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On the Connection between Turbulent Motions and Particle Acceleration in Galaxy Clusters

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Abstract

Giant radio halos are megaparsec-scale diffuse radio sources associated with the central regions of galaxy clusters. The most promising scenario to explain the origin of these sources is that of turbulent re-acceleration, in which MeV electrons injected throughout the formation history of galaxy clusters are accelerated to higher energies by turbulent motions mostly induced by cluster mergers. In this Letter, we use the amplitude of density fluctuations in the intracluster medium as a proxy for the turbulent velocity and apply this technique to a sample of 51 clusters with available radio data. Our results indicate a segregation in the turbulent velocity of radio halo and radio quiet clusters, with the turbulent velocity of the former being on average higher by about a factor of two. The velocity dispersion recovered with this technique correlates with the measured radio power through the relation $P_{\text{radio}} \propto \sigma_v^{3.3 \pm 0.7}$, which implies that the radio power is nearly proportional to the turbulent energy rate. In case turbulence cascades without being dissipated down to the particle acceleration scales, our results provide an observational confirmation of a key prediction of the turbulent re-acceleration model and possibly shed light on the origin of radio halos.

Key words: acceleration of particles – galaxies: clusters: general – radiation mechanisms: non-thermal – radio continuum: general – turbulence

1. Introduction

Giant radio halos are megaparsec-scale diffuse radio sources associated with the central regions of galaxy clusters (see Feretti et al. 2012 for a review). Radio halos are a transient phenomenon observed in only a fraction of galaxy clusters (e.g., Venturi et al. 2007, 2008; Basu 2012), and recent studies have shown that they occur only in dynamically disturbed systems (Buote 2001; Cassano et al. 2010; Rossetti et al. 2011), suggesting a connection between particle acceleration and major cluster mergers.

The kinetic energy injected during major mergers first generates a turbulent cascade down to small spatial scales (Dolag et al. 2005; Vazza et al. 2009, 2017; Miniati 2015), which may ultimately be dissipated into gas heating, magnetic field amplification, and cosmic-ray acceleration. Stochastic acceleration generated by turbulent motions in the intracluster medium (ICM) is now thought to be the most plausible particle acceleration mechanism to re-energize a pre-existing population of $\sim$ MeV electrons and explain the origin of radio halos (e.g., Brunetti et al. 2001; Petrosian & Bykov 2008; Brunetti & Lazarian 2011, 2016). The turbulent re-acceleration model reproduces a number of observed features, such as the existence of steep-spectrum radio halos (Brunetti et al. 2008; Macario et al. 2010), the curvature of the spectra at high frequency (Thierbach et al. 2003), and the link between radio halos and cluster mergers (Cassano et al. 2010). However, the connection between radio halos and ICM turbulence remains untested.

In this Letter, we use the power spectrum of X-ray surface-brightness fluctuations in the ICM to search for a connection between turbulent motions and radio properties in galaxy clusters. Recent theoretical progress (Schuecker et al. 2004; Churazov et al. 2012; Gaspari & Churazov 2013; Gaspari et al. 2014; Zhuravleva et al. 2014) has shown that gas density fluctuations act as a passive tracer of velocity fluctuations in the ICM and that the maximal amplitude of density fluctuations is linearly related to the turbulent Mach number. This method has been successfully applied to several clusters thus far (Zhuravleva et al. 2015; Arévalo et al. 2016; Hofmann et al. 2016; Khatri & Gaspari 2016) and it was found to reproduce the level of turbulence directly measured in the Perseus cluster by *Hitomi* (Hitomi Collaboration et al. 2016). We retrieve the amplitude of density fluctuations at a fixed spatial scale and use this quantity as a proxy for the expected level of turbulence. We then search for a connection between turbulence and radio emission. The Letter is organized as follows. In Section 2, we present the sample and describe the methodology adopted to analyze the X-ray data. In Section 3, we present our results and discuss them in the framework of the turbulent re-acceleration model in Section 4. In Section 5, we summarize our main conclusions.

