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The all-sky PLATO input catalogue[★]

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ABSTRACT

Context. The ESA PLANetary Transits and Oscillations of stars (PLATO) mission will search for terrestrial planets in the habitable zone of solar-type stars. Because of telemetry limitations, PLATO targets need to be pre-selected.

Aims. In this paper, we present an all sky catalogue that will be fundamental to selecting the best PLATO fields and the most promising target stars, deriving their basic parameters, analysing the instrumental performances, and then planing and optimising follow-up observations. This catalogue also represents a valuable resource for the general definition of stellar samples optimised for the search of transiting planets.

Methods. We used *Gaia* Data Release 2 astrometry and photometry and 3D maps of the local interstellar medium to isolate FGK ($V \leq 13$) and M ($V \leq 16$) dwarfs and subgiant stars.

Results. We present the first public release of the all-sky PLATO input catalogue (asPIC1.1) containing a total of 2 675 539 stars including 2 378 177 FGK dwarfs and subgiants and 297 362 M dwarfs. The median distance in our sample is 428 pc for FGK stars and 146 pc for M dwarfs, respectively. We derived the reddening of our targets and developed an algorithm to estimate stellar fundamental parameters (T_{eff} , radius, mass) from astrometric and photometric measurements.

Conclusions. We show that the overall (internal+external) uncertainties on the stellar parameter determined in the present study are ~ 230 K (4%) for the effective temperatures, $\sim 0.1 R_{\odot}$ (9%) for the stellar radii, and $\sim 0.1 M_{\odot}$ (11%) for the stellar mass. We release a special target list containing all known planet hosts cross-matched with our catalogue.

Key words. catalogs – astrometry – techniques: photometric – planets and satellites: terrestrial planets – stars: fundamental parameters – ISM: structure

1. Introduction

The PLATO mission (Rauer et al. 2014, 2016) is the third medium-class mission in the ESA Cosmic Vision programme. Its ambitious goals are the detection and characterisation of terrestrial planets around solar-type stars as well as the study of the properties of host stars. The focus of PLATO is on long-orbital-period terrestrial planets that are also in the habitable zone of solar-type stars. PLATO is designed for the discovery of Earth-analogue planets under potentially favourable conditions for the development of life.

PLATO will achieve its goals by detecting the low-amplitude dips in stellar brightness produced by planets transiting in front of the disk of their parent stars. The mission is designed and optimised to continuously monitor two Long-duration Observation Phase (LOP) fields, each one for up to three years, and has the capability to monitor some additional fields during shorter times, up to three months each (‘Step and stare’ Observation Phase, SOP). The precise observing strategy will be decided at a later stage (two years before launch at the latest), taking advantage of the advances in the field at that time.

PLATO will employ an array of 26 cameras (hereafter ‘cameras’ signifies the telescope optics and the focal plane assembly,

including all the ancillary devices like baffling, electronics, etc.) made up with refractive 20 cm-class optical elements with a very wide field (about 40 degrees in diameter) and a 12cm diameter equivalent aperture.

Each camera field of view is almost covered by an array of detectors totaling 1037 square degrees, and 24 cameras are arranged in four groups of six cameras each, co-aligned to four slightly overlapping fields of view. The S/N of each star will therefore be dependent on its precise location within the overall PLATO field of view. The combined field of view of the system of cameras amounts to a total of about 2100 square degrees (Ragazzoni et al. 2015; Magrin et al. 2018). The remaining two cameras are optimised for fast monitoring of very bright stars, their colour measurements (e.g., Grenfell et al. 2020), and fine guidance and navigation. Their focal plane is equipped with frame transfer detectors that allow for a coverage somewhat larger than 600 degrees. The geometry of the covered patches in the sky is kept symmetric, under nominal conditions, for 90 degree rotations around the line of sight to keep the same coverage of the field of view throughout a pointing and to allow continuous in-flight calibration among different cameras. Compliance with the performance specifications is computed using models for the degradation of the performances over the mission lifetime, mainly due to a moderate degradation of the throughput, due to the choice of rad-hardened glasses for the optical elements most exposed to cosmic radiation (Magrin et al. 2016; Corso et al. 2018), and to the worst scenario in which the detectors or electrical systems of two of the cameras could fail by the End Of Life (EOF).

[★] The catalogue described in this article is only available at MAST as a High Level Science Product via <https://dx.doi.org/10.17909/t9-8msm-xh08> and <https://archive.stsci.edu/hlsp/aspic>, in the SSDC tools page (<https://tools.ssdsc.asi.it/asPICtool/>) and at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via <http://cdsarc.u-strasbg.fr/cgi-bin/viz-bin/cat/J/A+A/653/A98>

Because of limits on the telemetry rate from the satellite to Earth, it is not possible to download full CCD images at a sufficient rate to detect planets; therefore, it is necessary to pre-select PLATO targets. For a subsample of targets, customised ‘postage stamps’ (imagettes) are downloaded to ground; for the remaining selected stars, centring and photometry will be performed on board and then downloaded to ground. Other, previous space missions followed the same approach. The CoRoT Exo-Dat catalogue (Deleuil et al. 2009), the *Kepler* input catalogue (KIC; Brown et al. 2011), the K2 Ecliptic Plane Input Catalogue (EPIC; Huber et al. 2016), and the TESS input catalogue (TIC; Stassun et al. 2018, 2019) were produced to select the best targets for exoplanet searches.

In this paper, we present the all-sky PLATO input catalogue (asPIC), whose main purpose is to serve as the primary database for the selection of PLATO targets. Establishing the properties of PLATO host stars is crucial to determining the properties of transiting planets; for example, the uncertainty on their radii is directly related to the uncertainty on the host-star radius. The classical properties of planet hosts are also essential to modelling the stars using asteroseismology. Also, the habitability of a given planet depends on the properties of the host star, among other factors. Moreover, knowing the values of stellar parameters permits the statistical study of the frequency of planets as a function of stellar type, metallicity, environment, and so on (e.g., Howard et al. 2012; Bryson et al. 2020).

The construction of such a catalogue not only requires the identification, across the entire sky, of stars with the required spectral type and luminosity class, but also requires that noise constraints be accounted for to facilitate the detection of transiting planets around the selected sources. Hence limits on the magnitude range of the stellar samples are set according to mission requirements and the degree of contamination of each target from neighbouring sources must be evaluated. The PLATO pixel scale will be 15 arcsec pix⁻¹ on axis, which lies between that of *Kepler* (4 arcsec pix⁻¹; Borucki et al. 2010) and that of TESS (21 arcsec pix⁻¹; Ricker et al. 2014).

The optical quality of the PSF, although slightly variable over the field of view and depending on the thermal status of the cameras, is expected to encompass an area of a few pixels (see also Gullieuszik et al. 2016; Umbriaco et al. 2018). Consequently, it is expected that several sources, especially in crowded fields, will be blended into a single brightness measurement area and they may become the cause of false positive signals (e.g., Santerne et al. 2015; Fressin et al. 2013). More specifically, the asPIC will be used to:

1. Select the optimal PLATO observing fields;
2. Select all FGKM dwarf and subgiant stars satisfying the magnitude and signal-to-noise constraints defined in the PLATO stellar sample requirements;
3. Estimate basic stellar parameters for all targets, such as temperatures, radii, and masses;
4. Identify known variable stars, binaries, members of multi-planet systems, active stars with bitmaps, and summarize the information available from existing catalogues, that is, mainly *Gaia* and TESS releases;
5. Supply a list of all target contaminants up to a specific angular distance from the target and up to a limiting magnitude, for the targets in the fields to be observed, in order to understand their impact on photometric performances and because of the challenge that background contaminants pose to planetary candidate validation;
6. Guide the organisation and optimisation of ground-based follow-up strategies.

Previous attempts to produce this catalogue showed that this is not a trivial task. In particular, it is hard to robustly discriminate between dwarfs and evolved stars without accurate knowledge of the absolute luminosities of sources. Reduced proper motions were employed in the past to this purpose (Nascimbeni et al. 2016), but this technique is now superseded by the availability of trigonometric parallaxes from *Gaia* Data Release 2 (DR2) (Gaia Collaboration 2016, 2018)¹, essentially for all targets of PLATO interest, permitting accurate determination of distances and therefore of absolute luminosities.

After ESA adoption (in June 2017), PLATO recently entered the full industrial development phases C-D (March 2020) to be ready for launch in 2026. At this development stage, the selection of the PLATO LOP fields is not yet finalised (it will be two years before launch). The present work is focused on the asPIC. This catalogue is made publicly available to the community and it will be updated whenever relevant changes occur, including new *Gaia* releases.

After briefly describing the mission requirements defining PLATO stellar samples in Sect. 2, we describe the target selection criteria in Sect. 3. Then, in Sect. 4, we describe how we account for reddening and absorption in our analysis. Our stellar parameter pipeline is presented in Sect. 5. In Sect. 6, we compare our results with those published in various literature sources. In Sect. 7, we describe a special list of known planet host stars included in the catalogue release. Section 8 provides a brief summary of the present work.

2. PLATO stellar sample requirements

The four main PLATO stellar samples (named P1, P2, P4, and P5)² will be observed in the wavelength range between 500 nm and 1000 nm, and an additional colour sample will be observed in two broad blue and red spectral bands spanning the 500–675 and 675–1000 nm wavelength ranges, respectively (Corso et al. 2018). The science requirements for the PLATO stellar samples are listed below and summarised in Table 1 according to the ESA Science Requirements Document (SciRD, PTO-EST-SCI-RS-0150_SciRD_7_0). The definition of the PLATO samples requires knowledge of the visual apparent magnitude *V* because the SciRD adopts this magnitude as reference to identify sufficiently bright targets for ground-based spectroscopic follow-up. The *V* magnitude we used to select the targets comes from our own calibrated transformation of *Gaia* DR2 magnitudes and colours to the Johnson *V* band (Eq. (1)). This choice ensures the maximum homogeneity of the *V* band photometry. The calibration procedure is described in Appendix A. The colour transformation relation is

$$(G - V)_0 = c_0 + c_1(G_{BP} - G_{RP})_0 + c_2(G_{BP} - G_{RP})_0^2 + c_3(G_{BP} - G_{RP})_0^3 + c_4(G_{BP} - G_{RP})_0^4 + c_5(G_{BP} - G_{RP})_0^5, \quad (1)$$

where *G*, *G_{BP}*, *G_{RP}* are the *Gaia* DR2 magnitudes and the null subscript indicates colours dereddened with the procedure described in Sect. 4. The best-fit coefficients of the transformation relation are reported in Table 2.

¹ Although *Gaia* EDR3 has now been released (Gaia Collaboration 2021) this work is based on *Gaia* DR2. Future updates of the asPIC will be based on *Gaia* DR3.

² For historical reasons, the sample P3 has been eliminated, but the numbering of the PLATO samples was left unchanged.

Table 1. Summary of science requirements for the PLATO stellar samples.

	P1	P2	P4	P5	Colour sample
Stars	$\geq 15\,000$ (goal 20 000)	≥ 1000	≥ 5000	$\geq 245\,000$	300
Spectral type	Dwarf and subgiants F5-K7	Dwarf and subgiants F5-K7	M Dwarfs	Dwarf and subgiants F5-late K	Anywhere in the HR diagram
Limit V	11	8.5	16	13	–
Random noise (ppm in 1 h)	< 50	< 50	–	–	–
Wavelength (nm)	500–1000	500–1000	500–1000	500–1000	Blue (500–675) and red (675–1000) spectral bands

Table 2. Coefficients of the polynomial best-fit model relating the $(G - V)_0$ and $(G_{BP} - G_{RP})_0$ colours and residuals of the fit (see also Appendix A).

	c_0	$(G_{BP} - G_{RP})_0$ c_1	$(G_{BP} - G_{RP})_0^2$ c_2	$(G_{BP} - G_{RP})_0^3$ c_3	$(G_{BP} - G_{RP})_0^4$ c_4	$(G_{BP} - G_{RP})_0^5$ c_5	σ_{fit}
$(G - V)_0$	–0.17276	0.47885	–0.71953	0.24374	–0.04458	0.00317	0.02745

Notes. The relation is valid within $0.5 < (G_{BP} - G_{RP})_0 < 5$.

2.1. Stellar sample 1 (P1)

- The total number of targets in stellar sample 1 (cumulative over all sky fields) shall be at least 15 000 dwarf and subgiant stars of spectral types from F5 to K7, with a goal of 20 000.
- The dynamic range of stellar sample 1 shall be $V \leq 11$.
- The random noise level for stellar sample 1 shall be below 50 ppm in 1 h.
- In stellar sample 1, the proportion of brighter targets ($V \leq 10.5$) shall be maximised.
- Stellar sample 1 shall be observed during a LOP.

2.2. Stellar sample 2 (P2)

- The total number of targets in stellar sample 2 (cumulative over all sky fields) shall be at least 1000 dwarf and subgiant stars of spectral types from F5 to K7.
- The dynamic range of stellar sample 2 shall be $V \leq 8.5$.
- The random noise level for stellar sample 2 shall be below 50 ppm in 1 h.
- Stellar sample 2 shall be observed during a LOP.

2.3. Stellar sample 4 (P4)

- The total number of targets in stellar sample 4 (cumulative over all sky fields) shall be at least 5000 cool late-type dwarfs (M dwarfs) monitored during a LOP.
- The dynamic range of stellar sample 4 shall be $V \leq 16$.

2.4. Stellar sample 5 (P5)

- The total number of targets in stellar sample 5 (cumulative over all sky fields) shall be at least 245 000 dwarf and subgiant stars of spectral types from F5 to late K.
- The dynamic range of stellar sample 5 shall be $V \leq 13$.
- Stellar sample 5 shall be observed during a LOP.

2.5. Colour sample

- Part of the payload must provide photometry in at least two separate colour broad-bands.
- PLATO shall have the capability to observe 300 stars (in at least two pointings) located anywhere on the HR diagram, with colour information, imagerettes, and a sampling time equal to 2.5 s.

The above stellar samples are defined only across the two selected PLATO LOP fields. The final decision on these fields will be taken two years before launch. So far, two preliminary LOP fields have been identified. The P1 sample is a subsample of P5. However, the P1 sample is also S/N-limited and, because of the particular configuration of PLATO pointing, it can be defined only once the LOP fields are selected.

In this work, we construct an all-sky catalogue, focusing on only two main PLATO stellar samples. One sample is related to FGK (limited to F5) dwarf and subgiant stars with apparent visual magnitude $V \leq 13$ (and formally corresponds to the P5 sample as defined above), while the other sample is related to M-dwarfs with $V \leq 16$ (which corresponds to P4). For the sake of clarity, hereafter we refer to these two samples as the FGK sample and the M sample, respectively. We note that although the asPIC is an all-sky catalogue, PLATO will not point towards the Galactic plane.

