	42.749
ν_{16}	A_{2u}	128.0	4.784	123.4	6.911
ν_{17}	B_{1g}	775.6	0.0	773.6	0.0
ν_{18}	B_{1g}	166.1	0.0	162.9	0.0
ν_{19}	B_{1u}	3175.4	0.0	3156.7	0.0

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TABLE IV – continued from previous page

mode	sym.	B97-1/6-31g*		B97-1/TZ2P	
		freq.	int.	freq.	int.
ν_{20}	B _{1u}	1585.9	0.0	1577.7	0.0
ν_{21}	B _{1u}	1453.2	0.0	1443.1	0.0
ν_{22}	B _{1u}	1175.1	0.0	1181.5	0.0
ν_{23}	B _{1u}	678.2	0.0	677.2	0.0
ν_{24}	B _{1u}	560.0	0.0	563.3	0.0
ν_{25}	B _{2g}	972.3	0.0	985.4	0.0
ν_{26}	B _{2g}	789.7	0.0	869.0	0.0
ν_{27}	B _{2g}	635.7	0.0	656.5	0.0
ν_{28}	B _{2g}	227.8	0.0	224.1	0.0
ν_{29}	B _{2u}	3190.7	0.0	3171.4	0.0
ν_{30}	B _{2u}	1527.8	0.0	1511.2	0.0
ν_{31}	B _{2u}	1381.8	0.0	1367.2	0.0
ν_{32}	B _{2u}	1218.5	0.0	1205.0	0.0
ν_{33}	B _{2u}	1168.5	0.0	1153.5	0.0
ν_{34}	B _{2u}	476.9	0.0	477.0	0.0
ν_{35}	E _{1g}	954.2	0.0	971.3	0.0
ν_{36}	E _{1g}	954.2	0.0	971.3	0.0
ν_{37}	E _{1g}	852.9	0.0	854.8	0.0
ν_{38}	E _{1g}	852.9	0.0	854.8	0.0
ν_{39}	E _{1g}	668.5	0.0	672.8	0.0
ν_{40}	E _{1g}	668.5	0.0	672.8	0.0
ν_{41}	E _{1g}	454.3	0.0	454.3	0.0
ν_{42}	E _{1g}	454.3	0.0	454.3	0.0
ν_{43}	E _{1g}	300.1	0.0	292.8	0.0
ν_{44}	E _{1g}	300.1	0.0	292.8	0.0
ν_{45}	E _{1u}	3193.0	107.828	3173.4	63.965
ν_{46}	E _{1u}	3193.0	107.849	3173.4	63.973
ν_{47}	E _{1u}	3173.6	6.729	3154.9	3.395
ν_{48}	E _{1u}	3173.6	6.728	3154.9	3.406
ν_{49}	E _{1u}	1660.5	12.129	1643.9	7.615
ν_{50}	E _{1u}	1660.5	12.136	1643.9	7.623
ν_{51}	E _{1u}	1540.8	1.630	1526.2	3.024
ν_{52}	E _{1u}	1540.8	1.630	1526.2	3.026
ν_{53}	E _{1u}	1435.1	0.815	1419.0	0.387
ν_{54}	E _{1u}	1435.1	0.815	1419.0	0.392
ν_{55}	E _{1u}	1339.0	23.446	1333.4	21.148
ν_{56}	E _{1u}	1339.0	23.448	1333.4	21.126
ν_{57}	E _{1u}	1241.4	1.029	1232.7	1.255
ν_{58}	E _{1u}	1241.4	1.028	1232.7	1.254
ν_{59}	E _{1u}	1161.6	7.599	1155.0	8.633
ν_{60}	E _{1u}	1161.6	7.603	1155.0	8.628
ν_{61}	E _{1u}	821.0	0.043	819.7	0.0014
ν_{62}	E _{1u}	821.0	0.043	819.7	0.0012
ν_{63}	E _{1u}	779.4	6.525	780.6	6.808
ν_{64}	E _{1u}	779.4	6.525	780.6	6.797
ν_{65}	E _{1u}	382.3	3.153	382.6	3.872
ν_{66}	E _{1u}	382.3	3.153	382.6	3.867
ν_{67}	E _{2g}	3191.5	0.0	3172.1	0.0
ν_{68}	E _{2g}	3191.5	0.0	3172.1	0.0
ν_{69}	E _{2g}	3174.8	0.0	3156.2	0.0
ν_{70}	E _{2g}	3174.8	0.0	3156.2	0.0
ν_{71}	E _{2g}	1660.0	0.0	1645.1	0.0
ν_{72}	E _{2g}	1660.0	0.0	1645.1	0.0
ν_{73}	E _{2g}	1487.9	0.0	1475.9	0.0