2. Analysis

2.1. The Sample

We base our analysis on the sample of 55 clusters with available radio information at the nearly uniform depth from Cassano et al. (2013). The sample is based on the GMRT radio halo survey (Venturi et al. 2007, 2008) with the addition of

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clusters with known radio emission from the literature. We searched the *XMM-Newton*, *Chandra*, and *ROSAT*/PSPC archives for available X-ray data and selected a subsample of 51 clusters for which the quality of the X-ray data is sufficient to retrieve the level of surface-brightness fluctuation over large spatial scales. A radio halo was detected for 25 clusters in our sample, whereas for the remaining 26 systems, upper limits to the radio flux at 1.4 GHz are available. For nearby clusters ($z < 0.1$), we use *ROSAT*/PSPC, as the wide field of view (FOV) of this instrument allows us to cover uniformly a circular region of 1 Mpc radius, i.e., comparable to the typical sizes of radio halos. In the redshift range $0.1 < z < 0.3$, we use *XMM-Newton* as our instrument of choice given its large collecting area and FOV, with the exception of a few cases (e.g., Bullet, A520) for which deep *Chandra* data are available. Beyond $z = 0.3$ only the angular resolution of *Chandra* is sufficient to resolve scales less than ~ 50 kpc; thus, we restrict ourselves to *Chandra* data. The final sample is presented in Table 1 together with its relevant properties and the adopted instrument. In all cases, we restrict ourselves to the [0.5–2] keV band to be sensitive only to density fluctuations. Given that our sample comprises only hot clusters, this choice of energy band is appropriate.

2.2. Data Reduction

2.2.1. XMM-Newton

We reduced the *XMM-Newton* data using XMMSAS v15.0 and the ESAS software package (Snowden et al. 2008). Time periods including flaring soft proton flux are filtered out to extract clean event files. We use the unexposed corners of the CCD chips to measure the quiescent background level in each observation. We then renormalize the filter-wheel-closed data sets to match the count rates measured in the CCD corners, which allows us to create model particle background images for each observation. We extract photon images and exposure maps in the [0.5–2] keV band from the cleaned event files. Finally, to avoid contamination from point sources, we run *ewavelet* in each observation and mask a circular region of 30 arcsec radius around each point source. For more details we refer the reader to Eckert et al. (2014).

2.2.2. Chandra

We analyzed the data using the CIAO v4.9 software package and CALDB v4.7.3. For each observation, the raw data are reprocessed with the latest calibration files by running the *chandra_repro* pipeline. Periods of flaring background are removed by using the *deflare* tool. We then extract photon images and exposure maps in the [0.5–2] keV band. To estimate the local background, we use the *blanksky* and *blanksky_image* tools (Hickox & Markevitch 2006) to estimate the count rate in the [9.5–12] keV band and renormalize blank-sky data sets to the appropriate level for each observation. Point sources are detected using *wavdetect* and masked during the analysis.

2.2.3. ROSAT/PSPC

We reduced the *ROSAT* data using the Extended Source Analysis Software package (Snowden et al. 1994). The analysis pipeline follows exactly the method presented in Eckert et al. (2012). We create photon images, exposure maps, and background in the *ROSAT* R37 band (corresponding to [0.42–2.01] keV) and mask the detected point sources

according to the size of the PSF at each radius. For more details on the analysis procedure we refer to Eckert et al. (2012).

2.3. Temperature Measurements

We extracted mean spectroscopic temperatures within the same region of 1 Mpc radius as for the extraction of the amplitude of surface-brightness fluctuations. For *XMM-Newton*, we follow the method outlined in Eckert et al. (2014). Briefly, we use a phenomenological model to describe the spectral shape of the non X-ray background, which we fit together with the source. The sky background is modeled as the sum of two APEC models (Smith et al. 2001) for the Galactic foregrounds and an absorbed power law for the cosmic X-ray background. In the case of *Chandra*, we use an offset region to describe the local background. In both cases, the source is described as a single-temperature APEC model absorbed by the Galactic N_H , which we fix to the 21 cm value (Kalberla et al. 2005). The spectra are then fit in XSPEC using C statistic (Cash 1979). The best-fit spectroscopic temperatures are reported in Table 1.

2.4. Surface-brightness Fluctuations

To compute the amplitude of density fluctuations, we create a two-dimensional model for the large-scale gas distribution and extract a residual map by dividing the observed emission by the model brightness. The amplitude of surface-brightness fluctuations is then recovered from the Fourier power spectrum of the fluctuations in the residual image at the desired scale.