3. Selection criteria

The selection of the PLATO targets was performed using the absolute colour-magnitude diagram (CMD) corrected for reddening. First, we employed theoretical models based on Galactic simulations to identify the region of the diagram where our targets are located (Sect. 3.1). On the same diagram, we then represented stars that, based on available spectroscopic and/or astrometric observations, have the properties required by the SciRD (Sect. 3.2). Finally, by taking into account both the theoretical

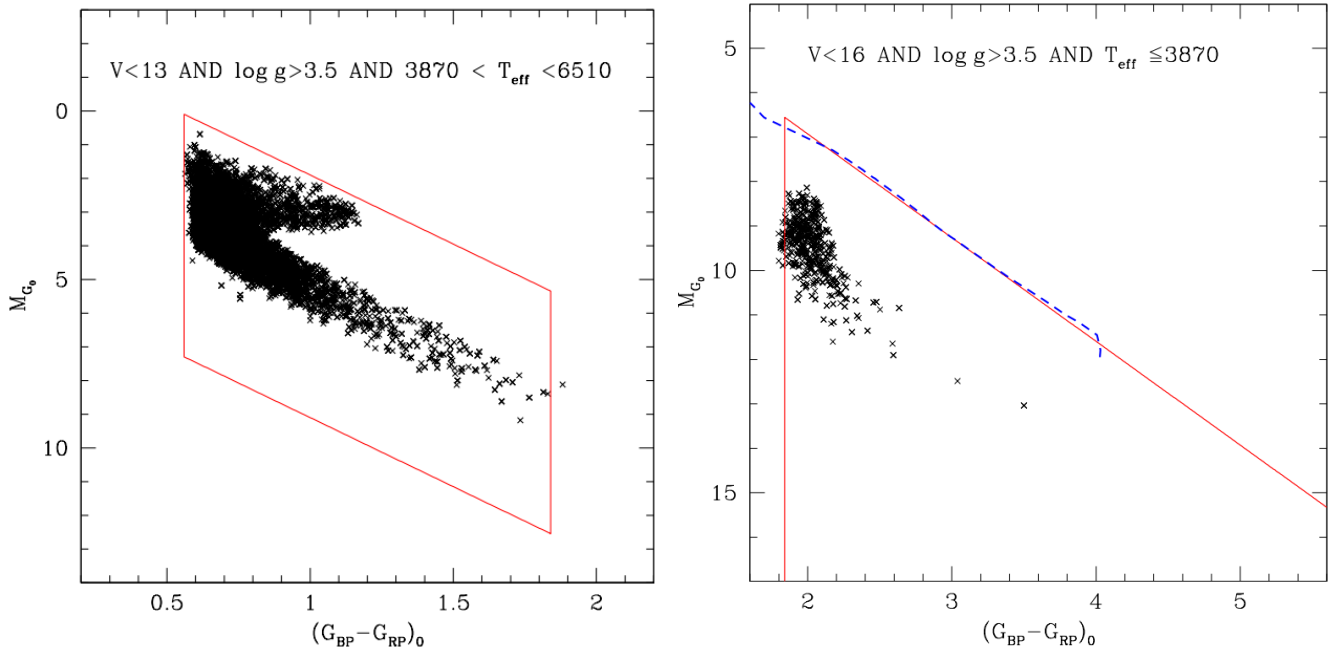


Fig. 1. Absolute magnitude vs. reddening-free colour obtained from TRILEGALv1.6 simulations. Crosses represent simulated stars from a total area of 100 square degrees in the NPF and the SPF, satisfying the conditions described by the labels at the top of each plot. The *left and right panels* also show the analytical selection (red lines) described in Sect. 3.3 for FGK dwarfs and subgiants (FGK sample) and for M dwarfs (M sample), respectively. In the diagram on the right, the blue dashed line represents a 10 Myr solar metallicity isochrone from the Padova database (Bressan et al. 2012).

and empirical distributions of stars in the CMD, we proceeded by defining an analytical selection region, and estimated its degree of completeness and contamination (Sect. 3.3). Finally, in Sect. 3.5 we show that our selection does not imply any bias in the metallicity distribution of the selected stars with respect to a sample of local stars (Sect. 3.4) and present the asPIC1.1 catalogue resulting from our selection.

As summarised in Sect. 2, the PLATO samples considered here include an FGK dwarf and subgiant sample (limited to F5), and an M dwarf sample. Considering the definitions reported in Table 5 of Pecaut & Mamajek (2013), spectral type F5 corresponds to $T_{\text{eff}} = 6510$ K, or to an unreddened *Gaia* colour $(G_{\text{BP}} - G_{\text{RP}})_0 = 0.587$, and spectral type M0 corresponds to $T_{\text{eff}} = 3870$ K or unreddened $(G_{\text{BP}} - G_{\text{RP}})_0 = 1.84$, where G_{BP} and G_{RP} are the blue and red *Gaia* magnitudes, respectively. In addition, dwarfs and subgiants have $\log g > 3.5$.

3.1. Theoretical models

To determine and calibrate our selection criteria in an observational CMD, we used Galactic simulations from TRILEGALv1.6, (Girardi et al. 2005)³. TRILEGAL is a population synthesis code for simulating the stellar photometry of any Galaxy field. It allows the user to simulate the photometry in several different broad-band photometric systems. Each simulation we performed corresponds to an area of ten square degrees. The simulations were centred at $(l, b) = (65^\circ, 30^\circ)$ and $(l, b) = (253^\circ, -30^\circ)$, and on four surrounding locations ten degrees apart from the central fields in Galactic longitude or latitude. The two central fields correspond to the two provisional PLATO long-duration fields in the Northern and Southern Galactic hemispheres, that is the North PLATO field and the South PLATO field (NPF and SPF, respectively). We considered the *Gaia* DR2

passbands of Maíz Apellániz & Weiler (2018) and determined the apparent visual magnitude V using the relation reported in Appendix A. We corrected the apparent magnitudes of each simulated star for extinction using the V band extinction tabulated in the simulations and the relations given in Sect. 5.2 in order to convert the extinction in the V band to the extinction in the *Gaia* bands. Figure 1 shows the intrinsic colour (reddening free) and absolute magnitude obtained from TRILEGALv1.6 simulations. The left panel shows stars with $V < 13$, $\log g > 3.5$ and $3870 \text{ K} < T_{\text{eff}} < 6510 \text{ K}$ corresponding to our adopted limits for the FGK sample, while the right panel shows stars with $V < 16$, $\log g > 3.5$ and $T_{\text{eff}} \leq 3870 \text{ K}$ (M sample). Figures 2 and 3 show the distributions of effective temperatures (left panel) radii (middle) and masses (right) of the simulated stars. Magnitude limited samples, like the ones we are considering, are dominated by hot dwarfs and subgiants, because these stars are intrinsically more luminous and therefore observable in a large volume.

We also repeated the simulations for fields located in the Galactic plane or at the Galactic poles. The intrinsic colours and magnitudes of the selected stars are perfectly compatible with the ones shown in Fig. 1. Larger uncertainties in the reddening estimates occur for the Galactic plane and would lead to higher levels of contamination, particularly for hot dwarfs and subgiants (see also Sect. 3.3).

3.2. Empirical control samples

For FGK dwarfs and subgiants, as an empirical control sample, we used stars from the RAVE DR5 survey (Kunder et al. 2017). RAVE provides two sets of stellar parameters: those that are direct outputs of the stellar parameter pipeline and the calibrated parameters. Because the direct outputs are based on a grid of synthetic spectra which may not match real spectra over the entire domain of the parameter space, a calibration based on reference stars is performed in order to minimise systematic

³ <http://stev.oapd.inaf.it/cgi-bin/trilegal>

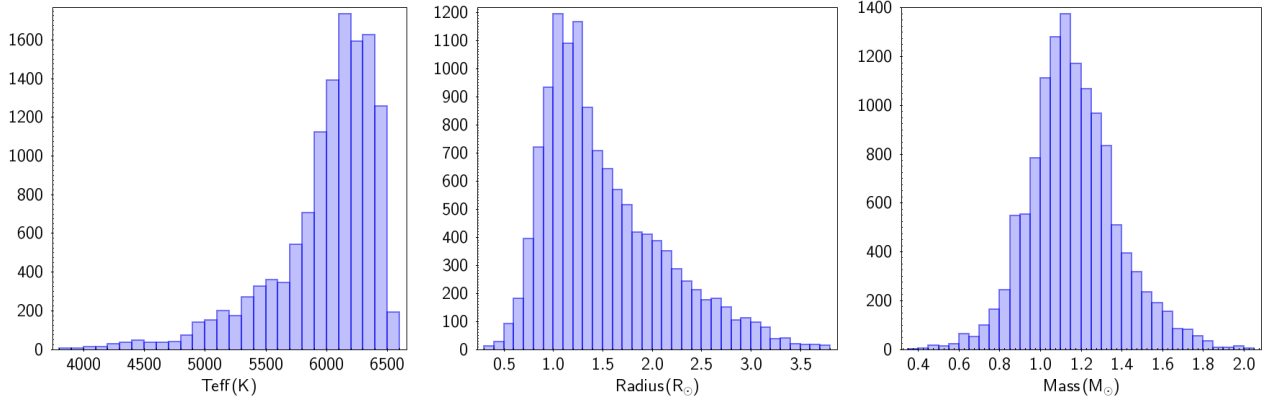


Fig. 2. Distribution of effective temperatures (*left panel*), stellar radii (*middle*), and masses (*right*) for FGK sample stars from TRILEGAL simulations.

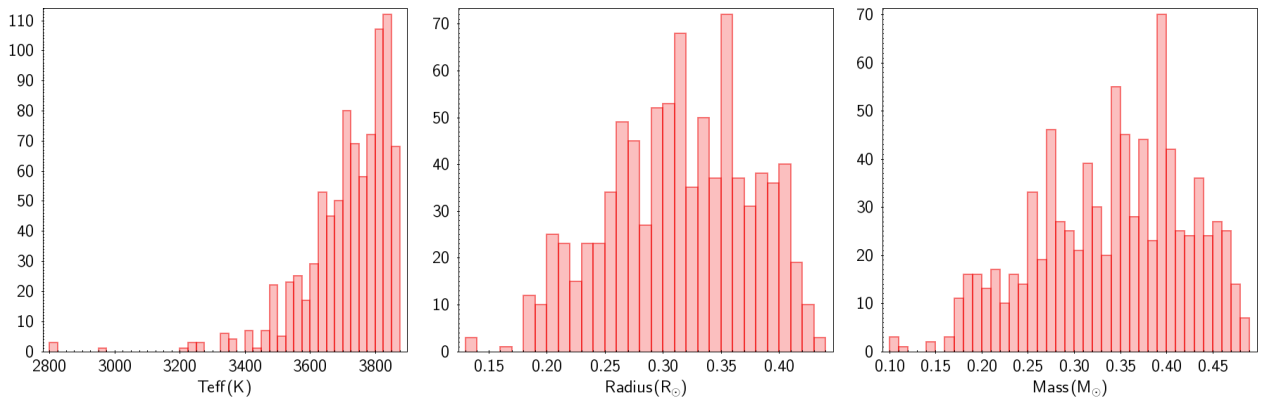


Fig. 3. Distribution of effective temperatures (*left panel*), stellar radii (*middle*), and masses (*right*) for M sample stars from TRILEGAL simulations.

offsets. For example, surface gravities are calibrated using both the asteroseismic $\log g$ values of 72 giants from [Valentini et al. \(2016\)](#) and the *Gaia* benchmark dwarfs and giants ([Heiter et al. 2015](#)). We then adopted the calibrated parameters. We considered stars with RAVE parameter $\text{ALGOCONV}=0$ and $S/N > 50$ to ensure the good quality of extracted parameters. A total of 122 167 stars were used. We analysed the distribution of RAVE stars in the absolute colour magnitude diagram constructed using [Bailer-Jones et al. \(2018\)](#) distances. We corrected the apparent magnitudes for extinction using the procedure presented in Sect. 4.4. To identify the selection regions of dwarfs and subgiants in the $(G_{\text{BP}} - G_{\text{BP}})_0$, $M_{G,0}$ plane, we considered a grid of $0.05 \text{ mag} \times 0.5 \text{ mag}$ bins in colour and magnitude across the plane. We considered stars with $V < 13$, where the V band magnitude was calculated with the calibration relationship reported in Appendix A. Stars for which RAVE parameters did not satisfy the conditions $\log g > 3.5$ and $3870 \text{ K} < T_{\text{eff}} < 6510 \text{ K}$ were considered contaminants. We then selected those grid rectangles with at least 50 stars and for which the number of dwarfs and subgiants was larger than 75% of the total number of stars. Figure 4 (left panel) shows the result of this procedure. The black rectangles are the regions belonging to the dwarf+subgiant class.

For M-dwarfs, we used The REsearch Consortium On Nearby Stars (RECONS⁵) catalogue ([Henry et al. 2018](#)), and

the [Mann et al. \(2015\)](#) catalogue. Among other information, the RECONS catalogue reports the apparent visual Johnson magnitude of the targeted stars. Whenever multiple measurements for the same object were present in the catalogue, we averaged their visual magnitudes. From the initial list of 348 entries we obtained a list of 289 measurements for the individual objects. The catalogue was cross-matched with *Gaia* DR2 and we then selected stars with $V < 16$ that satisfy Eqs. (1) and (2) of [Arenou et al. \(2018\)](#) *Gaia* astrometric and photometric quality criteria and are classified as M-dwarfs⁶, restricting the list to 195 stars. We then calculated the reddening as described in Sect. 4. The right panel of Fig. 4 shows the absolute colour–magnitude diagram corrected for reddening of the selected stars (open black circles) of this sample.

[Mann et al. \(2015\)](#) constructed a sample of nearby late K and M dwarfs (all within 40 pc from the Sun) for which optical and near-infrared spectra were obtained. We used the effective temperatures derived by these latter authors to isolate only M dwarfs by imposing $T_{\text{eff}} \leq 3870 \text{ K}$. *Gaia* DR2 was cross-matched with this catalogue and we restricted the sample to those stars satisfying Eqs. (1) and (2) of [Arenou et al. \(2018\)](#). The apparent visual magnitude in the Johnson system (necessary for our selection, and not reported in the catalogue) was calculated as described in Appendix A. By means of these selection criteria the initial sample of 179 stars was reduced to 153 stars represented by the magenta asterisks in Fig. 4.

⁴ We also compared our reddening $E(B - V)$ with the one calculated in RAVE and found perfectly consistent results: $\Delta E(B - V) = (-0.01 \pm 0.02)$.

⁵ <http://www.recons.org/>

⁶ RECONS does not report the effective temperature, but it reports the spectral type of the stars.

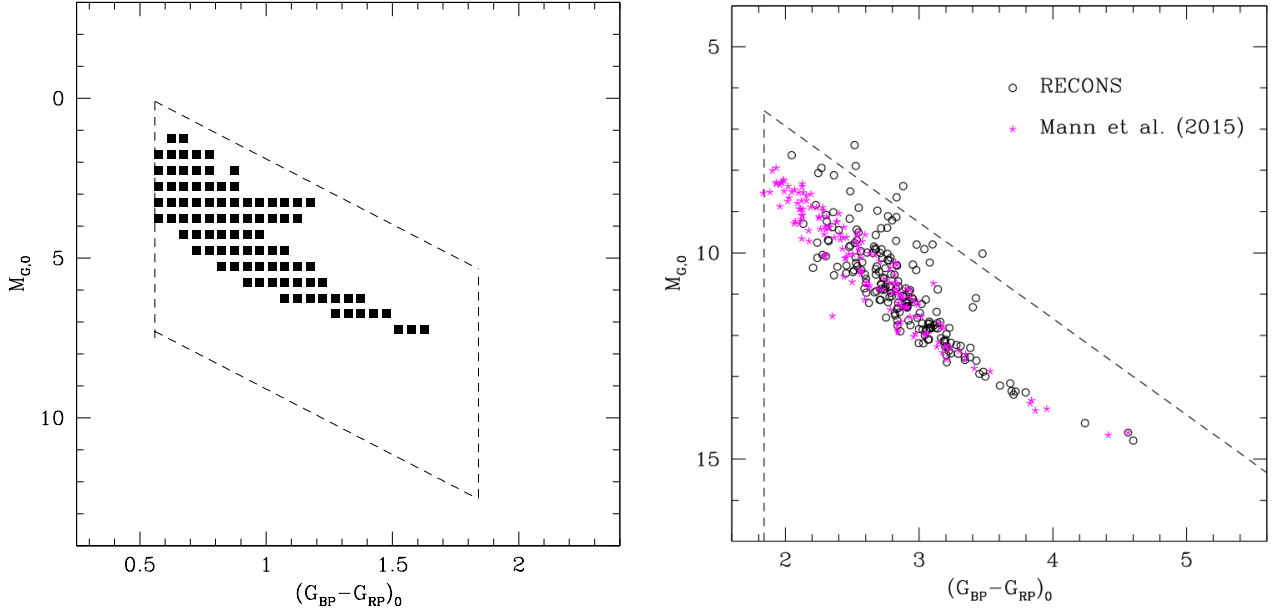


Fig. 4. *Left panel:* black rectangles represent FGK dwarfs and subgiants from RAVE (see Sect. 3.2, Kunder et al. 2017). *Right panel:* M-dwarfs (Henry et al. 2018, open circles) and (Mann et al. 2015, magenta asterisks) catalogues. Dashed lines in both panels show the analytical selection defined in Sect. 3.3 for the FGK (*left*) and M stellar samples (*right*).