Continued on next page

TABLE IV – continued from previous page

mode	sym.	B97-1/6-31g*		B97-1/TZ2P	
		freq.	int.	freq.	int.
ν_{74}	E _{2g}	1487.9	0.0	1475.9	0.0
ν_{75}	E _{2g}	1470.6	0.0	1457.9	0.0
ν_{76}	E _{2g}	1470.6	0.0	1457.9	0.0
ν_{77}	E _{2g}	1430.8	0.0	1418.6	0.0
ν_{78}	E _{2g}	1430.8	0.0	1418.6	0.0
ν_{79}	E _{2g}	1254.5	0.0	1242.3	0.0
ν_{80}	E _{2g}	1254.5	0.0	1242.3	0.0
ν_{81}	E _{2g}	1187.1	0.0	1177.3	0.0
ν_{82}	E _{2g}	1187.1	0.0	1177.3	0.0
ν_{83}	E _{2g}	1006.6	0.0	1004.8	0.0
ν_{84}	E _{2g}	1006.6	0.0	1004.8	0.0
ν_{85}	E _{2g}	687.2	0.0	686.7	0.0
ν_{86}	E _{2g}	687.2	0.0	686.7	0.0
ν_{87}	E _{2g}	492.3	0.0	493.3	0.0
ν_{88}	E _{2g}	492.3	0.0	493.3	0.0
ν_{89}	E _{2g}	367.9	0.0	368.0	0.0
ν_{90}	E _{2g}	367.9	0.0	368.0	0.0
ν_{91}	E _{2u}	965.3	0.0	980.7	0.0
ν_{92}	E _{2u}	965.3	0.0	980.7	0.0
ν_{93}	E _{2u}	814.8	0.0	829.2	0.0
ν_{94}	E _{2u}	814.8	0.0	829.2	0.0
ν_{95}	E _{2u}	770.2	0.0	780.6	0.0
ν_{96}	E _{2u}	770.2	0.0	780.6	0.0
ν_{97}	E _{2u}	552.9	0.0	551.6	0.0
ν_{98}	E _{2u}	552.9	0.0	551.6	0.0
ν_{99}	E _{2u}	302.9	0.0	299.2	0.0
ν_{100}	E _{2u}	302.9	0.0	299.2	0.0
ν_{101}	E _{2u}	89.1	0.0	86.7	0.0
ν_{102}	E _{2u}	89.1	0.0	86.7	0.0

It is well known that vibrational frequencies computed in the harmonic approximation can be brought in fairly good agreement with experimental values by scaling them with empirical scaling factors (see e. g. Ref. 5). The scaling factor(s) are typically obtained by fitting a large number of identified vibrational transitions, depend on the level of theory and basis set used, and may be different for different types of vibrations (i.e. stretches, bends, twists, the kinds of bonds involved etc.).^{6,7} Scaling factors for the B97-1 functional in combination with the TZ2P basis set were obtained by Cané, Miani, and Trombetti⁸ and confirmed by Maltseva *et al.*^{9,10} with minimal variation. For the sake of comparison with anharmonic calculations we here used the original ones taken from Ref. 8, namely 0.9660 for C-H stretches and 0.982 for all the other bands. Given the very small sample of bands these factors were calibrated on, their general validity is not adequately tested, and they should be used with caution.

III. ANHARMONIC BANDS OF PYRENE AS A FUNCTION OF r

The large set of bands listed in Table V results from anharmonic calculations with different r values, keeping fixed $h = 8$. State assignment was done, as far as this table is concerned, on the calculation with $r = 0.3$, where it is the least ambiguous. In most eigenstates a single harmonic basis function is dominant. In the few cases where several harmonic basis functions have weights in the same range, the latter are all listed. Only transitions with non-negligible intensity are included in the table.

When decreasing the r -value, some bands split in several ones due to resonances. Since this splitting may involve significant mixing with bands which are distinct at the $r = 0.3$ level, sometimes the same wave number appears several times in columns corresponding to smaller r values. This happens, for example, for the bands which are computed at 3034, 3037, and 3042 cm^{-1} at the $r = 0.3$ level. These give rise to a multitude of bands at smaller r values, of which only the strongest are listed. Note also that, in some cases, distinct bands are almost coincident in position. However, they are listed on different lines in the table.

TABLE V: Calculated anharmonic band positions for pyrene as a function of the threshold r and for fixed $h = 8$.

mode	$r=0.3$	$r=0.2$	$r=0.1$	$r=0.05$
$\nu_{44} + \nu_{48} + \nu_{53}$, $\nu_5 + \nu_{44}$	3144	3138, 3144, 3153	3134, 3144, 3146, 3154	3145, 3156
$\nu_{44} + \nu_{57}$	3107	3105, 3106	3103, 3104	3098, 3102
$\nu_{10} + \nu_{52} + \nu_{57}$	3102	3100	3099	3098, 3102
$\nu_4 + \nu_{45}$	3097	3097	3089, 3098	3092, 3098, 3102
$\nu_{34} + \nu_{57} + \nu_{62}$	3096	3095	3092, 3093	3083, 3092, 3095
$\nu_{26} + \nu_{57}$	3080	3079	3080	3062, 3073, 3078
ν_{23}	3055	3039, 3062, 3064, 3069	3010, 3029, 3035, 3069, 3104	3019, 3026, 3030, 3067, 3076
$\nu_4 + \nu_{28}$	3051	3023, 3030, 3036	3029, 3045, 3055, 3058	3056
$\nu_{45} + \nu_{56}$	3043	3029, 3039, 3062, 3064, 3069	3029, 3035, 3045, 3055, 3069	3030, 3043, 3067, 3076
ν_{42} , ν_{43}	3042	3042	3019, 3029, 3037, 3053	3023, 3040
$\nu_4 + \nu_{46}$	3037	3038	3029, 3037	3032, 3033, 3040, 3042, 3044, 3045
ν_{43} , ν_{42}	3034	3035	2946, 3029, 3053	2940, 3032, 3033, 3044, 3045
ν_{24}	3031	3030	2954, 2969, 3010, 3045, 3078, 3079	2956, 3018, 3019, 3026, 3076
$\nu_{49} + \nu_{61}$	2359	2359	2359	2359
$\nu_{14} + \nu_{35}$	1936	1936	1937	1937
$\nu_{35} + \nu_{66}$	1932	1932	1932	1934
$\nu_{14} + \nu_{36}$	1922	1923	1919, 1937	1921, 1937
$\nu_{36} + \nu_{66}$	1919	1919	1918	1919
$\nu_{14} + \nu_{19}$	1866	1866	1865	1866
$\nu_{15} + \nu_{35}$	1852	1552	1850	1848
$\nu_{15} + \nu_{36}$	1839	1839	1838	1842
$\nu_{35} + \nu_{67}$	1824	1825	1825	1823
$\nu_{36} + \nu_{67}$	1809	1810	1810	1766, 1801, 1823
$\nu_{37} + \nu_{66}$	1803	1803	1801, 1803	1801, 1809, 1810
$\nu_{15} + \nu_{19}$	1790	1788	1790	1792
$\nu_{14} + \nu_{20}$	1758	1761	1768	1766
$\nu_{19} + \nu_{67}$	1745	1745	1745	1745
$\nu_{17} + \nu_{18} + \nu_{35} + \nu_{71}$	1742	1742	1746	1741
$\nu_{38} + \nu_{66}$	1740	1745	1746	1743
$\nu_{15} + \nu_{37}$	1727	1727	1726	1726
$\nu_{36} + \nu_{68}$, $\nu_{35} + \nu_{68}$	1705	1700	1701	1659, 1669, 1697
$\nu_{15} + \nu_{20}$	1692	1692	1694	1659, 1669, 1697
$\nu_{35} + \nu_{68}$, $\nu_{35} + \nu_{69}$, $\nu_{38} + \nu_{66}$	1687	1692	1684	1659, 1697, 1766
$\nu_{36} + \nu_{69}$	1669	1665, 1668	1661, 1701	1659, 1669, 1697
$\nu_{16} + \nu_{35}$	1655	1656	1655	1651