To model the cluster gas distribution, we assume that the morphology of the cluster can be described by an elliptical beta model (Cavaliere & Fusco-Femiano 1976). We use a weighted principal component analysis method to determine the centroid of emission and the ellipticity parameters (major and minor axes, rotation angle). We mask obvious substructures associated with individual subclumps (e.g., the Bullet in 1E 0657–56) to avoid introducing additional power unrelated to turbulent motions. We then extract a surface-brightness profile using PROFIT (Eckert et al. 2011), and we fit the beta model to the data using C statistic. In the cases where a double beta model is statistically favored, we adopt the double beta model solution as our model of choice. We then create a model image for the two-dimensional gas distribution by folding the best-fit model with the exposure map and adding the background map.

To compute the Fourier power spectrum at the desired scale, we use the modified Δ -variance method introduced by Arévalo et al. (2012), which allows us to take holes and non-periodic boxes into account. In this method, the image and the mask are convolved with a Mexican hat filter, and the filtered image is corrected for the spurious features introduced by holes. The variance of the filtered image is proportional to the power at the chosen scale. To estimate the noise level, we simulate Poisson noise on top of the model image and apply the same procedure. The variance of the noise image is then subtracted from that of the true image. The uncertainty in the measurement of the power spectrum is estimated by splitting the filtered image into 20 subregions, computing the power in each region separately, and performing 10^4 bootstrap resamplings of the measured values. We then adopt the standard deviation of the bootstrap distribution as our 1σ error.