3.3. Analytical selection

For the selection of M and FGK targets from the *Gaia* CMD, we then proceeded by defining a set of inequations. The red colour limit corresponding to $(G_{BP} - G_{RP})_0 = 1.84$ was used to separate M-dwarfs from FGK dwarfs and subgiants following Pecaut & Mamajek (2013), while the blue colour limit of the FGK samples was set at $(G_{BP} - G_{RP})_0 = 0.56$ corresponding to a spectral type between F4V-F5V according to Pecaut & Mamajek (2013). This limit is bluer than the $(G_{BP} - G_{RP})_0 = 0.587$ proposed by Pecaut & Mamajek (2013) for F5 stars in order to (partially) account for reddening errors. By using TRILEGAL simulations, we perturbed the reddening correction according to the expected reddening uncertainties derived from our map (see Sect. 4), and estimated the number of contaminants (stars with spectral type earlier than F5) and targets in narrow colour intervals of 0.05 mag. Because our samples are magnitude limited, hotter stars are more numerous than cool ones, because they are intrinsically more luminous (see e.g., Fig. 2, left panel). The major source of contamination comes from massive dwarfs and subgiants of spectral type earlier than F5. Figure 5 (top panel) shows a colour-magnitude diagram from Galactic simulations presenting dwarfs and subgiant stars with $V < 13$ and $T_{\text{eff}} < 7220$ K (i.e. later than F0). At the limit of $(G_{BP} - G_{RP})_0 = 0.56$ (dashed red line) simulations show that we expect a loss of $\sim 10\%$ of good targets, but an inclusion of $\sim 90\%$ of contaminating stars. Extending the selection criteria to bluer colours would increase the number of contaminating stars (i.e. stars with spectral type earlier than F5), with negligible recovery of good targets (bottom panel of Fig. 5). The upper luminosity threshold of FGK subgiants was set to include stars down to $\log g = 3.5$. Such a definition also helps to reduce the contamination from evolved, reddened sources. The lower luminosity threshold of FGK dwarfs is set at about three magnitudes below the main sequence to avoid contamination from potentially spurious sources populating regions of the CMD where no plausible main sequence dwarf is expected to be found. With our selections, distant red giants are excluded by construction independently of their reddening and the same holds for white dwarfs.

In general, theoretical predictions and observational data agree reasonably well. For M-dwarfs, such as those visible in Figs. 1 and 4, simulated stars appear systematically offset toward larger absolute magnitudes than observations, and appear to span a smaller colour range. Some differences are also evident among the observational samples. Several RECONS stars are intrinsically brighter than the Mann et al. (2015) sample stars. The upper absolute luminosity limit of the M sample was set by performing a linear regression of a 10 Myr solar metallicity isochrone from the Padova database (Bressan et al. 2012), represented by the blue dashed line in Fig. 1. Considering the theoretical and observational uncertainties, this threshold is a good compromise to include all cool late-type dwarfs at any metallicity, including binaries and active stars. At the same time, by adopting this limit we are more likely to discard giant stars, which are expected to be brighter than this limit. To reduce the contamination from reddened stars lurking inside our selection region we also limited the distance of the stars in the M sample to 600 pc. The rationale behind this choice is related to our limiting magnitude $V = 16$. Assuming the conversion between Johnson V and *Gaia* magnitudes in Appendix A, the maximum distance at which we can detect an unreddened M0 type star with an apparent visual magnitude $V = 16$ and at the bright limit of our selection ($M_{G,0} = 6.55$) is 573 pc. Therefore, any other M-type star within our selection region cannot be found at a distance of greater than 600 pc. In summary, our adopted selection criteria are defined as follows:

$$\text{M sample} = \begin{cases} (G_{BP} - G_{RP})_0 \geq 1.84 \\ M_{G,0} > 2.334 (G_{BP} - G_{RP})_0 + 2.259 \\ \text{Distance} < 600 \text{ pc} \\ V \leq 16, \end{cases} \quad (2)$$

$$\text{FGK sample} = \begin{cases} 0.56 \leq (G_{BP} - G_{RP})_0 < 1.84 \\ M_{G,0} \leq 4.1 (G_{BP} - G_{RP})_0 + 5.0 \\ M_{G,0} \geq 4.1 (G_{BP} - G_{RP})_0 - 2.2 \\ V \leq 13. \end{cases} \quad (3)$$

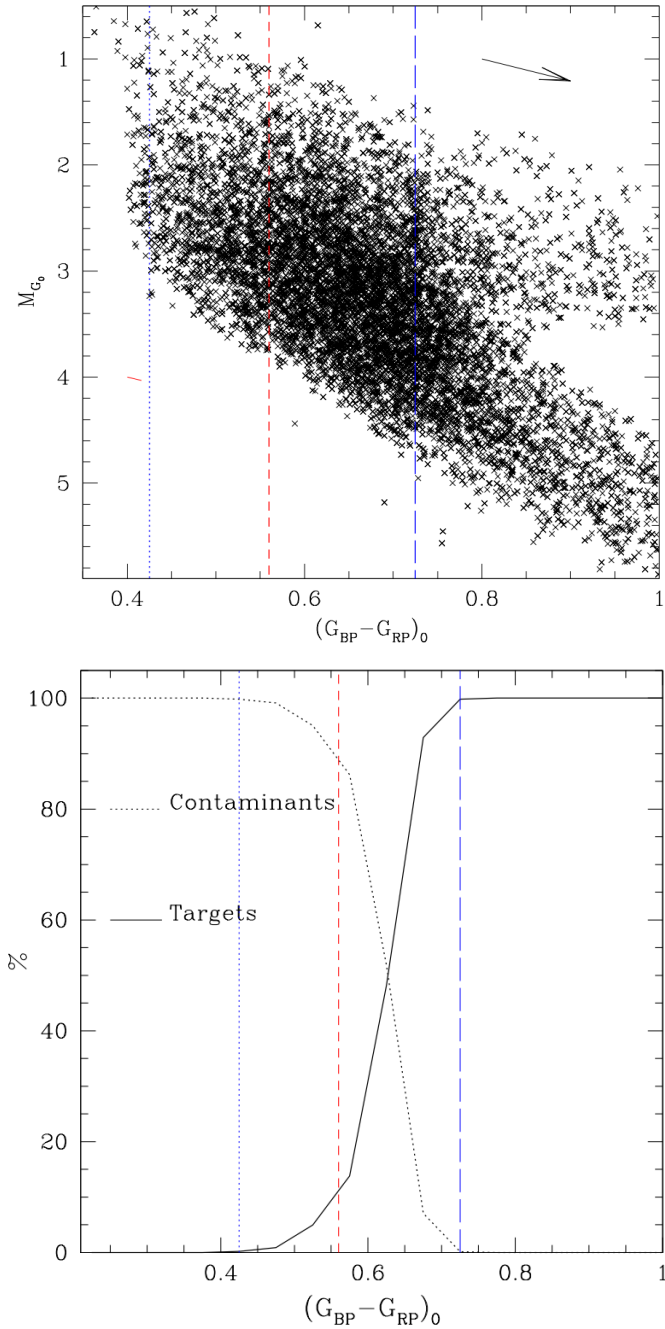


Fig. 5. *Top:* colour–magnitude diagram from Galactic simulations presenting dwarfs and subgiant stars with $V < 13$ and $T_{\text{eff}} < 7220$ K (i.e. later than F0). The arrow indicates the average direction of the reddening vector. *Bottom:* percentage of targets (continuous line) and contaminants (i.e. stars with spectral type earlier than F5, dotted line) within 0.05 mag wide colour intervals. The red dashed vertical lines denote our adopted blue colour limit of the PLATO FGK sample, the blue dotted vertical lines indicate the colour where the selection is 100% dominated by contaminants, while the blue long-dashed vertical lines show the colour where the selection is 100% dominated by the targets.

The quantities that appear in the relations above are the intrinsic colour $(G_{\text{RP}} - G_{\text{BP}})_0$ in the *Gaia* bands, the *Gaia* intrinsic absolute magnitude ($M_{G,0}$), the apparent visual magnitude (V), and the distance from Bailer-Jones et al. (2018). The intrinsic colour is determined by dereddening as discussed in Sect. 4. The selection regions are illustrated in Fig. 1 (red lines) and in Fig. 4 (dashed lines).

By adopting this analytic selection we find that 92% of the stars whose RAVE parameters are compatible with the selection criteria are included. For the M-dwarf samples we find that 97% (190 out of 195 stars) of the Henry et al. (2018) stars, and 99% (152 out of 153 stars) of the Mann et al. (2015) stars are included in the selection.

An alternative estimate of the completeness and contamination can be obtained using theoretical models. We used the same simulations described above in this section to determine the true positive rate (TPR) and the false positive rate (FPR) implied by our selection. We distinguished the cases of M and FGK selected stars. A star is considered a true positive (TP) if it satisfies inequation (2) for the M sample or inequation (3) for the FGK sample and the corresponding theoretical constraints ($V < 16$, $\log g > 3.5$, and $T_{\text{eff}} \leq 3870$ K for M sample and $V < 13$, $\log g > 3.5$, and $3870 \text{ K} < T_{\text{eff}} < 6510$ K for FGK sample); it is a false negative (FN) if it does not satisfy inequation (2) (or inequation (3)) but does satisfy the theoretical constraints; it is a false positive (FP) if inequation (2) (or inequation (3)) is satisfied but the theoretical constraints are not; and finally, it is a true negative (TN) if both inequation (2) (or inequation (3)) and the theoretical constraints are not satisfied. The TPR and FPR are therefore defined as $\text{TPR} = \frac{\text{TP}}{\text{TP} + \text{FN}}$ and $\text{FPR} = \frac{\text{FP}}{\text{FP} + \text{TN}}$. The TPR and the FPR are measures of the performances of the selection algorithm. The TPR of the FGK selection is 100% with a FPR of 12%. The TPR of the M selection is 88% with a FPR of 0%. The FGK selection appears more ‘permissive’ than the M selection, implying higher completeness but also higher contamination. Most of the decrement in the TPR of the M sample appears to be related to the sharp dependence on metallicity of stellar models of late-type stars.

3.4. Impact of selection on the metallicity distribution

It is important to understand if the selection criteria we adopted introduced any appreciable bias on the metallicity distribution of the target stars. To this purpose, we retrieved the Geneva-Copenhagen Survey (GCS) catalogue (Holmberg et al. 2009) and cross-matched it with asPIC by angular separation (accounting for proper motions) and brightness (by comparing the V magnitude of GCS with the V magnitude derived from *Gaia*; see Appendix A). In particular, we considered two samples, one with all stars in the GCS having $T_{\text{eff}} < 6500$ K and for which the value of the metallicity $[\text{Fe}/\text{H}]$ was reported (in total 12 652 stars), and another one with all stars in the GCS for which the asPIC selection criteria (Eqs. (2) and (3)) were satisfied (and had a metallicity estimate; in total 10 798 stars). We then compared the corresponding distributions of metallicities as shown in Fig. 6. We performed a Kolmogorov-Smirnov two-sided test to evaluate whether or not the two distributions were drawn from the same population. We consider a significance level equal to $\alpha = 0.01$ and obtain a p -value equal to 0.6625. The null hypothesis is therefore not rejected and the two distributions are statistically equivalent. This allows us to conclude that our selection criteria are not biasing the distribution of metallicities of the target stars (with respect to the distribution of metallicities of the GCS catalogue).

3.5. asPIC1.1

Figure 7 illustrates the selection in the $M_{G,0}$ versus $(G_{\text{BP}} - G_{\text{RP}})_0$ diagram while Fig. 8 (left panel) presents the distributions of distances for the M and FGK samples. The median values of M and FGK distances are 146 pc and 428 pc, respectively. The

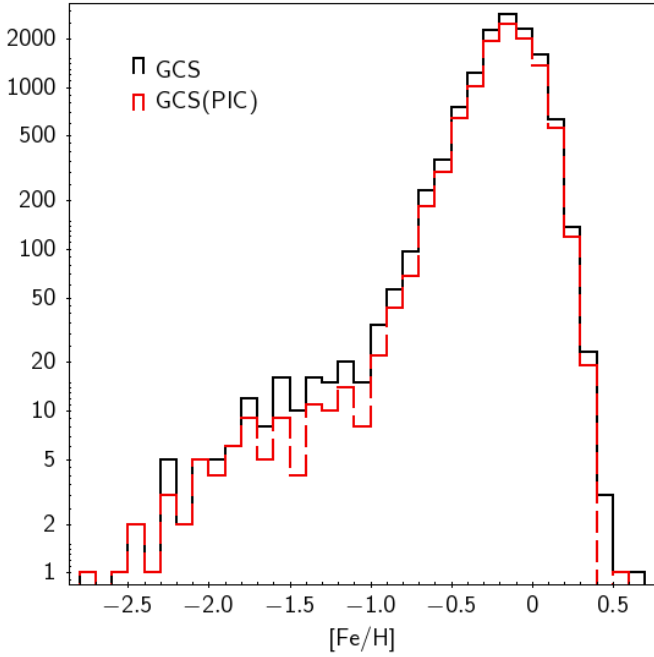


Fig. 6. Metallicity distribution of stars selected from the GCS catalogue (Holmberg et al. 2009). The black histogram shows all stars from the catalogue (with $T_{\text{eff}} < 6500$ K) while the red-dashed histogram shows all stars selected according to asPIC criteria (inequations (2) and (3)).

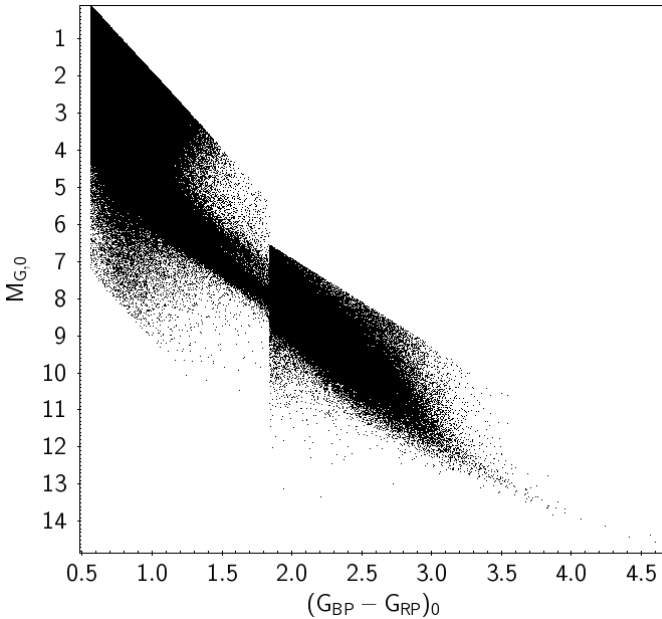


Fig. 7. Combined selection of FGK and M stars in asPIC1.1.

right panel of Fig. 8 shows the distribution of relative distance errors for the different stellar samples. The error on the distance is calculated from the upper and lower distance limits reported by Bailer-Jones et al. (2018) and is equal to the semi-difference of these values. In terms of percentage errors in distance, the median values for the M and FGK stars are 0.6 and 1.6%, respectively. Considering the median distances of the stellar samples reported above, this means that M and FGK distances are set with a precision of 0.9 pc and 6.8 pc, respectively. Figure 9 shows the distribution of the selected FGK dwarf and subgiant stars across the celestial sphere in a Galactic coordinate reference

system and the position of the two currently provisionally defined LOP fields (Nascimbeni et al. 2016). These fields are centred at Galactic coordinates $l = 65^\circ$, $b = 30^\circ$ (RA = $17^{\text{h}}:40^{\text{m}}:19.265^{\text{s}}$; Dec = $+39^\circ:35':01.47''$) for the NPF and $l = 253^\circ$, $b = -30^\circ$ for the SPF (RA = $05^{\text{h}}:47^{\text{m}}:11.700^{\text{s}}$; Dec = $-46^\circ:23':45.37''$). The definitive choice of the LOP fields, to be frozen at the latest two years before launch, involves a complex optimisation task to merge several constraints and priorities of both engineering and scientific nature and will be discussed in Nascimbeni et al. (in prep.). Nevertheless, the two provisional fields illustrated in Fig. 9 already satisfy all PLATO requirements (Sect. 2).

The final all-sky catalogue is named asPIC1.1 and contains 2 675 539 stars. There are 2 378 177 FGK dwarfs and subgiants ($V \leq 13$) and 297 362 M-dwarfs ($V \leq 16$). In Appendix B, we describe an alternative selection of FGK dwarf and subgiant stars based on Johnson B and V photometry, which was used to construct the first version of this catalogue (asPIC1.0). In Appendix C, we describe the catalogue content, while in Appendix D, we present the details of the implementation procedure that has been adopted to construct asPIC1.1.