Continued on next page

TABLE V – continued from previous page

mode	$r=0.3$	$r=0.2$	$r=0.1$	$r=0.05$
$\nu_{20} + \nu_{67}$	1646	1646	1645	1645
$\nu_{16} + \nu_{36}$	1636	1637	1639	1636
$\nu_{19} + \nu_{68}$	1617	1607, 1619	1606, 1618	1617
$\nu_{38} + \nu_{67}$	1613	1613	1612	1612
ν_{44}	1597	1591, 1599	1592, 1598	1588, 1596, 1600
$\nu_{53} + \nu_{60}$	1592	1592	1581, 1591	1590
ν_{26}	1585	1585	1586	1585
$\nu_{16} + \nu_{19}$	1578	1578	1576, 1578	1575, 1578, 1581, 1585, 1590
$\nu_{17} + \nu_{53} + \nu_{67},$ $\nu_{53} + \nu_{60}$	1544	1544	1544	1544
$\nu_{20} + \nu_{68}$	1537	1540	1540	1537
$\nu_{39} + \nu_{66}$	1521	1521	1521	1521
$\nu_{16} + \nu_{37}$	1514	1514	1513	1510
$\nu_{20} + \nu_{69}$	1509	1506	1509	1505
$\nu_{38} + \nu_{68}$	1489	1489	1486	1485
$\nu_{21} + \nu_{66}$	1475	1465, 1475	1461, 1463, 1474, 1486	1473
ν_{45}	1466	1464, 1465	1462, 1463	1447, 1453, 1459, 1460
$\nu_{14} + \nu_{40}$	1444	1444	1442	1438, 1442, 1447
$\nu_{51} + \nu_{64}$	1429	1417, 1425, 1429	1404, 1412, 1419, 1425, 1442	1424
ν_{27}	1419	1420	1414, 1418	1412, 1417
$\nu_{39} + \nu_{67}$	1417	1411, 1417, 1429	1412, 1419, 1425, 1429	1409, 1418, 1424
$\nu_{51} + \nu_{65}$	1412	1412	1410	1410
$\nu_{15} + \nu_{21}$	1311	1311	1310	1311
ν_{47}	1286	1286	1286	1286
$\nu_{20} + \nu_{70}$	1266	1266	1264	1264
$\nu_{21} + \nu_{68}$	1243	1243	1242	1241
$\nu_{40} + \nu_{68}$	1240	1240	1238	1236
ν_{29}	1190	1189	1189	1189
$\nu_3 + \nu_{64}$	1184	1184	1183	1179, 1180, 1181
ν_{49}	1181	1180	1180	1180, 1181
$\nu_{16} + \nu_{40}$	1157	1157	1155	1153
$\nu_{17} + \nu_{38}$	1096	1096	1097	1096
$\nu_{13} + \nu_{33}$	1093	1093	1091	1089
ν_{30}	1084	1084	1081, 1082	1080, 1081
$\nu_{12} + \nu_{34}$	1083	1083	1081, 1082	1080, 1081
$\nu_{53} + \nu_{63}$	1059	1059	1058	1058
$\nu_{39} + \nu_{70}$	999	997, 999	994, 995, 997	993, 996
$\nu_{52} + \nu_{65}, \nu_{31}$	997	995, 997, 999	986, 994, 997	985, 993, 996
$\nu_{31}, \nu_{52} + \nu_{65}$	962	962	962	962
ν_{66}	845	845	844	843
ν_{67}	820	820	817	816
ν_{32}	742	742	741	739
ν_{68}	712	712	711	710
ν_{69}	542	542	540	538
ν_{52}	497	497	495	493
ν_{34}	486	486	484	483
ν_{70}	349	349	349	342, 347
ν_{53}	202	202	200	199
ν_{71}	95	95	94	92
ν_{72}				

IV. ANHARMONIC SPECTRA OF PYRENE AND CORONENE

In this section we provide, for reference and easier comparison with experimental and/or other theoretical spec-

tra, complete plots of our best computed anharmonic spectra for Pyrene (from Fig. 1 to 13) and Coronene (from Fig. 14 to 27). Each figure covers a range of 300 cm^{-1} , and consecutive plots have an overlap of

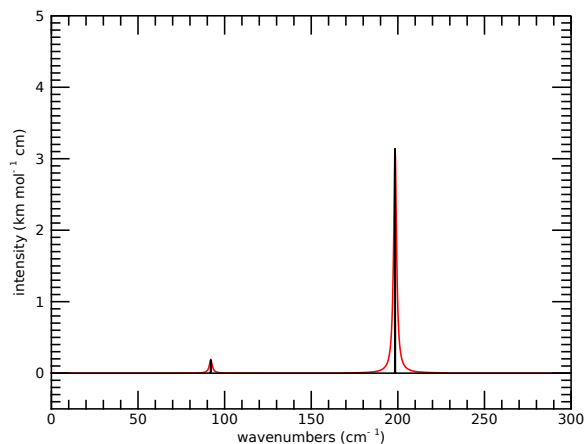


FIG. 1. Anharmonic spectrum of Pyrene computed with $r = 0.05$ and $h = 8$ in the range from 0 to 300 cm^{-1} . Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width.