Table 1
Properties of the Cluster Sample^a

Cluster	z	kT_{Mpc} [keV]	$P_{1.4}$ [W Hz ⁻¹]	A_{2D}	Mach _{1D}	σ_v [km s ⁻¹]	Instrument
Upper Limits							
A2697	0.232	6.94 ± 0.14	<0.41	0.06 ± 0.01	0.12 ± 0.01	250 ± 27	X
A141	0.23	5.9 ± 0.15	<0.43	0.08 ± 0.01	0.16 ± 0.02	325 ± 38	X
A3088	0.2537	6.42 ± 0.16	<0.43	0.07 ± 0.01	0.14 ± 0.02	291 ± 39	X
RXCJ0437.1+0043	0.285	6.45 ± 0.2	<0.65	0.04 ± 0.01	0.09 ± 0.02	179 ± 39	X
RXCJ1115.8+0129	0.3499	6.21 ± 0.11	<0.47	0.06 ± 0.01	0.11 ± 0.02	227 ± 41	X
A2631	0.2779	7.61 ± 0.3	<0.41	0.09 ± 0.01	0.16 ± 0.02	374 ± 48	X
A2645	0.251	6.28 ± 0.18	<0.59	0.08 ± 0.01	0.15 ± 0.01	315 ± 25	X
A2667	0.2264	5.99 ± 0.08	<0.45	0.08 ± 0.01	0.15 ± 0.01	310 ± 24	X
Z348	0.2537	2.94 ± 0.05	<0.65	0.05 ± 0.01	0.10 ± 0.02	144 ± 22	X
RXJ0142.0+2131	0.2803	7.37 ± 0.58	<0.45	0.09 ± 0.02	0.18 ± 0.03	402 ± 70	C
A267	0.231	5.64 ± 0.17	<0.34	0.08 ± 0.01	0.16 ± 0.01	306 ± 27	X
RXJ0439.0+0715	0.23	6.23 ± 0.48	<0.46	0.07 ± 0.02	0.13 ± 0.04	272 ± 77	C
RXJ0439.0+0520	0.208	6.53 ± 0.25	<0.32	0.01 ± 0.04	0.02 ± 0.08	41 ± 167	C
A611	0.288	6.63 ± 0.45	<0.43	0.07 ± 0.03	0.13 ± 0.05	273 ± 105	X
Z2089	0.2347	3.82 ± 0.13	<0.26	0.07 ± 0.01	0.13 ± 0.01	205 ± 16	X
A781	0.2984	6.21 ± 0.12	<0.36	0.11 ± 0.01	0.22 ± 0.02	457 ± 42	X
Z2701	0.214	5.41 ± 0.22	<0.35	0.07 ± 0.01	0.13 ± 0.02	250 ± 42	C
A1423	0.213	7.65 ± 0.39	<0.38	0.07 ± 0.01	0.13 ± 0.01	306 ± 32	C
A1576	0.279	7.73 ± 0.46	<0.64	0.10 ± 0.02	0.19 ± 0.03	429 ± 69	C
RXJ1532.9+3021	0.345	4.89 ± 0.07	<0.66	0.04 ± 0.00	0.07 ± 0.01	128 ± 15	X
A2146	0.2343	6.59 ± 0.17	<0.39	0.08 ± 0.01	0.15 ± 0.02	319 ± 35	C
A2261	0.224	8.01 ± 0.31	<0.32	0.08 ± 0.01	0.15 ± 0.02	357 ± 42	X
A2537	0.2966	7.46 ± 0.34	<0.51	0.06 ± 0.02	0.12 ± 0.03	263 ± 78	X
RXJ0027.6+2616	0.3649	5.29 ± 0.58	<0.74	0.08 ± 0.01	0.15 ± 0.03	284 ± 53	X
Z5768	0.266	3.34 ± 0.23	<0.36	0.09 ± 0.01	0.17 ± 0.02	260 ± 35	X
S780	0.2357	6.67 ± 0.22	<0.38	0.09 ± 0.01	0.17 ± 0.01	354 ± 28	C
Radio Halos							
A2744	0.307	8.69 ± 0.32	18.62 ± 0.94	0.14 ± 0.02	0.26 ± 0.04	629 ± 89	C
A209	0.206	6.67 ± 0.12	1.99 ± 0.21	0.07 ± 0.00	0.13 ± 0.01	273 ± 15	X
A2163	0.203	15.22 ± 0.16	22.91 ± 1.16	0.12 ± 0.01	0.22 ± 0.02	711 ± 65	X
RXCJ2003.5-2323	0.3171	9.35 ± 0.53	10.71 ± 1.73	0.13 ± 0.01	0.24 ± 0.02	603 ± 45	C
A520	0.199	7.21 ± 0.16	2.45 ± 0.18	0.13 ± 0.01	0.24 ± 0.02	536 ± 55	C
A773	0.217	7.65 ± 0.19	1.48 ± 0.16	0.10 ± 0.01	0.20 ± 0.02	451 ± 35	X
A1758a	0.28	7.22 ± 0.17	5.75 ± 0.98	0.14 ± 0.02	0.27 ± 0.03	589 ± 65	C
A2219	0.2281	10.02 ± 0.25	5.63 ± 0.8	0.09 ± 0.01	0.17 ± 0.02	432 ± 40	X
A521	0.2475	6.41 ± 0.21	1.45 ± 0.13	0.12 ± 0.01	0.22 ± 0.02	462 ± 44	C
A697	0.282	9.66 ± 0.85	1.51 ± 0.14	0.09 ± 0.02	0.18 ± 0.03	458 ± 81	X
A1300	0.3075	7.85 ± 0.25	3.8 ± 1.43	0.14 ± 0.01	0.26 ± 0.03	595 ± 62	C
CL0016+16	0.541	9.42 ± 0.3	5.01 ± 0.31	0.10 ± 0.01	0.19 ± 0.01	472 ± 35	C
A1914	0.1712	11.38 ± 0.67	5.62 ± 0.43	0.12 ± 0.01	0.24 ± 0.02	657 ± 53	C
A665	0.1819	7.53 ± 0.17	2.51 ± 0.21	0.12 ± 0.01	0.24 ± 0.02	539 ± 39	X
A545	0.154	6.57 ± 0.09	1.41 ± 0.22	0.10 ± 0.01	0.19 ± 0.01	399 ± 24	X
Coma	0.0231	8.28 ± 0.13	0.76 ± 0.06	0.09 ± 0.01	0.18 ± 0.02	420 ± 45	R
A2256	0.0581	7.65 ± 0.63	0.85 ± 0.08	0.13 ± 0.01	0.25 ± 0.02	571 ± 56	R
Bullet	0.296	14.58 ± 0.4	23.44 ± 1.51	0.17 ± 0.02	0.31 ± 0.03	977 ± 99	C
A2255	0.0806	5.81 ± 0.2	0.81 ± 0.17	0.09 ± 0.01	0.18 ± 0.01	365 ± 22	R
A2319	0.0557	9.6 ± 0.3	2.45 ± 0.19	0.11 ± 0.01	0.21 ± 0.02	529 ± 46	R
MACSJ0717.5+3745	0.548	13.59 ± 0.68	52.48 ± 20.56	0.21 ± 0.02	0.40 ± 0.03	1206 ± 96	C
A1995	0.3186	7.58 ± 0.41	1.66 ± 0.23	0.08 ± 0.01	0.14 ± 0.03	326 ± 59	C
MACSJ1149.5+2223	0.544	13.14 ± 0.9	2.29 ± 0.95	0.15 ± 0.02	0.29 ± 0.03	853 ± 94	C
PLCKG171.9-40.7	0.27	12.78 ± 0.79	4.9 ± 1.35	0.08 ± 0.01	0.15 ± 0.01	438 ± 40	X
A754	0.0542	8.33 ± 0.05	0.63 ± 0.07	0.15 ± 0.01	0.28 ± 0.02	676 ± 46	R

Note.