4. Reddening and absorption

Because of the presence of interstellar matter, stars appear redder and fainter than they are. This has an impact on the selection of PLATO targets which is based on the use of photometric quantities like colours and magnitudes, as discussed in Sect. 3. In particular, as the PLATO samples are magnitude limited, F-type stars, being intrinsically more luminous are located at larger distances on average than later type stars, and are therefore the most affected by reddening. By neglecting the reddening and absorption correction we estimate that 100% of the stars at the blue limit of our colour selection ($G_{\text{RP}} - G_{\text{BP}} = 0.56$, Eq. (3)) would be contaminant dwarfs and subgiants (i.e. hotter than spectral type F5). It is therefore important to account for the impact of interstellar matter before selecting the target stars.

Different approaches can be considered to determine reddening and extinction. One of the drivers for the choice of the best approach to follow for the PLATO input catalogue is to consider that the most important targets lie in the solar neighbourhood. Considering FGK targets, reddening should be estimated up to a distance of about 1 kpc.

Figure 10 shows the cumulative distributions of reddening (left) and HIPPARCOS distances (right) for all stars in the Geneva-Copenhagen Survey (GCS)⁷. The figure shows that 90% of these stars are within 150 pc and the reddening is $E(b - y) \leq 0.02$ corresponding to $E(B - V) \leq 0.04$. In the GCS, reddening is estimated from the calibration by Olsen (1988) of the Strömgren photometry intrinsic colour index ($b - y$). The stated precision is 0.009 mag. The same figure also shows that around 20% of the stars have negative reddening, and are too close to allow a reliable reddening determination. At the small distances of PLATO targets, determining the reddening on a purely empirical photometric basis is challenging and requires measurements of very high precision. Such estimates can be obtained from measurement of interstellar neutral sodium absorption lines imprinted on spectra of background stars (e.g., Vergely et al. 2001). Such measurements are best performed on early-type stars and extrapolation techniques can be used to reconstruct tomographic maps of the local interstellar

⁷ We used here the GCSII version (Holmberg et al. 2007) which reports the reddening.

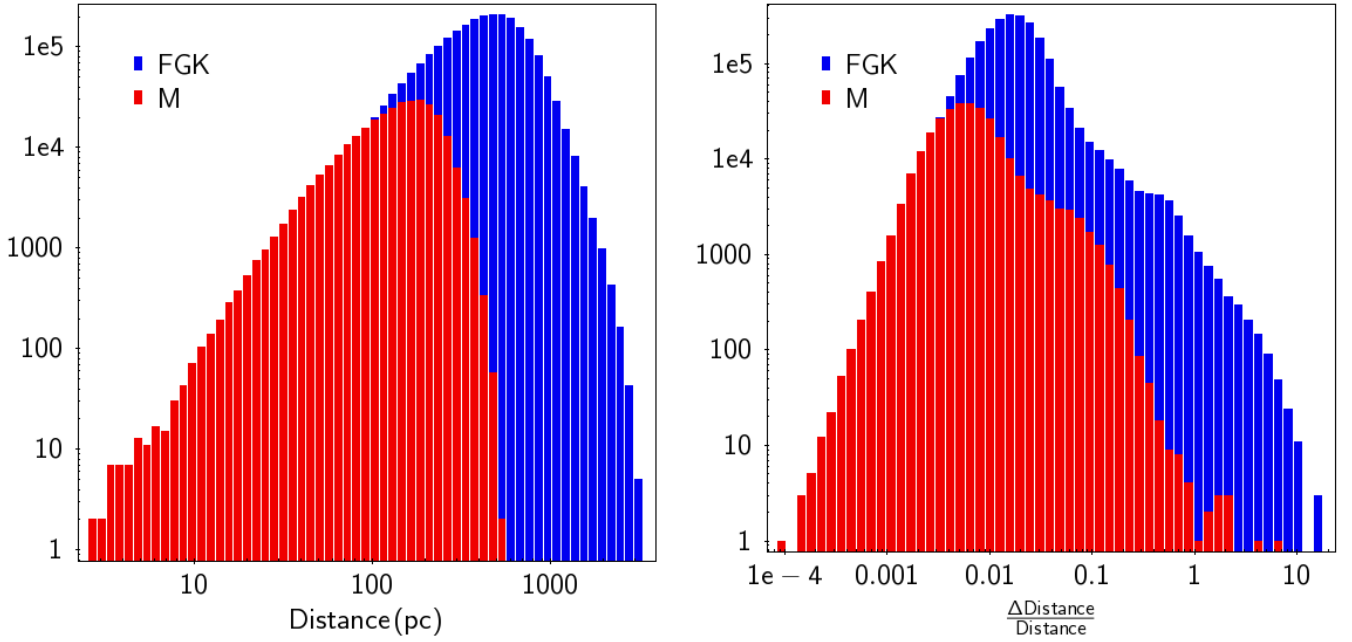


Fig. 8. *Left:* distributions of distances for M (red histogram) and FGK (blue histogram) stars as asPIC1.1. *Right:* distribution of relative errors on distances in asPIC1.1.

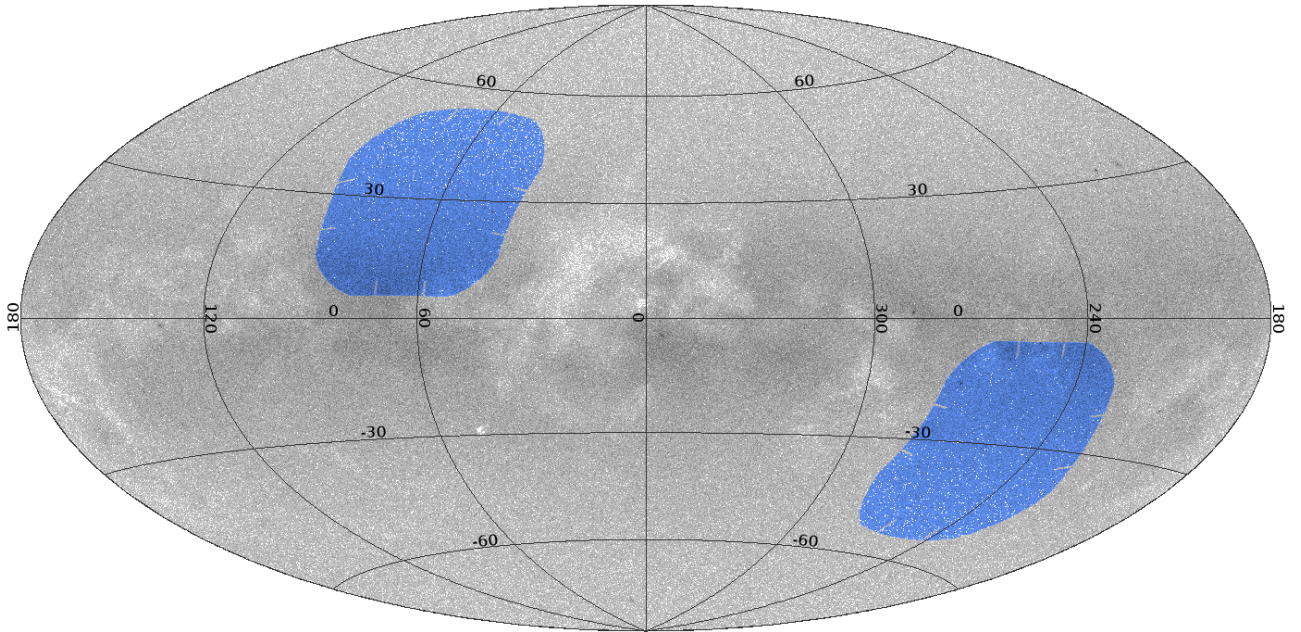


Fig. 9. Distribution of asPIC1.1 dwarf and subgiant FGKM stars across the celestial sphere (grey points) in a Galactic coordinates reference system. The plot also shows the position of the two provisional long-duration fields (NPF and SPF).

medium in all directions. Such spectroscopic measurements have a superior precision than purely photometric estimates, but the empirical database is still limited (Welsh et al. 2010), and the resulting reddening maps have a limited spatial extent (<300 pc from the Sun) and would be largely incomplete for our samples.

In general, spectroscopic measurements can be used to determine the effective temperature and therefore the intrinsic colours of stars, which can then be compared with the observed colours in several optical and near-infrared bands to infer colour excesses (Chen et al. 2019). These methods suffer from the completeness problem mentioned above (although the availability of

large spectroscopic surveys mitigates this problem), from the inhomogeneity of the spectroscopic samples adopted and from the different accuracies and precisions of the effective temperature determinations.

A reasonable compromise among the different techniques can be obtained by coupling empirical determinations with appropriate priors derived from theoretical models or observational constraints. Several authors began to produce 3D reddening maps using optical and/or near-infrared photometry and spectroscopy from large surveys adopting different assumptions on dust and stellar distributions (Green et al. 2018; Gontcharov 2010). Among these, we chose the one recently presented in

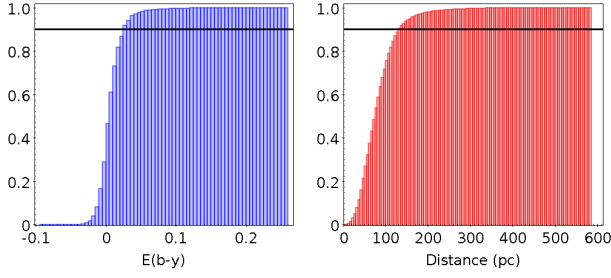


Fig. 10. *Left:* cumulative distribution of reddening values in GCS. *Right:* cumulative distribution of HIPPARCOS distances in GCS. Horizontal lines denotes the 90th percentile of the cumulative distributions.

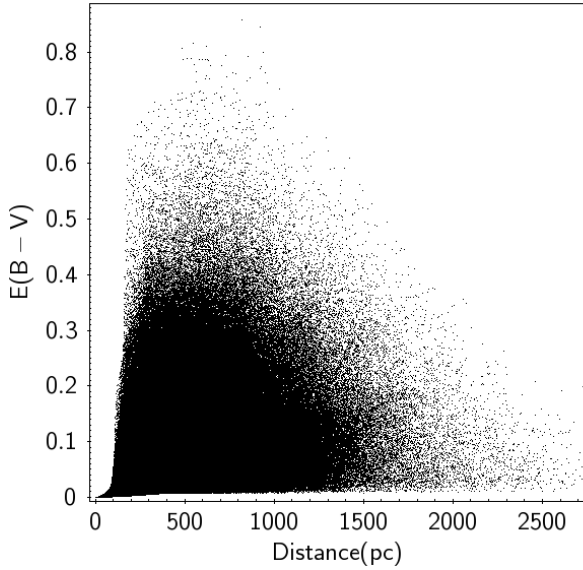


Fig. 11. Reddening $E(B - V)$ vs distance for all stars in asPIC1.1.

Lallement et al. (2018)⁸ because of its accurate description of the local interstellar medium and its overall spatial coverage which allowed us to determine the reddening for a large fraction of PIC stars in a homogeneous way. The 3D reddening map of Lallement et al. (2018) is described in Appendix E. The median reddening of asPIC1.1 stars inside the reddening map is equal to $E(B - V) = 0.04$ and the median uncertainty is $\sigma E(B - V) = 0.02$. For the M sample, 99.8% of the stars are contained in the reddening map and, for the FGK sample, this figure is 81.8%. For the stars falling outside the map, reddening is calculated up to the edge of the map, and then a correction is added, as described in Appendix E.1. Figure 11 presents the $E(B - V)$ versus distance diagram for all stars in asPIC1.1.

5. Stellar parameters

This section describes the algorithms we used to estimate effective temperatures, radii, and masses.

5.1. Intrinsic colour-effective temperature relation

We first determined the relationship between the intrinsic colour in the *Gaia* bands and the effective temperature. For FGK spec-

⁸ Very recently Lallement et al. (2019) published an updated version of the map, which is not included in this version of the catalogue but will be considered in the next version.

tral types, we used the stellar sample from Casagrande et al. (2010). Effective temperatures reported in this catalogue are calculated using the infrared flux method (IRFM). We cross-matched this sample with *Gaia* DR2 and considered only stars closer than 40 pc from the Sun. Casagrande et al. (2010) derived the IRFM effective temperatures for all stars with HIPPARCOS parallaxes and distances ≤ 70 pc from the Sun assuming they are unaffected by reddening. To be consistent with this choice we assumed that the sample of calibration stars we used (all located within 40 pc) have null reddening. The astrometric and photometric quality conditions expressed by Eqs. (1) and (2) in Arenou et al. (2018) were verified for all the selected stars⁹. Our final list comprises 110 stars out of the 423 stars of Casagrande et al. (2010).

For M-dwarfs, we considered the sample of Mann et al. (2015). All the 179 (see Sect. 3.2) stars in the sample are located within 40 pc from the Sun. We also checked that Eqs. (1) and (2) of Arenou et al. (2018) were satisfied, selecting in this way a total of 171 stars. Their effective temperatures were estimated from optical and near-infrared spectra. Finally, for hot stars ($T_{\text{eff}} > 6510$ K), we considered the temperature and colour estimates reported in Pecaut & Mamajek (2013).

As shown in Fig. 12, from the resulting sample of 281 stars we interpolated a single fifth-order polynomial relating the observed *Gaia* colour of the calibration stars and their effective temperature. The coefficients of the fit are reported below, and the RMS of the residuals of the fit is equal to 65 K

$$T_{\text{eff}}(K) = 9453.14 - 6859.40 (G_{\text{BP}} - G_{\text{RP}})_0 + 3542.16 (G_{\text{BP}} - G_{\text{RP}})_0^2 - 1053.09 (G_{\text{BP}} - G_{\text{RP}})_0^3 + 165.635 (G_{\text{BP}} - G_{\text{RP}})_0^4 - 10.5672 (G_{\text{BP}} - G_{\text{RP}})_0^5, \quad (4)$$

where the relation is valid for $0.5 < (G_{\text{BP}} - G_{\text{RP}}) < 5$. Hereafter, we use Eq. (4) as the relation between the intrinsic colour and the effective temperature. Using Eq. (4) we obtained a first guess of the effective temperature of all asPIC stars. These stars are generally located further, at larger distances than the sample of calibration stars, and therefore reddening cannot be neglected.

As reported in Sect. 4, we estimate reddening using the reddening map of Lallement et al. (2018). The customised algorithm described in Sect. E furnishes the interpolated reddening $E(B - V)$ for any 3D star position. As input, it requires the Galactic longitude l and the Galactic latitude b which are taken directly from the input catalogue along with the distance from the Bailer-Jones et al. (2018) catalogue.

5.2. Conversion from $E(B - V)$ to $E(G_{\text{BP}} - G_{\text{RP}})$ and determination of extinction A_G

The $E(B - V)$ reddening is converted to $E(G_{\text{BP}} - G_{\text{RP}})$ and to the extinction A_G in the *Gaia* G-band as follows:

$$A_G = R_G E(B - V), \quad (5)$$

$$A_{\text{BP}} = R_{\text{BP}} E(B - V), \quad (6)$$

$$A_{\text{RP}} = R_{\text{RP}} E(B - V), \quad (7)$$

⁹ With the exception of one source, the bright spectroscopic binary HD112758, for which the astrometric condition expressed by Eq. (1) of Arenou et al. (2018) was not satisfied, but that was kept in the sample because it is known to host a low-stellar-mass companion (Reffert & Quirrenbach 2011).