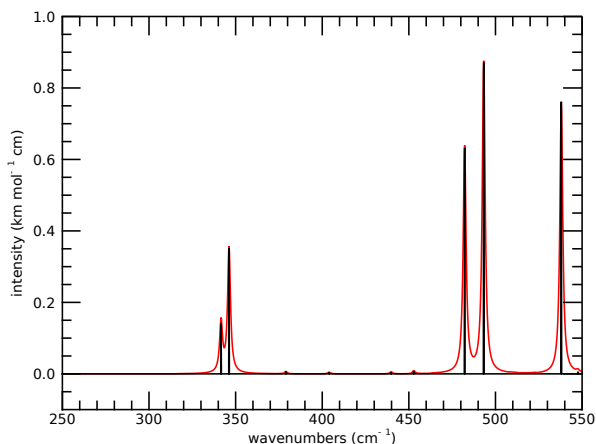


FIG. 2. Same as Fig. 1 in the range from 250 to 550 cm^{-1} .

50 cm^{-1} . The overall spectrum, in red, was obtained by convolving individual bands with a Lorentzian profile with a 1 cm^{-1} width. Intensities of the convolved spectra are absolute and expressed in $\text{km mol}^{-1}\text{cm}$, the standard units of quantum chemistry calculations. The precise positions of individual bands are marked in the spectra by black bars. The complete, anharmonic, tabulated spectra of Pyrene and Coronene can be found in the online PAH database by Mallocci, Joblin, and Mulas.¹¹

¹P. S. Thomas and T. C. Jr., *J. Chem. Phys.* **146**, 204110 (2017).

²D. Bégué, N. Gohaud, C. Pouchan, P. Cassam-Chenaï, and J. Liévin, *J. Chem. Phys.* **127**, 164115 (2007).

³P. S. Thomas and T. Carrington, *J. Phys. Chem. A* **119**, 13074 (2015).

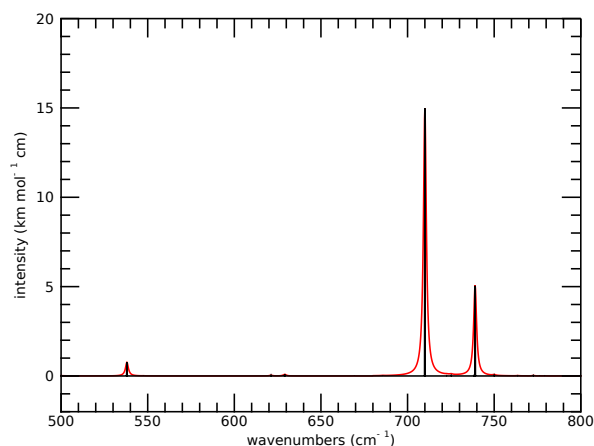


FIG. 3. Same as Fig. 1 in the range from 500 to 800 cm^{-1} .

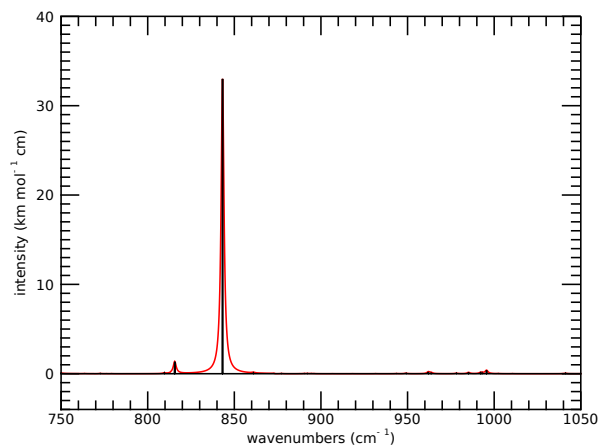


FIG. 4. Same as Fig. 1 in the range from 750 to 1050 cm^{-1} .

⁴M. Odunlami, V. L. Bris, D. Bégué, I. Baraille, and O. Coulaud, *J. Chem. Phys.* **146**, 214108 (2017).

⁵K. K. Irikura, R. D. Johnson, and R. N. Kacker, *J. Phys. Chem. A* **109**, 8430 (2005).

⁶F. Pauzat, D. Talbi, and Y. Ellinger, *Astron. Astrophys.* **293**, 263 (1995).

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⁸E. Cané, A. Miani, and A. Trombetti, *J. Phys. Chem. A* **111**, 8218 (2007).

⁹E. Maltseva, A. Petrignani, A. Candian, C. J. Mackie, X. Huang, T. J. Lee, A. G. G. M. Tielens, J. Oomens, and W. J. Buma, *Astrophys. J.* **831**, 11 (2016).

¹⁰E. Maltseva, A. Petrignani, A. Candian, C. J. Mackie, X. Huang, T. J. Lee, A. G. G. M. Tielens, J. Oomens, and W. J. Buma, *Astrophys. J.* **814**, 23 (2015).

¹¹G. Mallocci, C. Joblin, and G. Mulas, *Chem. Phys.* **332**, 353 (2007), astro-ph/0701254.