^a (1) Cluster name. (2) Redshift. (3) Spectroscopic temperature extracted within a circular region of 1 Mpc radius. (4) Radio power at 1.4 GHz in units of 10^{24} W Hz⁻¹ (from Cassano et al. 2013). (5) Amplitude of two-dimensional X-ray brightness fluctuations. (6) 1D Mach number $M_{1D} = 2.3 \delta\rho/\rho$. (7) Turbulent velocity dispersion $\sigma_v = M_{3D} c_s$, with $c_s \approx (\gamma kT_{\text{Mpc}}/\mu m_p)^{0.5}$. (8) X-ray instrument used for the analysis (X = *XMM-Newton*, C = *Chandra*, R = *ROSAT*).

2.5. Turbulent Velocity

As shown by Arévalo et al. (2012), filtering the image at a given scale ℓ allows us to select fluctuations at a wavenumber

$k = \frac{1}{\sqrt{2\pi^2}\ell}$. For the present work, we choose a smoothing scale $\ell = 150$ kpc, corresponding to $k^{-1} = 660$ kpc. Such a scale roughly corresponds to the typical size of groups accreting onto

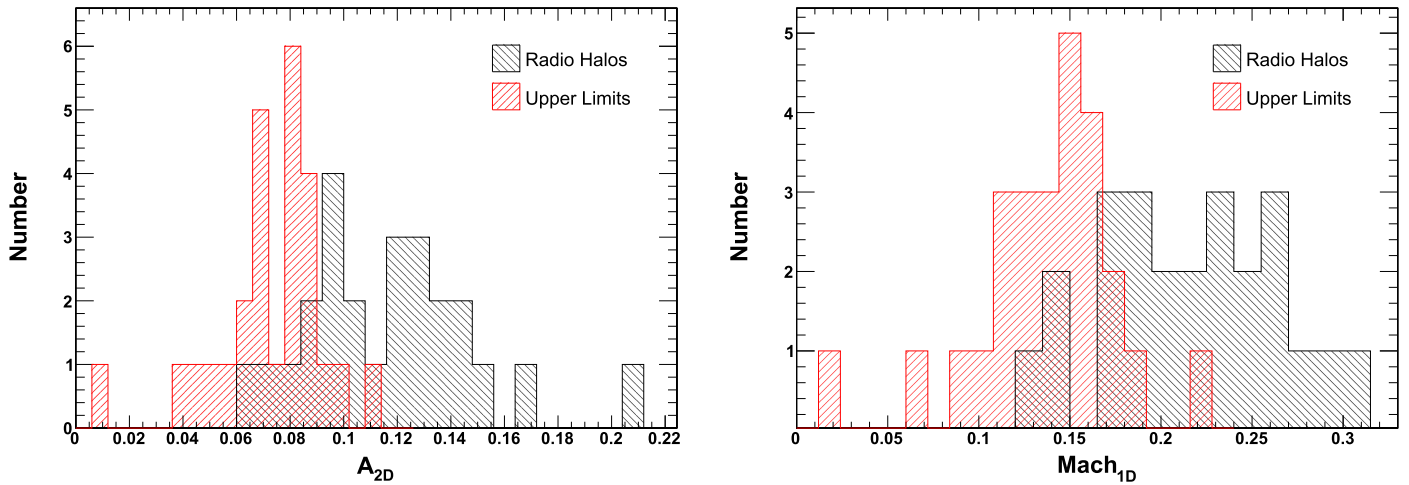


Figure 1. Left: distribution of two-dimensional fractional amplitude of surface-brightness fluctuations for radio halo clusters (black) and systems without a radio halo (red). Right: same as in the left panel, but for the 1D turbulent Mach number estimated from the deprojected density fluctuations.

clusters, i.e., it should be close to the injection scale of turbulent motions, at which the amplitude of fluctuations is expected to peak. Finally, the projected two-dimensional amplitude is converted into three-dimensional fluctuations $\delta\rho/\rho$ by using the best-fit beta model parameters and computing numerically the power induced by the projection of the emissivity distribution along the line of sight (see Equation (11) of Churazov et al. 2012). For consistency, we checked that the amplitudes recovered from *Chandra* and *XMM-Newton* data are in agreement when data from both telescopes are available. We found an excellent agreement between the two X-ray telescopes, giving us confidence that our method is robust.