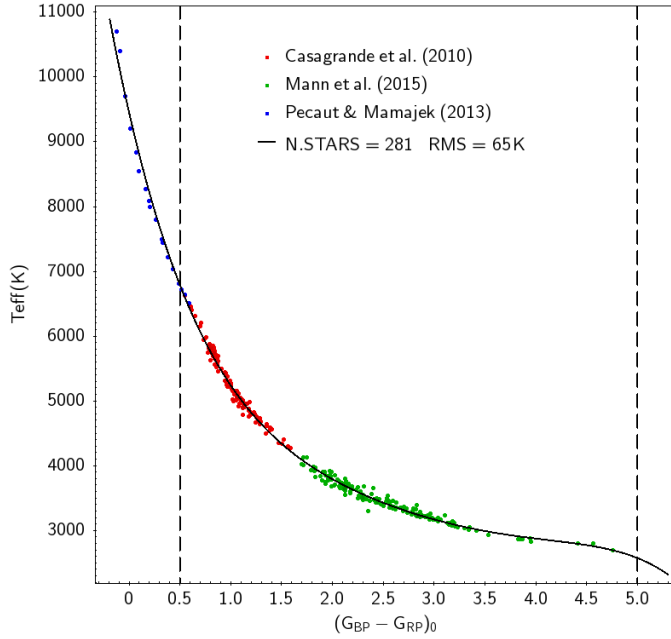


Fig. 12. Relation between effective temperature T_{eff} and intrinsic colour $(G_{\text{BP}} - G_{\text{RP}})_0$. The continuous black line denotes our best-fit model (Eq. (4)), while the dashed vertical lines denote the limits of validity of the relation. Coloured dots represent the samples of Casagrande et al. (2010), Mann et al. (2015), and Pecaut & Mamajek (2013) discussed in the text in Sect. 5.1.

where the *Gaia* extinction coefficients depend on the effective temperature as follows

$$R_G = -0.5335 + 12.9373(T_4) - 13.9514(T_4)^2 - 13.8012(T_4)^3 + 40.9902(T_4)^4 - 23.6648(T_4)^5, \quad (8)$$

$$R_{\text{BP}} = -2.4689 + 59.5802(T_4) - 253.9922(T_4)^2 + 526.5333(T_4)^3 - 523.9970(T_4)^4 + 201.2829(T_4)^5, \quad (9)$$

$$R_{\text{RP}} = 0.0407 + 9.8825(T_4) - 16.8207(T_4)^2 + 9.1283(T_4)^3 + 2.5051(T_4)^4 - 2.4158(T_4)^5, \quad (10)$$

and

$$T_4 = 10^{-4}(T_{\text{eff}}). \quad (11)$$

Such relations were derived using Padova stellar models (Bressan et al. 2012), assuming solar metallicity, solar composition, and temperatures within the interval $2500 \text{ K} < T_{\text{eff}} < 7200 \text{ K}$. In particular, we modelled the ratio of the absorption in a given photometric band to the reddening $E(B - V)$ considering an unreddened 1 Gyr isochrone, and the same isochrone reddened using $A_V = 0.5 \text{ mag}$, where reddening was applied on a star-by-star basis. We then interpolated apparent magnitudes versus effective temperature on the same temperature grid for both the reddened and unreddened isochrones, and calculated the absorption on each *Gaia* band as a function of temperature. In Fig. 13, we show our interpolating relations presented in Eqs. (8)–(10) along with the results obtained using the bolometric correction program of Casagrande & Vandenberg (2018)¹⁰. In particular, we used this program to calculate the bolometric corrections for unreddened and reddened stars interpolating across the grid of $\log g$ and effective temperatures of our

adopted isochrone (assuming solar metallicity and solar composition), and from the bolometric corrections we calculated the ratios between the extinction in each band and the reddening $E(B - V)$ as a function of the effective temperature¹¹. The results we obtained from our procedure agree well with those of Casagrande & Vandenberg (2018) in general. Some small differences are visible, particularly for late-type stars and for the G and G_{BP} bands which may be explained by differences in the underlying model atmospheres, because the Padova group colours and magnitudes are based on the libraries of stellar spectra of Castelli & Kurucz (2003), whereas Casagrande & Vandenberg (2018) used the MARCS library of theoretical stellar fluxes (Gustafsson et al. 2008).

5.3. Correction for reddening and extinction

The observed $G_{\text{BP}} - G_{\text{RP}}$ colour and G magnitude were corrected for reddening and extinction to obtain the intrinsic colour $(G_{\text{BP}} - G_{\text{RP}})_0$ and the intrinsic magnitude G_0 :

$$(G_{\text{BP}} - G_{\text{RP}})_0 = (G_{\text{BP}} - G_{\text{RP}}) - E(G_{\text{BP}} - G_{\text{RP}}), \quad (12)$$

$$G_0 = G - A_G. \quad (13)$$

5.4. Effective temperature: the iteration cycle

We recalculated the effective temperature using Eq. (4) and the new intrinsic colour $(G_{\text{BP}} - G_{\text{RP}})_0$ in Eq. (12). We then used the new effective temperature to estimate the reddening coefficients (Sect. 5.2) and then a new estimate of the intrinsic colour (Sect. 5.3). We then used this to estimate the effective temperature again. A subsequent iteration lead to effective temperatures that differ by less than 10 K from last estimate; therefore two iterations were sufficient.

5.5. Bolometric correction in the G -band, BC_G

The bolometric correction of the G -band is obtained from the relations given in Andrae et al. (2018):

$$BC_G = 6.0 \times 10^{-2} + 6.731 \times 10^{-5} \Delta T_{\text{eff}} - 6.647 \times 10^{-08} \Delta T_{\text{eff}}^2 + 2.859 \times 10^{-11} \Delta T_{\text{eff}}^3 - 7.197 \times 10^{-15} \Delta T_{\text{eff}}^4, \quad (14)$$

for a temperature interval $4000 \text{ K} \leq T_{\text{eff}} \leq 8000 \text{ K}$ and

$$BC_G = 1.749 + 1.977 \times 10^{-3} \Delta T_{\text{eff}} + 3.737 \times 10^{-7} \Delta T_{\text{eff}}^2 - 8.966 \times 10^{-11} \Delta T_{\text{eff}}^3 - 4.183 \times 10^{-14} \Delta T_{\text{eff}}^4, \quad (15)$$

for a temperature interval $3300 \text{ K} \leq T_{\text{eff}} < 4000 \text{ K}$ and where $\Delta T_{\text{eff}} = T_{\text{eff}} - T_{\text{eff},\odot}$ and $T_{\text{eff},\odot} = 5772 \text{ K}$.

5.6. Determination of the absolute magnitude $M_{G,0}$ and of the absolute luminosity L

The intrinsic absolute magnitude $M_{G,0}$ and the luminosity L were calculated as follows:

$$M_{G,0} = G_0 - 5 \log_{10} d + 5, \quad (16)$$

$$\frac{L}{L_{\odot}} = 10^{-0.4(M_{G,0} + BC_G - M_{\text{BOL},\odot})}, \quad (17)$$

where d is the distance of the star in parsecs from Bailer-Jones et al. (2018) and $M_{\text{BOL},\odot} = 4.74$ and $L_{\odot} = 3.828 \times 10^{26} \text{ W}$.

¹⁰ <https://github.com/casaluca/bolometric-corrections>

¹¹ $\frac{A_{\gamma}}{E(B-V)} = \frac{(BC_{A_V=0} - BC_{A_V})}{A_V} \times 3.1$, where $\gamma = G_{\text{BP}}, G, G_{\text{RP}}$ and $A_V = 0.5$.

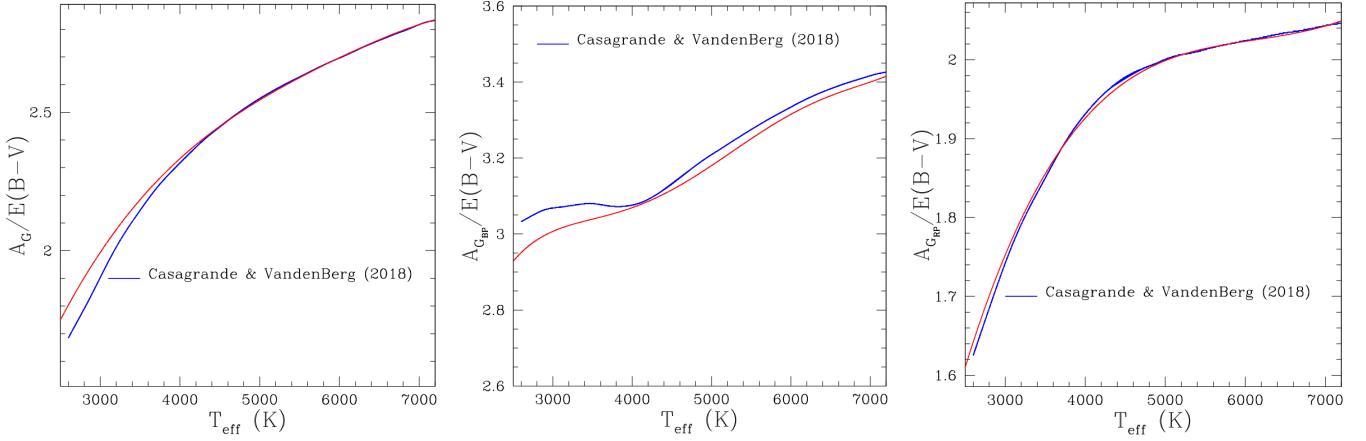


Fig. 13. Relations between $A_G/E(B-V)$ (left), $A_{GBP}/E(B-V)$ (middle), $A_{GBP}/E(B-V)$ (right) and effective temperature T_{eff} derived with the procedure described in the text (Sect. 5.2). The red lines represent our interpolating relations and the blue lines models from Casagrande & Vandenberg (2018).

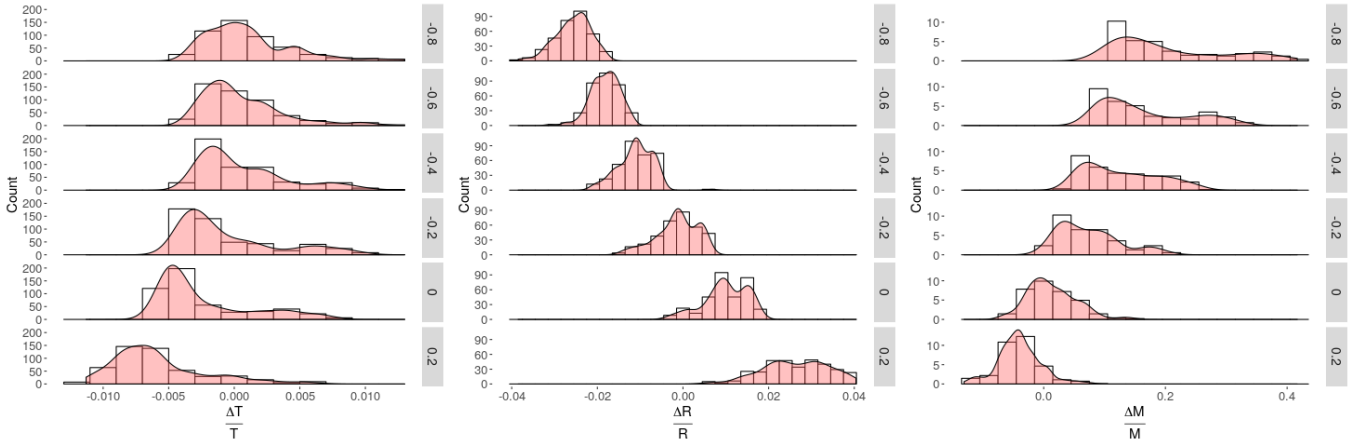


Fig. 14. Relative differences between stellar parameters estimated using asPIC1.1 and theoretical parameters from stellar isochrones in different metallicity intervals. The rightmost labels of each panel denote the centres of the $[M/H]$ metallicity intervals having a width of ± 0.1 dex as also listed in Table 3.

5.7. Determination of the stellar radius

We then calculated the stellar radius R from the Stefan-Boltzmann law:

$$\frac{R}{R_{\odot}} = \left(\frac{T_{\text{eff}}}{T_{\text{eff},\odot}} \right)^{-2} \sqrt{\frac{L}{L_{\odot}}}. \quad (18)$$

The presence of absorption and reddening influences the determination of the stellar radius because it affects both the estimation of the effective temperature and of the absolute luminosity. Therefore it is important to apply dereddening procedures like those we describe above. Reddening and absorption have two opposite and partially compensating effects on the stellar radius. A more detailed discussion about this point is provided in Appendix F.

5.8. Determination of the stellar mass

The stellar mass was determined using the following relation from Moya et al. (2018):

$$\log \left(\frac{M}{M_{\odot}} \right) = -0.119 + 2.14 \times 10^{-5} T_{\text{eff}} + 0.1837 \times \log \left(\frac{L}{L_{\odot}} \right). \quad (19)$$

Table 3. Median relative differences (as a percentage) and standard deviations between asPIC1.1 estimated stellar parameters and theoretical stellar parameters in different metallicity intervals.

$[M/H]$	$\frac{\Delta T}{T} (\%)$	$\frac{\Delta R}{R} (\%)$	$\frac{\Delta M}{M} (\%)$
$-0.9 \leq [M/H] < -0.7$	0.02 ± 0.33	-2.5 ± 0.4	17 ± 9
$-0.7 \leq [M/H] < -0.5$	-0.04 ± 0.30	-1.8 ± 0.3	14 ± 7
$-0.5 \leq [M/H] < -0.3$	-0.06 ± 0.33	-1.1 ± 0.4	11 ± 6
$-0.3 \leq [M/H] < -0.1$	-0.21 ± 0.38	-0.1 ± 0.5	6 ± 5
$-0.1 \leq [M/H] < 0.1$	-0.41 ± 0.36	1.0 ± 0.5	0.5 ± 3.9
$0.1 \leq [M/H] \leq 0.3$	-0.66 ± 0.36	2.6 ± 0.8	-4.1 ± 3.5

The range of validity is $4780 \text{ K} \leq T_{\text{eff}} \leq 10990 \text{ K}$ and $-0.717 \leq \log(L/L_{\odot}) \leq 2.01$. This relationship is valid for dwarfs and subgiant stars.

5.9. Impact of metallicity on the determination of stellar parameters

The procedure used to determine stellar parameters described in Sect. 5 does not take into account metallicity because we do not have metallicity measurements for most of our stars. In this

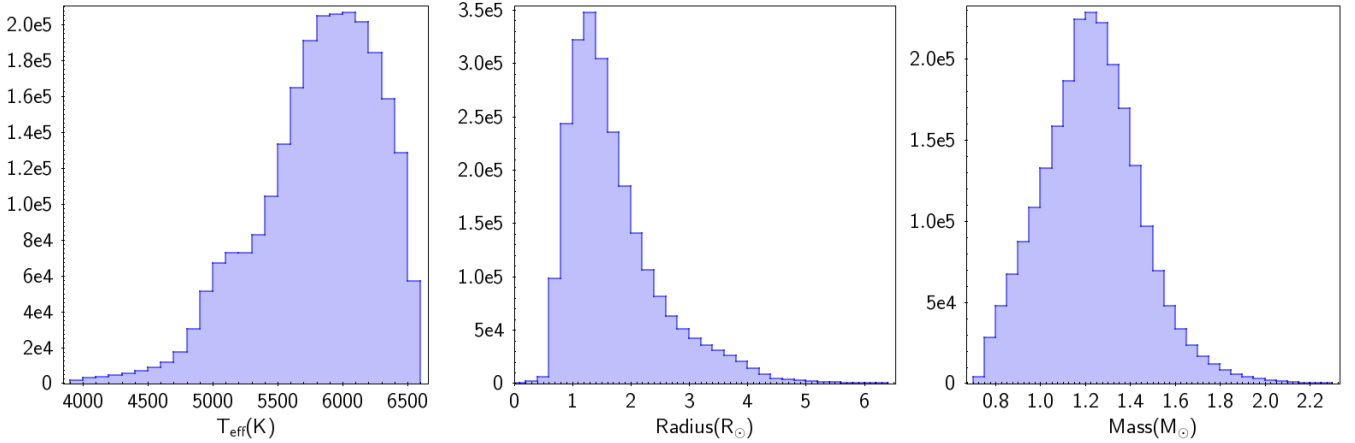


Fig. 15. Distributions of effective temperature (*left panel*), radius (*middle panel*), and mass (*right panel*) for the FGK sample.

section, we estimate the impact that metallicity has on the determination of stellar parameters. We retrieved a grid of isochrones with age between $7 < \log(\text{Age/yr}) < 10$ in steps of 0.2 dex and metallicity between $-1 < [M/H] < 0.3$ in steps of 0.1 dex from the Padova stellar database. For each model in the grid, we determined the stellar parameters with the procedure described in Sect. 5 and compared them with the theoretical parameters reported in the isochrones for six different intervals of metallicity. In particular, we calculated the relative differences ($\frac{\Delta T}{T}$, $\frac{\Delta R}{R}$, $\frac{\Delta M}{M}$) between our parameters ($T_{\text{eff,asPIC1.1}}$, $\text{Radius}_{\text{asPIC1.1}}$, $\text{Mass}_{\text{asPIC1.1}}$) and the theoretical parameters ($T_{\text{eff,ISO}}$, $\text{Radius}_{\text{ISO}}$, Mass_{ISO})

$$\frac{\Delta T}{T} = \frac{T_{\text{eff,asPIC1.1}} - T_{\text{eff,ISO}}}{T_{\text{eff,ISO}}}, \quad (20)$$

$$\frac{\Delta R}{R} = \frac{\text{Radius}_{\text{asPIC1.1}} - \text{Radius}_{\text{ISO}}}{\text{Radius}_{\text{ISO}}}, \quad (21)$$

$$\frac{\Delta M}{M} = \frac{\text{Mass}_{\text{asPIC1.1}} - \text{Mass}_{\text{ISO}}}{\text{Mass}_{\text{ISO}}}. \quad (22)$$

The results are presented in Fig. 14 and in Table 3. We find that the effective temperature has a weak dependence on metallicity with relative differences ranging from -0.6% to 0.02% . The radius presents relative differences of between -2.5% and 2.5% while the mass shows the largest relative differences ranging from -4% to 17% . We note that the mass is also the most uncertain among the theoretical parameters associated to the isochrones. Both the temperature and the mass relative differences are negatively correlated with the metallicity while the radius is positively correlated. We also calculated the standard deviations of the relative differences in each metallicity interval. These differences range between 0.3% and 0.4% for the temperature, between 0.3% and 0.8% for the radius, and between 4% and 9% for the mass.