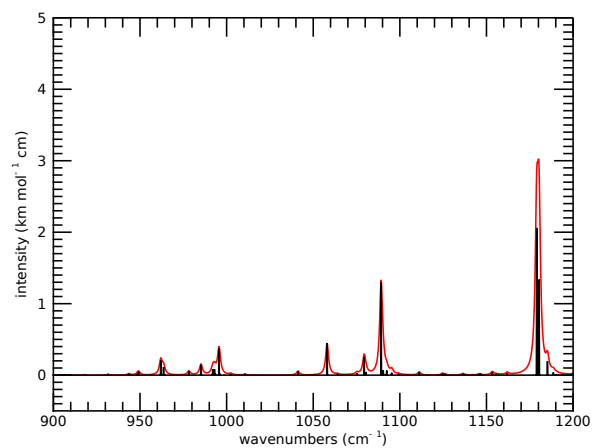


FIG. 5. Same as Fig. 1 in the range from 900 to 1200 cm^{-1} .

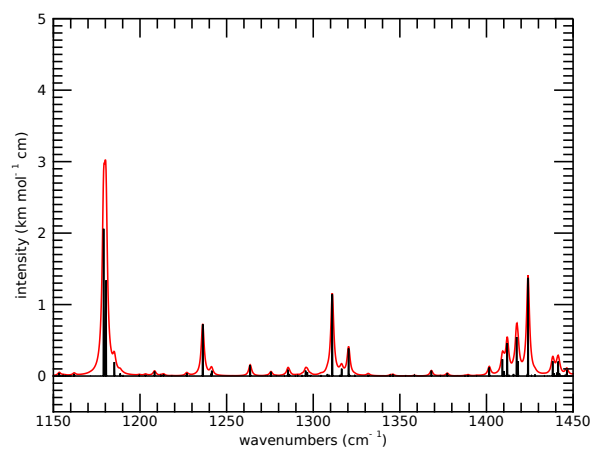


FIG. 6. Same as Fig. 1 in the range from 1150 to 1450 cm^{-1} .

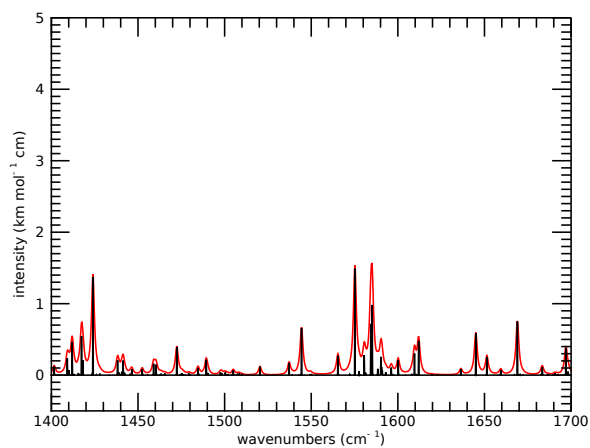


FIG. 7. Same as Fig. 1 in the range from 1400 to 1700 cm^{-1} .

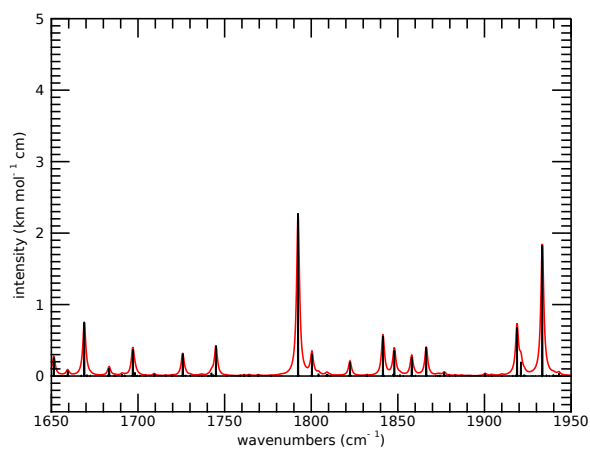


FIG. 8. Same as Fig. 1 in the range from 1650 to 1950 cm^{-1} .

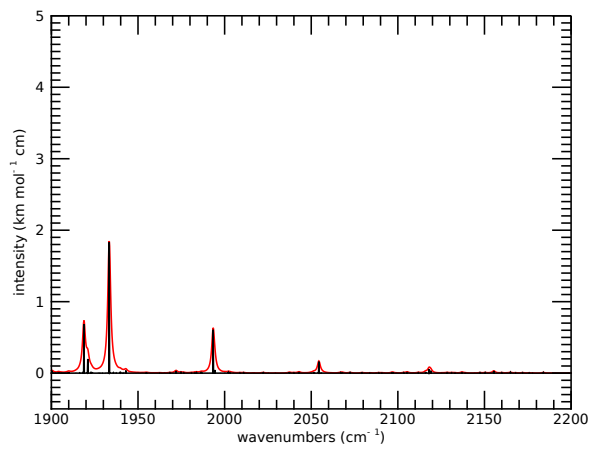


FIG. 9. Same as Fig. 1 in the range from 1900 to 2200 cm⁻¹.

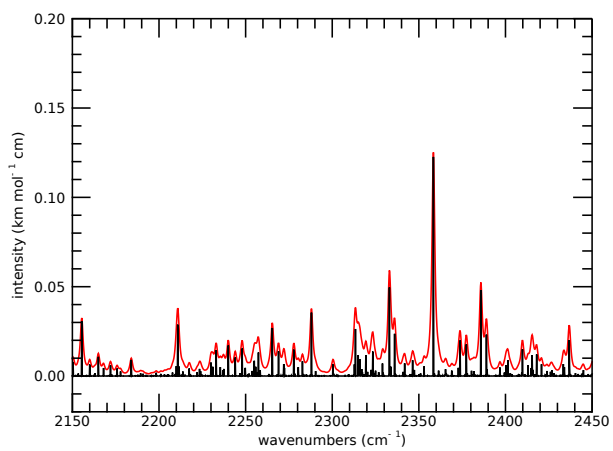


FIG. 10. Same as Fig. 1 in the range from 2150 to 2450 cm⁻¹.