As shown in Gaspari & Churazov (2013), in the subsonic regime the maximum amplitude of density fluctuations is linearly related to the turbulent Mach number $M = \frac{\sigma_v}{c_s}$, with σ_v the turbulent velocity dispersion and c_s the sound speed in the medium. The relation reads $M_{1D} \approx 2.3 \frac{\delta\rho}{\rho}$ or $M_{3D} \approx 4 \frac{\delta\rho}{\rho} (L/500 \text{ kpc})^{-0.25} \approx 3.7 \frac{\delta\rho}{\rho}$ for our choice of scale. The intrinsic scatter of the relation is expected to be $\sim 30\%$ (Zhuravleva et al. 2014). For each cluster, we compute the average spectroscopic temperature within the same circular region of 1 Mpc radius to estimate the average sound speed $c_s = (\gamma kT / \mu m_p)^{1/2}$ with $\gamma = 5/3$ and convert the Mach number into turbulent velocity.

3. Results

3.1. Amplitude of Surface-brightness Fluctuations

We estimated the two-dimensional amplitude of surface-brightness fluctuations A_{2D} at a fixed scale $k^{-1} = 660$ kpc and the corresponding Mach number and turbulent velocity dispersion for all clusters in our sample. The resulting values are provided in Table 1. In the left-hand panel of Figure 1, we show the distribution of the values of A_{2D} ($k = 1/660$ kpc) split between the radio halo and non-radio halo populations. We see that the radio halo clusters exhibit on average a higher value of A_{2D} than the upper limit population, with mean values $A_{2D,RH} = 0.117$ and $A_{2D,UL} = 0.072$ for the radio halo and upper limit populations, respectively. We performed a Kolmogorov-Smirnov (K-S) test to determine the probability

that the two sets of values are drawn from a common parent population. The K-S test returns a probability of 4×10^{-6} , meaning that the two distributions are different at 4.6σ confidence. If instead we focus only on the high-mass objects ($M_{500} > 5 \times 10^{14} M_\odot$) for which the existing radio data should be deep enough to detect a radio halo, the K-S probability reduces to 6×10^{-5} (4.0σ).

Using the deprojection factors estimated in Section 2.5, we then computed the amplitudes of 3D fluctuations of $\delta\rho/\rho$, which we converted into 1D Mach numbers using the relation $M_{1D} \approx 2.3 \delta\rho/\rho$ (Gaspari & Churazov 2013). The distribution of the values of M_{1D} is shown in the right-hand panel of Figure 1, again splitting the sample into the radio halo and radio quiet populations. A clear segregation is observed between the two populations, most radio halo clusters exhibiting a value of M_{1D} in the range 0.15–0.3 (which is consistent with the sample of Hofmann et al. 2016), whereas for the clusters where no radio emission was detected, the typical value of M_{1D} is ~ 0.12 and always less than ~ 0.2 . In this case, the K-S test returns a probability $p = 2 \times 10^{-7}$ that the two data sets are drawn from the same parent distribution, i.e., the result is significant at 5.2σ .

3.2. Correlation with Radio Power

For each cluster, we used the average temperatures reported in Table 1 to determine the sound speed in the medium. We then estimated the average three-dimensional turbulent velocity dispersion through the relation $\sigma_v = M_{3D} c_s \approx 3.7 c_s \delta\rho/\rho$ (Gaspari & Churazov 2013, Equation (22)). Note that our average temperatures are spectroscopic emission-weighted temperatures (e.g., Mazzotta et al. 2004), and thus our estimation of the sound speed does not exactly match the mean mass-weighted sound speed in the system. In Figure 2, we show the main result of the current study plotting the cluster 1.4 GHz radio power versus the estimated turbulent velocity dispersion. Again, a segregation is observed between the radio halo and upper limit populations. For radio halo clusters, the radio power appears to correlate with the velocity dispersion, with the Pearson coefficient $\rho = 0.80 \pm 0.02$ for 25 data points, indicating that the two quantities are significantly correlated.