5.10. Errors estimate

We determined the errors on stellar parameters using Monte Carlo simulations, perturbing all observing quantities according to their associated errors, and assuming a Gaussian perturbation. Beyond colours, magnitudes, distances, and reddening, we also perturbed the bolometric corrections adopting the errors reported by Andrae et al. (2018). The error on the reddening was taken from the reddening map of Lallement et al. (2018) for stars

falling inside its boundaries. For stars falling outside this map, the error on the reddening was assumed equal to twice the error on the median reddening of all stars inside the map, which is equal to $\sigma E(B - V) = 0.04$ (see also Appendix E). The effective temperature deduced by the colour–effective temperature relation was perturbed considering a conservative error of 200 K on its determination. We performed 100 simulations for each star and then calculated the errors on the effective temperature, mass, and radius as the half interval between the 16th and the 84th percentile of the cumulative distribution of the simulated values. The median value of the effective temperature error is 227 K (4%), and this value is $0.1 R_{\odot}$ (7%) for the radius, and $0.1 M_{\odot}$ (8%) for the mass.

5.11. Stellar parameter distributions

In Fig. 15 we present the distributions of effective temperature, radius, and mass for the FGK sample. We note that stellar parameters have not yet been calculated for the M dwarfs sample. They will be included in the next version of the catalogue.

6. Comparisons

We compared stellar parameters and reddening in asPIC1.1 with the values reported in the TIC (Stassun et al. 2019). We used the Candidate Target List v8.01 (CTLv8.01) and cross-matched it with asPIC1.1 using the *Gaia* source ID reported in the CTL catalogue. We considered only those stars for which all stellar parameters and the reddening are defined in both catalogues. We identified 2022 913 stars in common between the two catalogues. The median differences and the standard deviations of the differences between asPIC1.1 and TIC temperatures, radii, masses, and reddenings are: $\Delta T_{\text{eff}} = (-100 \pm 300)$ K, $\Delta R_{\star} = (0.05 \pm 0.07) R_{\odot}$, $\Delta M_{\star} = (0.05 \pm 0.17) M_{\odot}$, and $\Delta E(B - V) = (-0.002 \pm 0.055)$. We note that the distribution of stellar masses appears markedly asymmetric with asPIC having a tendency to overestimate the mass with respect to the TIC. This fact may arise from the different choices adopted for the calibration of stellar masses in the two catalogues. In our case, we adopted the calibration from Moya et al. (2018) which accounts both for the temperature and absolute luminosity of the stars, while Stassun et al. (2018, 2019) adopt a calibration based only on the temperatures. The differences between the asPIC1.1 and TICv8 estimated parameters are shown in the histograms in Fig. 16.

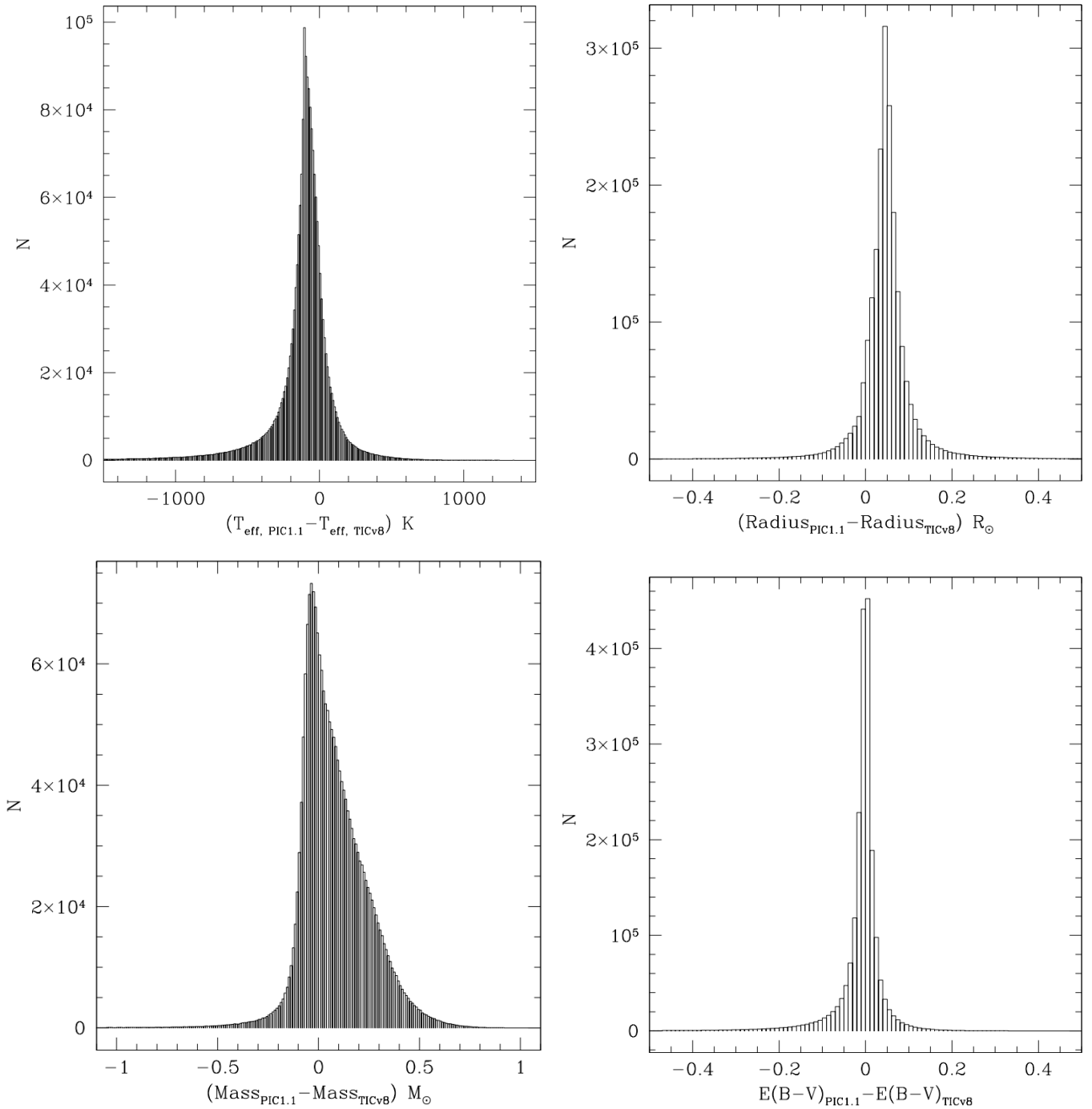


Fig. 16. Distributions of effective temperature, stellar radius, mass, and interstellar reddening differences between asPIC1.1 and CTL.

We then considered the Galactic Archaeology with HERMES (GALAH) survey second data release which contains 342 682 spectroscopically observed stars (Buder et al. 2018). We cross-matched GALAH with asPIC1.1 using the *Gaia* source ID reported in the catalogue. We found 64 061 matched sources. We compared the effective temperatures and the $\log g$ in GALAH with the asPIC1.1 values. In asPIC1.1 we report stellar masses (M) and radii (R). The $\log g$ is estimated using $\log g = \log M - 2 \log R + 4.4374$ (Smalley 2005). The median differences between the asPIC1.1 and the GALAH effective temperatures and gravities are: $\Delta T_{\text{eff}} = (4 \pm 100) \text{ K}$ and $\Delta \log g = (0.1 \pm 0.2) \text{ dex}$, as shown in Fig. 17

We also considered the sample of 1111 FGK dwarf stars from the HARPS GTO program, homogeneously analysed in Adibekyan et al. (2012). We matched the catalogue with asPIC1.1 using angular distances (accounting for proper motions) and found 1102 common sources. The median differences between the asPIC1.1 and the HARPS effective temperatures and gravities are: $\Delta T_{\text{eff}} = (-70 \pm 90) \text{ K}$ and $\Delta \log g = (-0.03 \pm 0.15) \text{ dex}$ (Fig. 18).

Finally, we compared our parameters with the ones reported in the first APOKASC catalogue of spectroscopic and asteroseismic data for dwarfs and subgiants (Serenelli et al. 2017). The cross-match with asPIC1.1 yielded 385 common stars for which all parameters were defined in both catalogues. The analysis

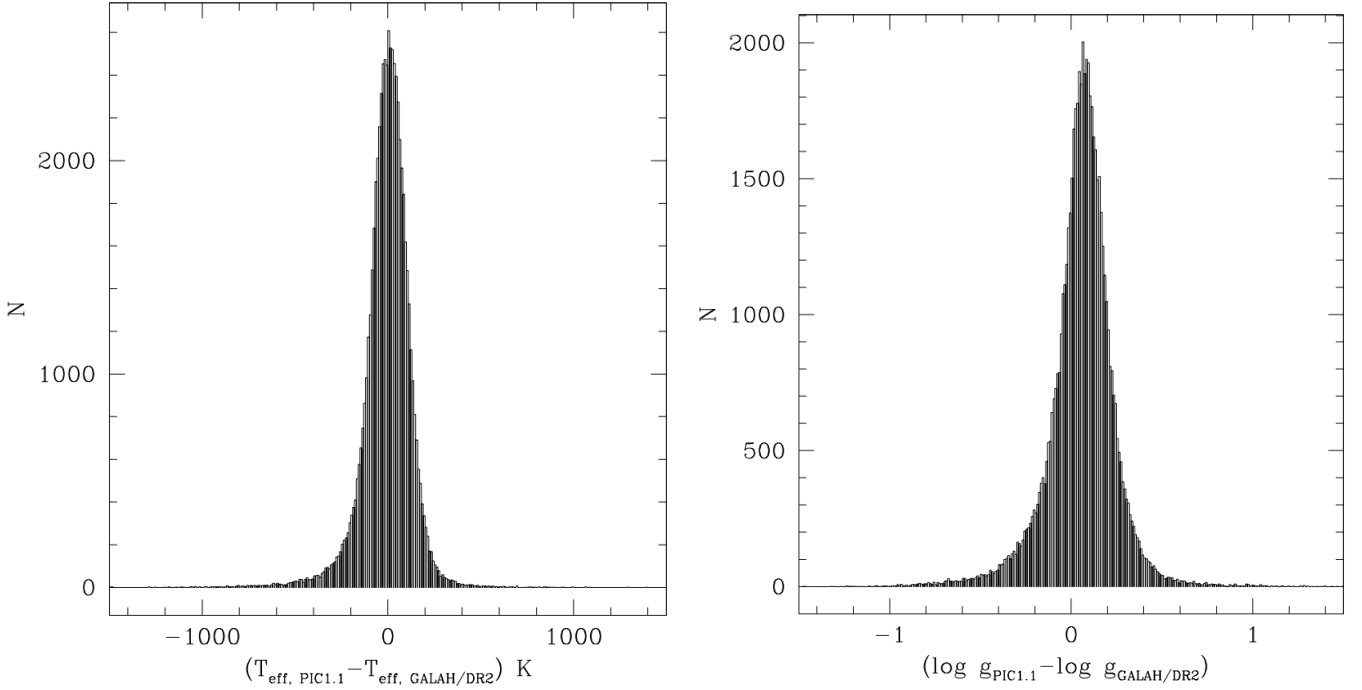


Fig. 17. Distributions of effective temperatures, and stellar gravity differences between asPIC1.1 and GALAH/DR2.

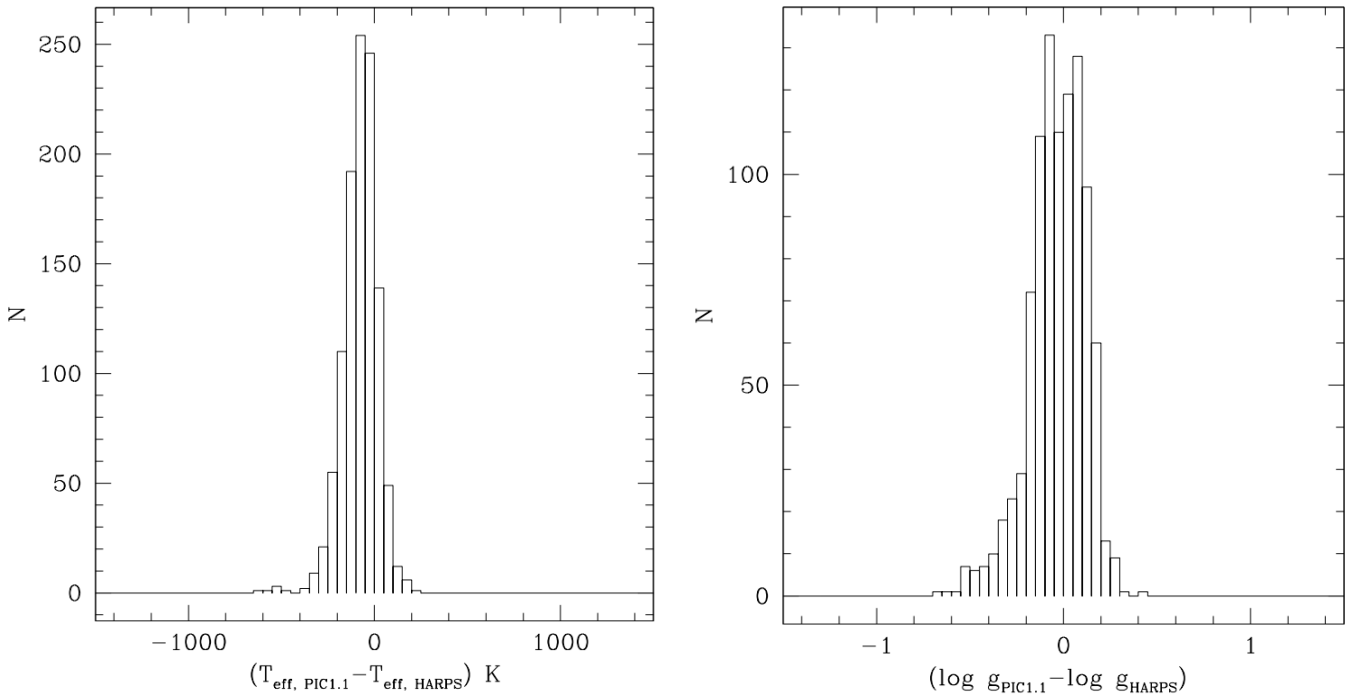


Fig. 18. Distributions of effective temperatures, and stellar gravity differences between asPIC1.1 and the HARPS sample.

was performed both for the SDSS temperature scale and for the spectroscopic temperature scale adopted in the APOGEE Stellar Parameters and Chemical Abundances pipeline (ASPCAP). In the first case, the median differences and the standard deviations of the differences between asPIC1.1 and APOKASC temperatures, radii, masses, and gravities are: $\Delta T_{\text{eff}} = (-130 \pm 90)$ K, $\Delta R_{\star} = (0.1 \pm 0.2) R_{\odot}$, $\Delta M_{\star} = (0.1 \pm 0.1) M_{\odot}$, and $\Delta \log g = (-0.02 \pm 0.06)$, while in the second case: $\Delta T_{\text{eff}} = (90 \pm 130)$ K, $\Delta R_{\star} = (0.1 \pm 0.1) R_{\odot}$, $\Delta M_{\star} = (0.1 \pm 0.1) M_{\odot}$, and $\Delta \log g = (-0.01 \pm 0.06)$ (see also Fig. 19).