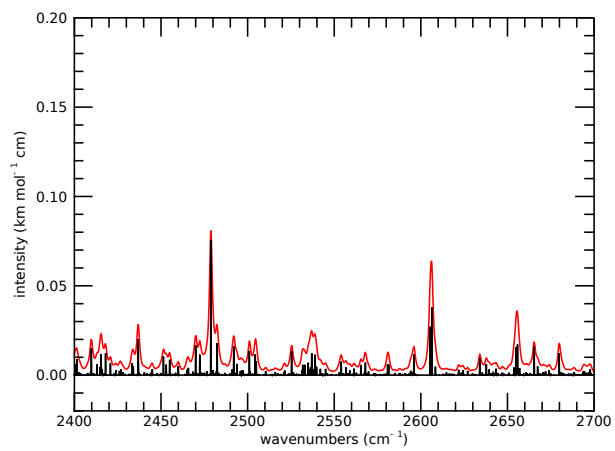


FIG. 11. Same as Fig. 1 in the range from 2400 to 2700 cm^{-1} .

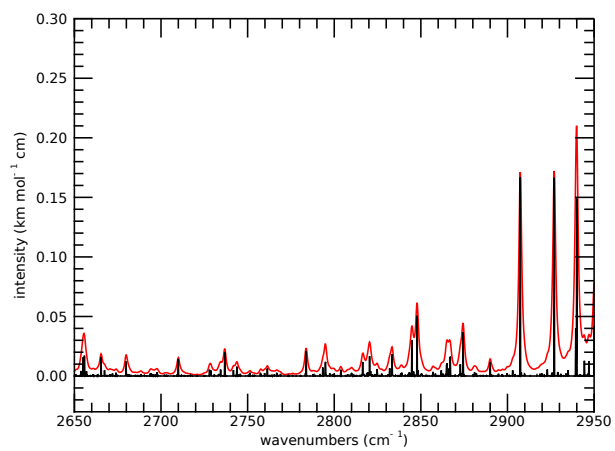


FIG. 12. Same as Fig. 1 in the range from 2650 to 2950 cm^{-1} .

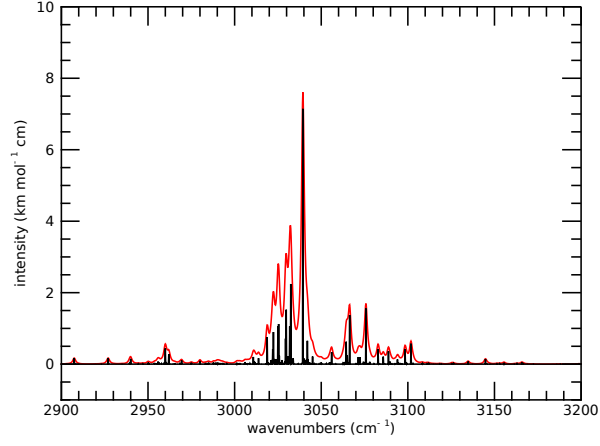


FIG. 13. Same as Fig. 1 in the range from 2900 to 3200 cm^{-1} .

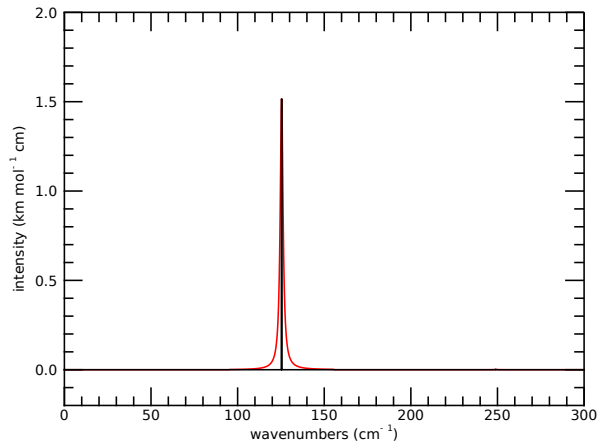


FIG. 14. Anharmonic spectrum of Coronene computed with $r = 0.1$ and $h = 6$ in the range from 0 to 300 cm^{-1} . Black bars indicate the precise position of individual bands, whereas red envelopes are convolved with Lorentzian profiles with a 1 cm^{-1} width.

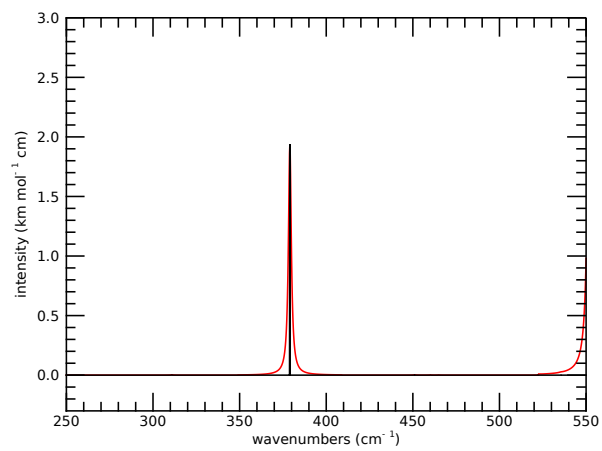


FIG. 15. Same as Fig. 14 in the range from 250 to 550 cm^{-1} .

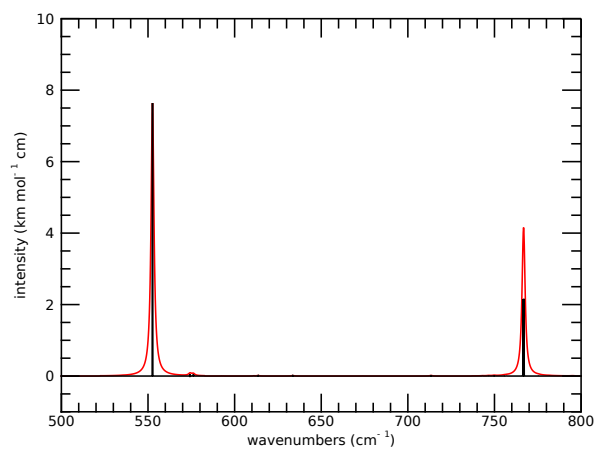


FIG. 16. Same as Fig. 14 in the range from 500 to 800 cm^{-1} .