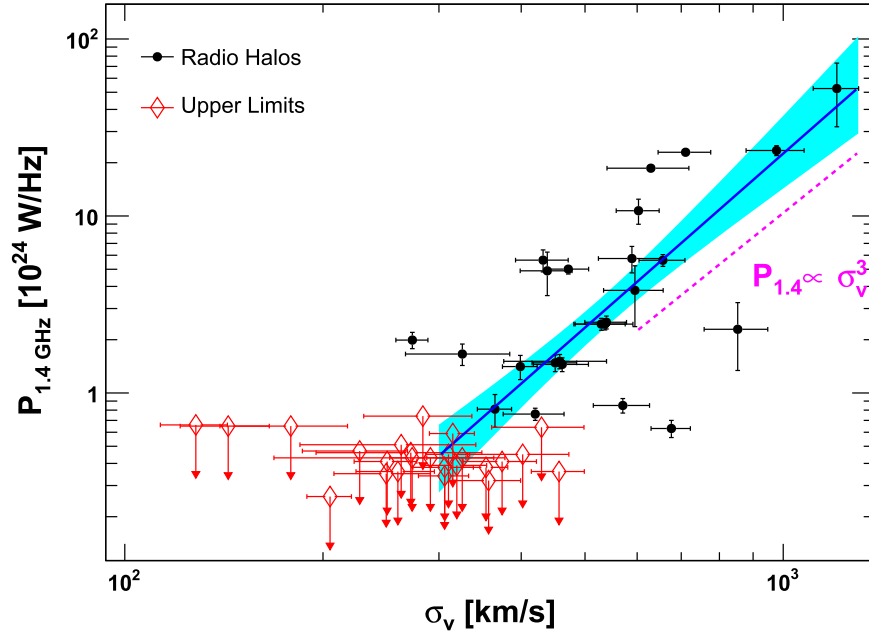


Figure 2. Radio power at 1.4 GHz as a function of velocity dispersion $\sigma_v = M_{3D} c_s$ for radio halo clusters (black) and radio upper limits (red). The blue curve and shaded area show the best fit to the data with a power law and its 1σ error envelope.

To describe the relation between σ_v and $P_{1.4}$, we fitted the data with a power law. Namely, we modeled the data as

$$\log\left(\frac{P_{1.4 \text{ GHz}}}{10^{24} \text{ W Hz}^{-1}}\right) = \log P_0 + \alpha \log\left(\frac{\sigma_v}{500 \text{ km s}^{-1}}\right). \quad (1)$$

We used the Bayesian routine `linmix_err` (Kelly 2007) to fit the data, taking both the data points and radio upper limits into account. The best-fitting values for the parameters are $\alpha = 3.27^{+0.71}_{-0.61}$ and $P_0 = 2.34^{+0.53}_{-0.49}$, with an intrinsic scatter $\sigma_{\ln P|\sigma_v} = 0.44^{+0.18}_{-0.13}$. We note that the relation between the projected surface-brightness fluctuations and three-dimensional velocity field should exhibit a substantial scatter; thus, the scatter of $\sim 40\%$ observed here is not surprising. Note that the relation reported in Figure 2 includes a hidden correlation, as the radio power is a steep function of cluster mass (e.g., Cassano et al. 2013) and the sound speed c_s also depends on cluster mass, albeit much more mildly.

4. Discussion

The main results of this work have potentially important implications on the origin of radio halos, which we discuss here. While previous works have found that radio halos arise almost only in morphologically disturbed systems (Buote 2001; Cassano et al. 2010, 2013; Cuciti et al. 2015), as evidenced, e.g., by their centroid shifts, our analysis shows that radio halo clusters remain on average more perturbed even after subtracting the large-scale gas distribution and masking substructures. We expect that surface-brightness fluctuations on top of the large-scale gas distribution trace residual gas motions. Numerical simulations predict that ICM turbulence should be the dominant source of residual motions and hence of ICM fluctuations (e.g., Lau et al. 2009; Vazza et al. 2009), although other sources of fluctuations (e.g., shocks, cold fronts, ram-pressure stripping) play a role to some extent. Our work thus establishes a connection between turbulent motions and particle acceleration in galaxy clusters, which corroborates a

key prediction of the turbulent re-acceleration scenario (e.g., Brunetti et al. 2001).