Table 4 summarises the results of the comparisons between asPIC1.1 and the stellar catalogues considered in this section and we conclude that the average difference between our effective temperatures and those of the other catalogues is -40 K, the average radius difference is $0.08 R_{\odot}$, and the average mass difference is $0.08 M_{\odot}$. Considering the internal errors we estimated in Sect. 5, and adding them in quadrature to the estimated external errors coming from the comparisons from other catalogues (see above), we conclude that our overall (internal+external) uncertainties on the stellar parameters determination is 230 K (4%) for

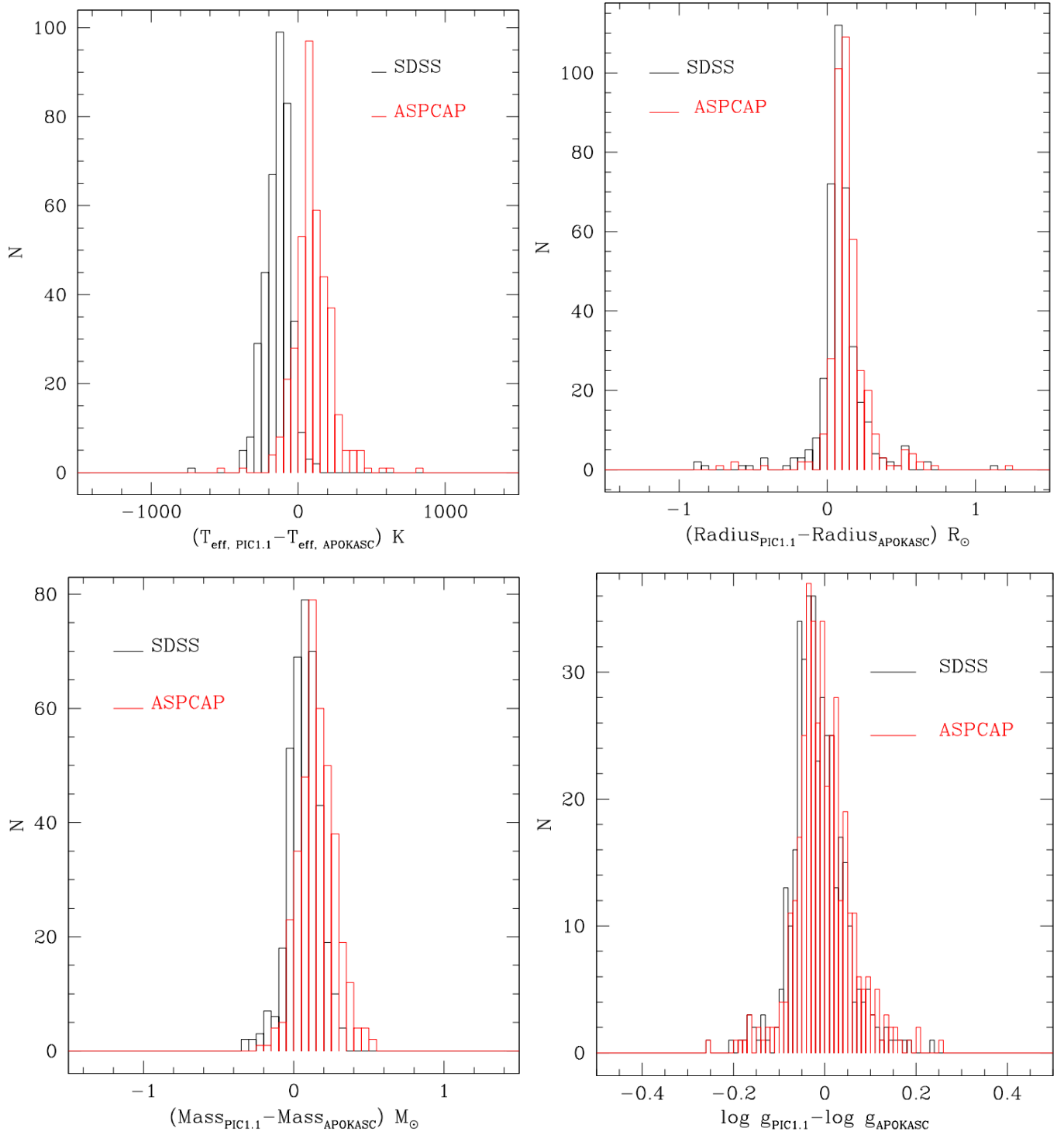


Fig. 19. Distributions of effective temperatures, stellar radii, masses and gravities differences between asPIC1.1 and APOKASC sample and two different temperature scales (SDSS, black and ASPCAP, red).

the effective temperatures, $0.13 R_{\odot}$ (9%) for the stellar radii, and $0.13 M_{\odot}$ (11%) for the stellar masses.

7. Special target list

With the asPIC 1.1 catalogue we also release a special list of objects that consists of all currently known confirmed planet hosts included in the asPIC1.1. This list has been constructed using the Exo-MerCat tool (Alei et al. 2020), which accesses

four main public catalogues of exoplanets, and obtain an updated list of the known planets along with their physical and orbital parameters. In particular, this tool considers the Extrasolar Planets Encyclopaedia¹², the NASA Exoplanet Archive¹³, the Exoplanets Orbit Database¹⁴, and the Open Exoplanet catalogue¹⁵,

¹² <http://exoplanet.eu/>

¹³ <https://exoplanetarchive.ipac.caltech.edu/>

¹⁴ <http://exoplanets.org/>

¹⁵ <http://www.openexoplanetcatalogue.com/>

Table 4. Median differences and standard deviation of the differences between asPIC1.1 and the stellar parameters and interstellar reddening values from other catalogues.

Catalogue	ΔT_{eff} (K)	$\Delta R_{\star}$ ($R_{\odot}$)	$\Delta M_{\star}$ ($M_{\odot}$)	$\Delta \log g$ (dex)	$\Delta E(B - V)$ (mag)
TIC (CTL _{v8})	-100 ± 300	0.05 ± 0.07	0.05 ± 0.17	–	-0.002 ± 0.055
GALAH DR2	4 ± 100	–	–	0.1 ± 0.2	–
HARPS	-70 ± 90	–	–	-0.03 ± 0.15	–
APOKASC (SDSS)	-130 ± 90	0.1 ± 0.2	0.1 ± 0.1	-0.02 ± 0.06	–
APOKASC (ASPCAP)	90 ± 130	0.1 ± 0.1	0.1 ± 0.1	-0.01 ± 0.06	–

each one employing its own cataloguing system. For this reason, before obtaining a uniform merging of the objects of interest, Exo-MerCat performs a proper discrimination among aliases and a standardisation of the different entries provided by the catalogues. The list of the known exoplanets is matched with the asPIC1.1 catalogue.

The special targets list is a living catalogue: it will be continuously updated at any new release of asPIC, because the number of newly discovered planets will constantly increase over the years, not only before but also during the PLATO mission lifetime.

8. Conclusions

In this paper, we present the asPIC1.1 catalogue, a public all-sky catalogue of dwarf and subgiant stars of interest for the PLATO survey, based on the *Gaia* DR2 data release. The asPIC catalogue will be fundamental to identifying the best fields for the PLATO space mission and the most promising targets, analysing the instrumental performances as well as planning and optimising ground-based follow-up studies. This catalogue also represents a valuable resource for the construction of stellar samples optimised for the search of transiting planets.

The catalogue includes a total of 2 675 539 stars among which 2 378 177 are FGK dwarfs and subgiants and 297 362 are M-dwarfs. The median distance of FGK stars in our sample is 428 pc and that for M dwarfs is 146 pc.

We also show that our selection criteria do not bias the statistical distribution of metallicities, and we analyse the impact that metallicity has on the derivation of stellar parameters. We derived the reddening of our targets and developed an algorithm to infer stellar fundamental parameters from astrometric and photometric measurements. We show that the overall (internal+external) uncertainties on the stellar parameters determined by our analysis are ~ 230 K (4%) for the effective temperatures, $\sim 0.1 R_{\odot}$ (9%) for the stellar radii, and $\sim 0.1 M_{\odot}$ (11%) for the stellar masses. We also release a special target list containing all known planet hosts cross-matched with our catalogue.

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Appendix A: Johnson and Gaia DR2 colour transformation

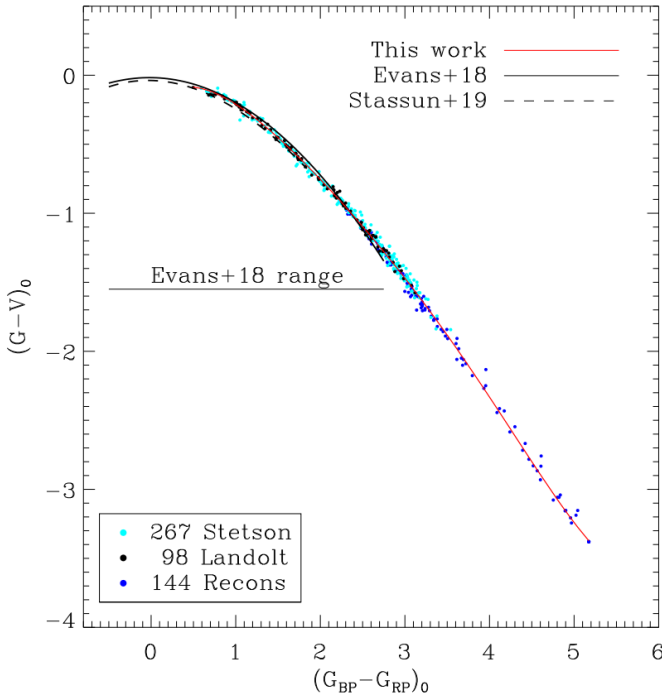


Fig. A.1. Relationship between the intrinsic $(G - V)_0$ and $(G_{BP} - G_{RP})_0$ colours calibrated in this work (solid red line). We also show the corresponding Evans et al. (2018) (solid black line) and Stassun et al. (2019) (dashed black line) relationships.

The definition of PLATO samples requires knowledge of the visual apparent magnitude V . The V magnitude we used to select the targets comes from a transformation of *Gaia* DR2 magnitudes and colours to the Johnson V band. We decided to adopt *Gaia* DR2 colours and magnitudes as the basis of our calculation of the V band magnitude because *Gaia* offers a homogeneous, precise, deep, multi-band, all sky photometry. Evans et al. (2018) give colour–colour relations between several commonly used passbands and *Gaia* DR2 colours. However, the colour range within which these relationships are applicable does not include all the spectral types of interest for PLATO, in particular the M-dwarfs. In order to extend the Evans et al. (2018) relationship to the entire M-type regime, we used a sample of main sequence stars covering a broad range of spectral types for which accurate and precise photometry in the Johnson V is given. To this aim, we used photometry of standard stars given in the catalogues of Stetson (2000) and Landolt (2009). In addition, for M-dwarfs we used the RECONS catalogue (Henry et al. 2018) which provides an accurate census of a volume-limited sample of neighbour stars and is a valuable resource in particular for M dwarfs which are poorly represented in the other catalogues. The Stetson, Landolt, and RECONS catalogues were cross-matched with *Gaia* DR2. We then corrected each star for reddening using the procedure described in Sect. 4, and isolated main sequence stars within 200 pc from the Sun with error on the *Gaia* G -band photometry $\sigma_G < 0.003$, reddening $E(B - V) < 0.05$, and satisfying the astrometric and photometric quality criteria defined by Eqs. (1) and (2) of Arenou et al. (2018). In total we selected 267 stars from Stetson (2000), 98 stars from Landolt (2009) and 144 stars from Henry et al. (2018). We interpolated the $(G - V)_0$ versus $(G_{BP} - G_{RP})_0$ colours with a fifth-order

polynomial relation and the best-fit coefficients are reported in Table 2 of the main text. Figure A.1 shows a graphical representation of the relationship, overlaying the best-fit model (solid red line) and the stellar samples that have been used (coloured dots), the Evans et al. (2018) relation (solid black line), and the Stassun et al. (2019) calibration (dashed black line) used in the TICv8, and based on PHOENIX model atmospheres (provided by K. Stassun, priv. comm.). The standard deviation of the residuals from our best-fit model is $\sigma_{\text{fit}} = 0.02745$ mag. The relation is valid within $0.5 < (G_{BP} - G_{RP})_0 < 5$. We note that this relation is applied to the intrinsic (i.e., reddening-free) magnitudes, and then absorption is added (see Sect. 4) to obtain the apparent magnitude.

Appendix B: Selection of FGK dwarfs and subgiants using Johnson B, V photometry.

For completeness of our discussion we report here an alternative selection of FGK dwarf and subgiant stars used during the construction of the first version of the catalogue (named as PIC1.0). The procedure is similar to the one described in Sect. 3, but is based on the Johnson photometric system. The selection criteria are then defined in the absolute and intrinsic $M_{V,0}$ versus $(B - V)_0$ colour magnitude diagram and are given by the following inequations where we further distinguish dwarfs from subgiants:

Dwarfs:

$$\begin{cases} 0.42 < (B - V)_0 < 1.38 \\ M_{V,0} \geq 5 (B - V)_0 + 0.4 \\ M_{V,0} < 5 (B - V)_0 + 3.5 \\ V < 13 \end{cases}$$

Subgiants:

$$\begin{cases} 0.42 < (B - V)_0 \leq 0.8 \\ M_{V,0} < 5 (B - V)_0 + 0.4 \\ M_{V,0} > 5 (B - V)_0 - 2 \\ V < 13 \end{cases}$$

OR

$$\begin{cases} 0.8 < (B - V)_0 \leq 1.0 \\ M_{V,0} < 4.5 \\ M_{V,0} > 5 (B - V)_0 - 2 \\ V < 13. \end{cases}$$

The selection is depicted in Fig. B.1. One of the limitations of this selection is that it cannot be easily extended in the M-dwarf domain because of saturation effects of the $(B - V)$ colour. With the release of *Gaia* DR2 photometry we decided to work directly in the *Gaia* photometric system, as described in Sect. 3. Nevertheless the relationships reported here can be useful to isolate FGK dwarfs and subgiants using Johnson B and V photometry. The TPR and FPR of this selection are respectively 100% and 45%, and therefore such a selection has the same effectiveness as that presented in Sect. 3.3 in recovering true positives, but implies a larger rate of false positives.

Table B.1. Description of the columns of the asPIC1.1 catalogue.