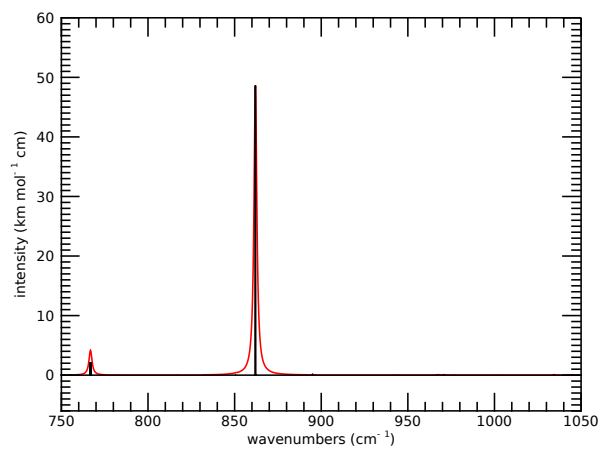


FIG. 17. Same as Fig. 14 in the range from 750 to 1050 cm⁻¹.

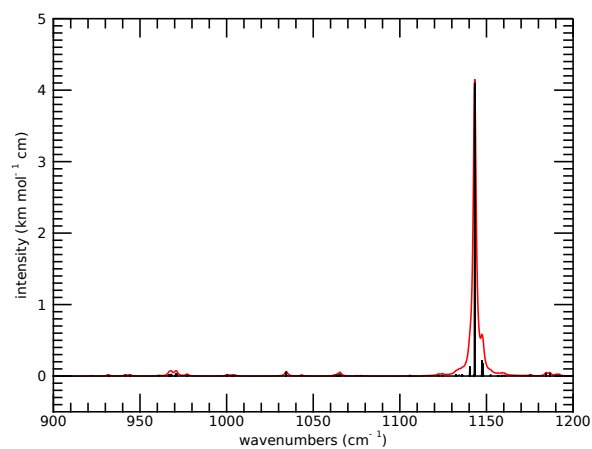


FIG. 18. Same as Fig. 14 in the range from 900 to 1200 cm⁻¹.

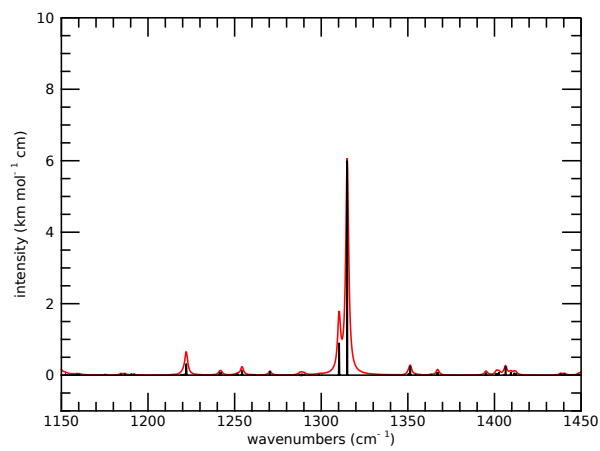


FIG. 19. Same as Fig. 14 in the range from 1150 to 1450 cm^{-1} .

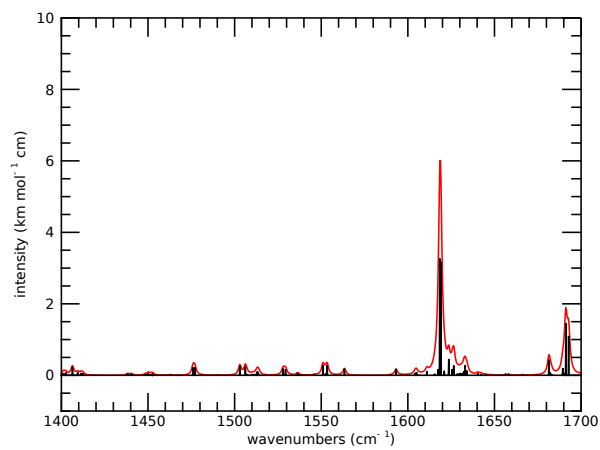


FIG. 20. Same as Fig. 14 in the range from 1400 to 1700 cm^{-1} .

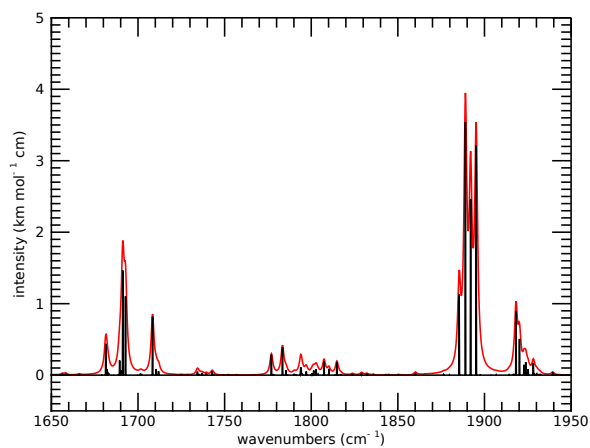


FIG. 21. Same as Fig. 14 in the range from 1650 to 1950 cm^{-1} .