In case the density fluctuations measured here can be entirely ascribed to turbulence, Figure 2 shows that the recovered turbulent velocity dispersion correlates with the radio power at 1.4 GHz, which can be interpreted in the framework of the turbulent re-acceleration scheme. The energy rate per unit volume related to turbulence can be estimated as

$$P_{\text{turb}} \approx 9.8 \times 10^{-25} \left(\frac{\sigma_v}{500 \text{ km s}^{-1}}\right)^3 \left(\frac{n_{\text{gas}}}{10^{-2} \text{ cm}^{-3}}\right) \times \left(\frac{L_{\text{inj}}}{500 \text{ kpc}}\right)^{-1} \text{ erg s}^{-1} \text{ cm}^{-3}. \quad (2)$$

In the classical Kolmogorov model, the kinetic power across the cascade is preserved, regardless of the specific dissipation mechanism. While the specific scales and mechanisms for the ultimate dissipation of turbulent motions in the ICM are currently unknown (see, e.g., Brunetti & Lazarian 2011), it is reasonable to assume that a fraction of the kinetic power through the cascade goes into the re-acceleration of radio emitting electrons. The exact relation between P_{radio} and P_{turb} is uncertain as it depends on many unknown factors. Yet, our data support (within 1σ) a simple $P_{\text{radio}} \propto \sigma_v^3 \propto P_{\text{turb}}$, suggesting that A_{2D} can capture the bulk of the energy involved in the re-acceleration process. However, the variety of mechanisms superseding the dissipation of plasma turbulence in the ICM (heating, magnetic field amplification, acceleration of cosmic-ray electrons and protons) is yet unknown; thus, the relation between the emitted power and the turbulent kinetic energy should be further investigated.

Although the segregation between radio halo and radio quiet clusters in Figures 1 and 2 is clear, there are obviously a number of outliers. Among these, we note the cases of A209 and A781. A209 is the radio halo cluster with the lowest values of M_{ID} (0.13 ± 0.01) and σ_v ($273 \pm 15 \text{ km s}^{-1}$). Such values

would classify this object firmly within the “upper limit” quadrant. A209 could be a fading radio halo in which turbulence is now being dissipated. Alternatively, projection effects may hide large-scale fluctuations, which would lead to an underestimate of the turbulent Mach number. Conversely, A781 exhibits a fairly high Mach number $M_{1D} = 0.22 \pm 0.02$ and $\sigma_v = 457 \pm 42$. This system does not host a radio halo (Venturi et al. 2011), although residual radio emission at the level of 15–20 mJy at 320 MHz could indicate a steep-spectrum radio halo. One possibility is that this system is in the early stage of a merger and that the turbulent cascade is not fully developed, such that the onset of particle acceleration at the microscale has not yet taken place.

As a word of caution, we note that the conversion between 2D amplitude and turbulent velocity dispersion assumes that all the observed fluctuations are induced by turbulent motions. Additional perturbing phenomena such as unresolved gas clumps or shock fronts will lead to an increase in the fractional amplitude of density perturbations. Although deviating regions such as clear substructures were masked during this analysis, the values reported here might overestimate the intrinsic turbulent Mach number.

Future X-ray imaging spectrometers such as *ATHENA* (Nandra et al. 2013) or the *Hitomi* recovery mission *XARM* will allow us to measure the velocity dispersion directly and will set strong constraints on the models describing the acceleration of relativistic electrons by turbulence in the ICM.

5. Conclusions

In this Letter, we presented measurements of the projected amplitude of gas density fluctuations in the ICM at a single scale $k^{-1} = 660$ kpc and, for the first time, correlated it with the radio emission for a sample of 51 galaxy clusters. Our results can be summarized as follows:

1. We observe a clear segregation between the level of ICM fluctuations between the clusters exhibiting a radio halo and the ones where no radio emission has been detected. The difference between the two populations is significant at the 4.6σ level.
2. If the measured fluctuations are interpreted as being entirely due to residual gas motions, the 1D Mach number of turbulent motions is found to be larger in the radio halo populations by about a factor of two compared to the clusters where no radio emission was detected.
3. The turbulent velocity dispersion $\sigma_v = M_{3D} c_s$ of radio halo systems correlates with the observed radio power (correlation coefficient 0.80 ± 0.02). The best-fit relation reads $P_{1.4} \propto \sigma_v^{3.3 \pm 0.7}$ with 44% scatter. Thus, the radio power is nearly proportional to the turbulent energy rate $P_{\text{turb}} \propto \sigma_v^3$.
4. Provided that surface-brightness fluctuations are probing gas motions in the ICM and that the measured turbulence cascades with no break down to the resonant scales where particles can be accelerated, our results corroborate stochastic acceleration via turbulence as the most likely mechanism to boost the emergence of radio halos in galaxy clusters.

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