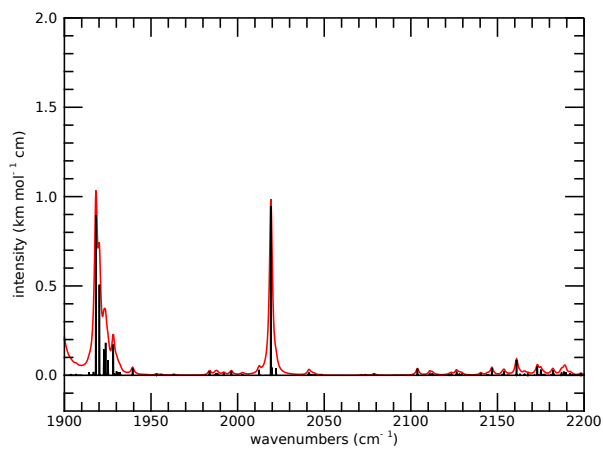


FIG. 22. Same as Fig. 14 in the range from 1900 to 2200 cm^{-1} .

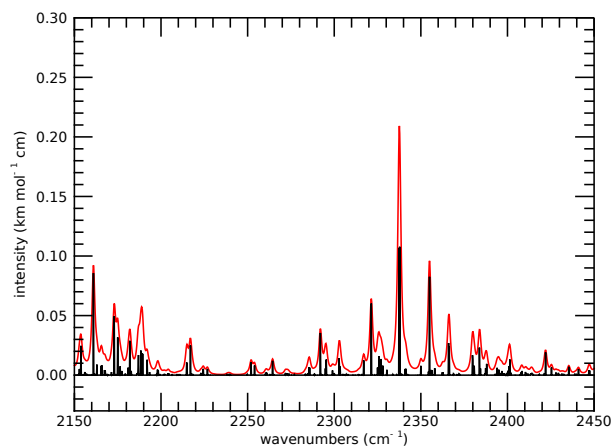


FIG. 23. Same as Fig. 14 in the range from 2150 to 2450 cm^{-1} .

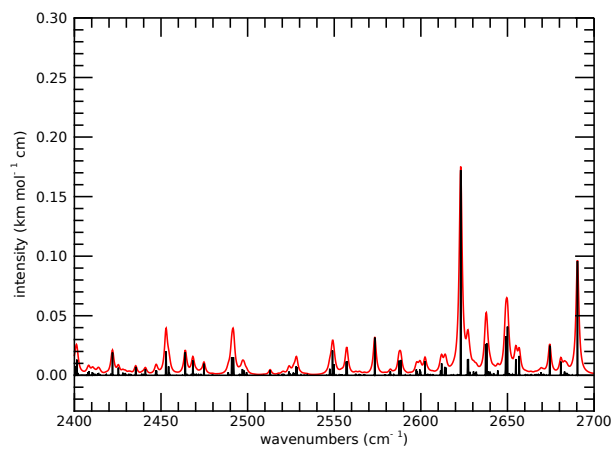


FIG. 24. Same as Fig. 14 in the range from 2400 to 2700 cm^{-1} .

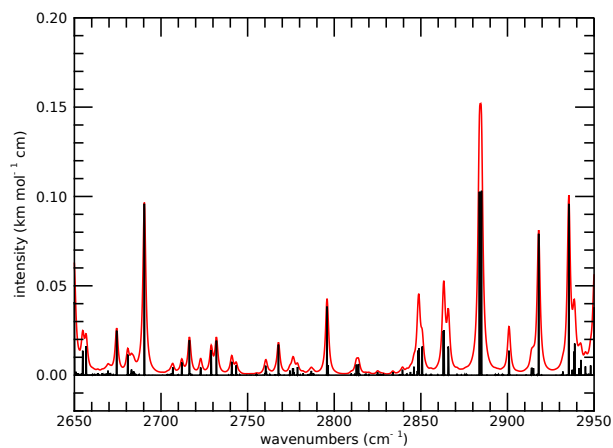


FIG. 25. Same as Fig. 14 in the range from 2650 to 2950 cm^{-1} .

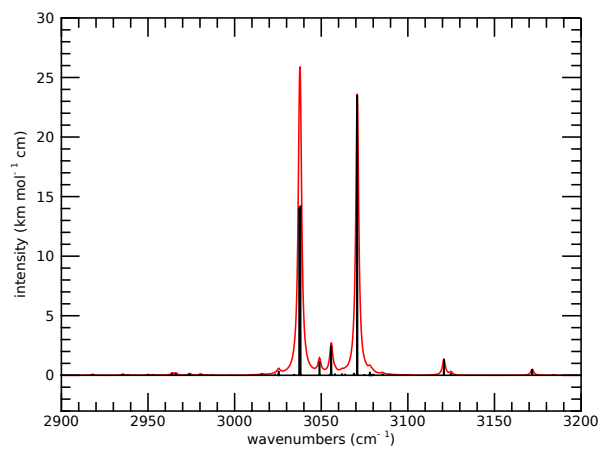


FIG. 26. Same as Fig. 14 in the range from 2900 to 3200 cm^{-1} .

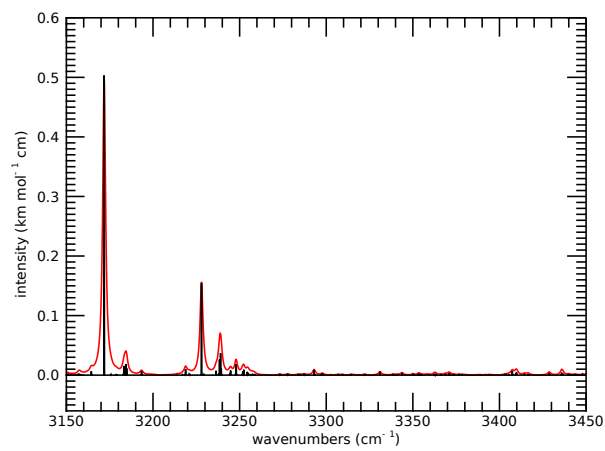


FIG. 27. Same as Fig. 14 in the range from 3150 to 3450 cm^{-1} .