Col. number	Col. name	Units	Description
1	PICidDR1	–	PLATO identifier for asPIC1.1
2	PICnameDR1	–	PLATO identifier for asPIC1.1 with version number
3	sourceId	–	<i>Gaia</i> DR2 sourceId
4	RAdeg	deg	Barycentric Right Ascension (α) of the source in ICRS at the reference epoch
5	eRAdeg	mas	Standard error $\sigma_{\alpha*} = \sigma_{\alpha} \cos \delta$ of the right ascension in ICRS at the reference epoch
6	DEdeg	deg	Barycentric Declination (δ) of the source in ICRS at the reference epoch
7	eDEdeg	mas	Standard error σ_{δ} of the declination in ICRS at the reference epoch
8	Plx	mas	Parallax
9	ePlx	mas	Standard error on parallax
10	pmRA	mas yr ⁻¹	Proper motion in Right Ascension direction $\mu_{\alpha*} = \mu_{\alpha} \cos \delta$ in ICRS at the reference epoch
11	epmRA	mas yr ⁻¹	Standard error on proper motion in right ascension direction
12	pmDE	mas yr ⁻¹	Proper motion in Declination direction μ_{δ} in ICRS at the reference epoch
13	epmDE	mas yr ⁻¹	Standard error on proper motion in declination direction
14	PM	mas yr ⁻¹	Total proper motion calculated as $PM = \sqrt{pmRA^2 + pmDE^2}$
15	ePM	mas yr ⁻¹	Error on the total proper motion
16	Epoch	yr	Reference epoch to which the astrometric source parameters are referred
17	Gflux	e ⁻ s ⁻¹	Mean flux in the <i>G</i> band
18	eGflux	e ⁻ s ⁻¹	Error on <i>G</i> band mean flux
19	Gmag	mag	Mean magnitude in the <i>G</i> band
20	eGmag	mag	Error on mean magnitude in the <i>G</i> band
21	BPflux	e ⁻ s ⁻¹	Mean flux in the integrated BP band
22	eBPflux	e ⁻ s ⁻¹	Error on mean flux in the integrated BP band
23	BPmag	mag	Mean magnitude in the integrated BP band
24	eBPmag	mag	Error in the mean magnitude in the integrated BP band
25	RPflux	e ⁻ s ⁻¹	Mean flux in the integrated RP band
26	eRPflux	e ⁻ s ⁻¹	Error on mean flux in the integrated BP band
27	RPmag	mag	Mean magnitude in the integrated RP band
28	eRPmag	mag	Error in the mean magnitude in the integrated RP band
29	BPRPexcess	–	BP/RP excess factor
30	BPRP	mag	$G_{BP} - G_{RP}$ colour
31	BPG	mag	$G_{BP} - G$ colour
32	GRP	mag	$G - G_{RP}$ colour
33	GLON	deg	Galactic longitude of the object at the reference epoch
34	GLAT	deg	Galactic latitude of the object at the reference epoch
35	ELON	deg	Ecliptic longitude of the object at the reference epoch
36	ELAT	deg	Ecliptic latitude of the object at the reference epoch
37	rest	pc	The estimated distance (Bailer-Jones et al. 2018)
38	rlo	pc	Lower bound of the confidence interval of the estimated distance (Bailer-Jones et al. 2018)
39	rhi	pc	Higher bound of the confidence interval of the estimated distance (Bailer-Jones et al. 2018)
40	erest	pc	Error on the estimated distance calculated as $restError = 0.5(rhi - rlo)$
41	AG	mag	Estimate of the extinction A_G in the <i>G</i> band
42	eAG	mag	Uncertainty on A_G estimate
43	EBPRP	mag	Estimate of reddening $E(G_{BP} - G_{RP})$
44	eEBPRP	mag	Uncertainty on reddening $E(G_{BP} - G_{RP})$
45	EBV	mag	Estimate of reddening $E(B - V)$
46	eEBV	mag	Uncertainty on $E(B - V)$ reddening
47	extStatus	–	extStatus value. If extStatus=1 the star is outside the reddening map of Lallement et al. (2018)
48	BPRP0	mag	Dereddened ($G_{BP} - G_{RP}$) colour
49	eBPRP0	mag	Error on dereddened ($G_{BP} - G_{RP}$) colour (BPRP0)
50	gaiaV0	mag	Dereddened apparent visual magnitude <i>V</i> in the Johnson-Cousin system
51	egaiaV0	mag	Error on intrinsic visual magnitude (gaiaV0)
52	BJgaiaMV0	mag	Absolute intrinsic visual magnitude obtained as: $BJgaiaMV0 = gaiaV0 - 5 \log_{10} rest + 5$
53	eBJgaiaMV0	mag	Error on absolute intrinsic visual magnitude (BJgaiaMV0)
54	gaiaV	mag	Visual apparent magnitude in the Johnson-Cousin system obtained as: $gaiaV = gaiaV0 + 3.1 EBV$
55	egaiaV	mag	Error on apparent visual magnitude <i>V</i> (gaiaV)
56	BJgaiaMG0	mag	Absolute intrinsic <i>G</i> magnitude obtained as: $BJgaiaMG0 = Gmag - 5 \log_{10} rest + 5 - AG$
57	eBJgaiaMG0	mag	Error on intrinsic absolute <i>G</i> magnitude (BJgaiaMG0)
58	Teff	K	Estimate of stellar effective temperature (T_{eff})
59	eTeff	K	Uncertainty on T_{eff} estimate
60	Radius	$R_{\odot}$	Estimate of stellar radius
61	eRadius	$R_{\odot}$	Uncertainty on stellar radius estimate
62	Mass	$M_{\odot}$	Estimate of stellar mass
63	eMass	$M_{\odot}$	Uncertainty on mass estimate
64	sourceFlag	–	Bitmask with the following meaning: sourceFlag = 1 – FGK star sourceFlag = 2 – M star sourceFlag = 4 – Planet Host
65	contGaiaMag60	mag	Total <i>G</i> magnitude of the <i>Gaia</i> contaminants within 60 arcsec
66	contNumber60	–	Total number of <i>Gaia</i> contaminants within 60 arcsec
67	contGaiaMag45	mag	Total <i>G</i> magnitude of the <i>Gaia</i> contaminants within 45 arcsec
68	contNumber45	–	Total number of <i>Gaia</i> contaminants within 45 arcsec
69	contGaiaMag30	mag	Total <i>G</i> magnitude of the <i>Gaia</i> contaminants within 30 arcsec
70	contNumber30	–	Total number of <i>Gaia</i> contaminants within 30 arcsec

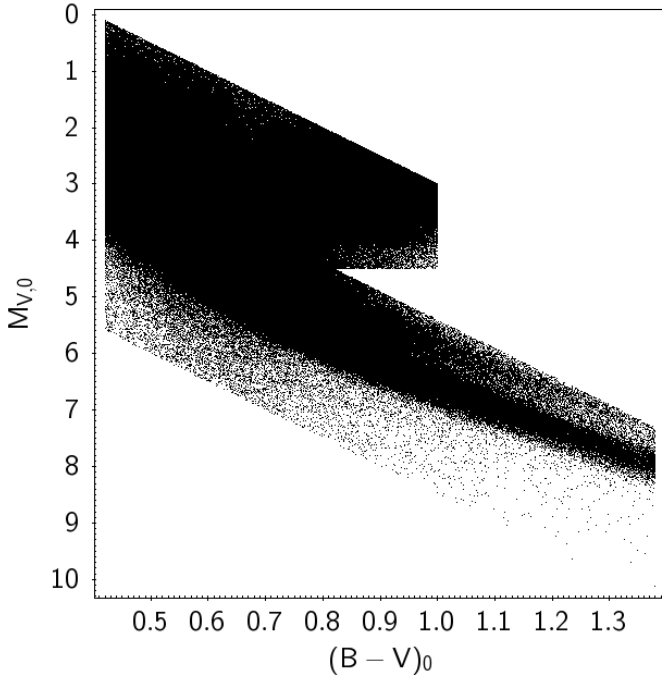


Fig. B.1. Selection of FGK dwarf and subgiant stars in the Johnson B , V absolute intrinsic colour-magnitude diagram.

Appendix C: The catalogue

C.1. Naming convention

The asPIC filename (asPIC.n1.n2) contains two dot-separated integer numbers which are used to distinguish the different versions of the catalogue.

The first number (n1) indicates an asPIC major release and is linked to the *Gaia* data release used to build the catalogue. When a new *Gaia* release triggers an asPIC major update, n1 is increased by one. The second number (n2) indicates a substantial update, which is not major and thus does not require changing n1. Examples of substantial updates are, for example, a change in (a) the number of objects, (b) the number of columns, (c) the format of one or more columns, or (d) the content of one or more columns. When the first number n1 changes, n2 is set to zero.

C.2. Content

In Table B.1 we report the asPIC1.1 column names and a brief description of their meaning.

Appendix D: Implementation procedure

In this Appendix we describe the implementation procedure adopted to build the asPIC1.1 catalogue. The implementation of PLATO asPIC1.1 is done using the PLATO Data Processing Centre - ASI (PDPC-A) infrastructure dedicated to the *Gaia* and PLATO missions.

- The procedure starts using the *Gaia* DR2 catalogue main table (gaiaSource) and the Gaia estimated distances table from Bailer-Jones et al. (2018). A asPIC main table (asPIC1.1) is created and all sources from gaiaSource where: $\text{photGMeanMag} \leq 15.5$ AND bpRp is not null AND rest is not null are added. The asPIC1.1 table includes 52 557 711 sources.

- The asPIC1.1 table is updated and the following columns are added: pmtotal , pmtotalError , photGmeanMagError , $\text{photBpMeanMagError}$, $\text{photRpMeanMagError}$, restError .
- We created an input file to be fed into the stellar parameters' code (Sect. 5). The input file includes all sources in the asPIC1.1 table. The code is run in parallel. The results of the stellar parameters' code (47 966 591 sources) are then fed back into the database in a dedicated table. The stellar parameters' code output columns are: ag , agError , ebprp , ebprpError , ebv , ebvError , extStatus , teff , teffError , radius , radiusError , mass , massError and dwsGFlag . The flag dwsGFlag discriminates between FGK ($\text{dwsGFlag}=1$), M ($\text{dwsGFlag}=2$) stars and stars that are neither FGK or M ($\text{dwsGFlag}=0$). The main table (asPIC1.1) is updated to include these results.
- The main table is further updated to include bpRp0 , bpRp0Error , BJgaiaMG0 , BJgaiaMG0Error , PICidDR1 , gaiaV0 , gaiaV0Error , gaiaV , gaiaVError and gaiaMV0 , gaiaMV0Error . Only 47 288 759 sources out of 47 966 591 are updated, since a few sources are left out because they do not fall within the Gaia-Johnson calibration limits.
- The definition of the samples in the asPIC110 table is updated according to the final selection criteria:
 1. FGK: ($\text{dwsGFlag}=1$ AND $\text{gaiaV} \leq 13.0$ AND $\text{BpRp0} > 0.56$) which outputs 2 378 177 sources;
 2. M: ($\text{dwsGFlag}=2$ AND $\text{gaiaV} \leq 16.0$ AND ($\text{BJgaiaMG0} > 2.334 * (\text{BpRp0}) + 2.259$) AND $\text{rest} < 600.0$) which outputs 297 362 sources.
- Photometric contaminants are searched for and counted in the *Gaia* DR2 catalogue for all the candidate targets listed in asPIC using a cone search with 30, 45, and 60 arcsec radius respectively. To compute the total G magnitude we summed the entire flux of each contaminant. As the contaminants may fall close to the circle border, the computed total magnitude is over-estimated. This information was used to compute the columns contGaiaMag30 , contNumber30 , contGaiaMag45 , contNumber45 , contGaiaMag60 and contNumber60 .

Appendix E: The 3D reddening map

The 3D reddening map presented in Lallement et al. (2018) is produced using a hierarchical approach merging individual photometric colour excess measurements of stars with colour excesses estimated from diffuse interstellar bands absorption. These measurements are coupled with a non-homogeneous, large-scale prior reddening distribution deduced from massive photometric surveys. The 3D map from Lallement et al. (2018) was retrieved from the dedicated website¹⁶. The map is extended in a region between -2 kpc and 2 kpc from the Sun position in the Heliocentric Galactic X and Y coordinates and between -300 pc and 300 pc in the Galactic Z coordinate. It has a resolution of 5 pc^3 and units are $\text{mag} \cdot \text{pc}^{-1}$. The map is particularly useful at small distances since it reveals how local cavities and cloud complexes are spatially distributed around the Sun.

We used the map of Lallement et al. (2018) in conjunction with Bailer-Jones et al. (2018) distances from *Gaia* DR2 to estimate the amount of interstellar reddening to our target stars. We evaluated the reddening $E(B - V)_i$ on a sample of N , equally spaced, nodal points along each target direction by considering the weighted average of the 27 surrounding voxels reddening (a

¹⁶ <http://stilism.obspm.fr>

$3 \times 3 \times 3$ voxels cube around the nodal point). The reddening was weighted by the distance of the nodal point from each one of the neighbouring voxels' centers while the distance between two adjacent nodal points was 2 pc. We therefore summed up the reddening along the boresight to derive the integrated reddening to the target star

$$E(B - V) = \sum_{i=1}^{i=N} \frac{dE(B - V)_i}{dR} \times \Delta R \text{ (pc)}, \quad (\text{E.1})$$

where

$$\Delta R = 2 \text{ pc}$$

and

$$\frac{dE(B - V)_i}{dR}$$

is the reddening per unit length¹⁷ in the i th voxel. The median reddening of asPIC1.1 stars inside the reddening map is equal to $E(B - V) = 0.04$ and the median uncertainty is $\sigma E(B - V) = 0.02$. For the M sample 99.8% of the stars are contained in the reddening map and, for the FGK sample, 81.8%. For the stars falling outside the map, reddening is calculated up to the edge of the map, and then a correction is added as described in Sect. E.1.

E.1. Reddening map extension

As discussed in the previous section, a small fraction of asPIC stars falls outside the limits of the [Lallement et al. \(2018\)](#) map. For these objects we applied a reddening correction $\Delta E(B - V)$ based on the following model

$$\Delta E(B - V)(l, b, d, d_0) = E(B - V)_\infty(l, b) \frac{\int_{d_0}^d \rho(r) dr}{\int_0^\infty \rho(r) dr}, \quad (\text{E.2})$$

where l and b are the Galactic longitude and latitude, and the integral is calculated along the (l, b) direction from the distance d_0 at the border of the map and the distance d at which a given star is found. The function ρ is a model of the density distribution of the dust. As in [Binney et al. \(2014\)](#), we used

$$\rho(\mathbf{x}) = \exp \left[\frac{R_0 - R}{h_R} - \frac{|z - z_w|}{k_{fl} h_z} \right], \quad (\text{E.3})$$

where $\mathbf{x}$ is the position vector of a star with respect to the Galactic centre, R the galactocentric distance along the Galactic plane, z the orthogonal distance from the Galactic plane, and $k_{fl}(R)$, $z_w(R)$ describe, respectively, the flaring and warping of the gas disk,

$$k_{fl}(R) = 1 + \gamma_{fl} \min(R_{fl}, R - R_{fl}), \quad (\text{E.4})$$

$$z_w(R, \phi) = \gamma_w \min(R_w, R - R_w) \sin \phi, \quad (\text{E.5})$$

where ϕ is the Galactocentric azimuth that increases in the direction of Galactic rotation with the Sun at $\phi = 0$. Table E.1 reports the parameter values. The reddening at infinity $E(B - V)_\infty$ in any given direction is taken from the [Schlegel et al. \(1998\)](#) map, multiplied by the correction factor reported in [Binney et al. \(2014\)](#):

$$f(E(B - V)_\infty) = 0.6 + 0.2 \left[1 - \tanh \left(\frac{E(B - V)_\infty - 0.15}{0.3} \right) \right]. \quad (\text{E.6})$$

The correction described in this section is applied only if the reddening calculated up to the [Lallement et al. \(2018\)](#) map edge was not larger than the reddening $E(B - V)_\infty$ estimated from the [Schlegel et al. \(1998\)](#) map.

¹⁷ The map reports precisely this quantity.

Table E.1. Parameters of the analytical dust's density distribution.

h_r	h_z	R_{fl}	γ_{fl}	R_w	γ_w
4.2	0.088	$1.12 R_0$	0.0054	8.4	0.18

Notes. Distances are in kiloparsecs.

Appendix F: The impact of reddening and absorption on the determination of the stellar radius

It is of fundamental importance to estimate the impact of reddening and absorption on the determination of the stellar radius, a crucial parameter in the context of transiting planet searches. By applying the Stefan-Boltzmann law it follows that the relative variation of the apparent stellar radius R , estimated neglecting the presence of dust, with respect to the true stellar radius R_0 can be expressed as

$$\frac{\Delta R}{R_0} = \left(\frac{T_{\text{eff},0}}{T_{\text{eff}}} \right)^2 10^{-0.2(A+\Delta BC)} - 1, \quad (\text{F.1})$$

where $T_{\text{eff},0}$ is the true effective temperature, and T_{eff} the estimated temperature of the target star, while A and ΔBC are the absorption and differential bolometric correction in a given photometric band, respectively.

Considering a reddening law, a colour-effective temperature relation, and an effective temperature-bolometric correction relation, it is possible to express the relative radius variation as a function of reddening only. For example, adopting the relations discussed in Sect. 5, valid for the *Gaia* photometric bands, we obtain that, in correspondence to a reddening equal to $E(G_{\text{BP}} - G_{\text{RP}}) = 0.3$ mag, the radius of a solar-type star would be overestimated by about 3% if both reddening and absorption were neglected in the radius calculation, as shown in Fig. F.1 (solid line). The individual contributions resulting from neglecting either the reddening or the absorption are illustrated respectively by the dashed and dotted lines in the same figure. Therefore, in correspondence to the same amount of reddening, $E(G_{\text{BP}} - G_{\text{RP}}) = 0.3$ mag, neglecting the effect of the reddening leads to underestimation of the effective temperature which in turn leads to overestimation of the stellar radius by about 32% (dashed line). On the other hand, neglecting the effect of absorption produces an underestimation of the stellar luminosity, which implies an underestimation of the stellar radius by about 22% (dotted line). Therefore, reddening and absorption have two opposite effects on the apparent stellar radius which are partially compensating each other.

These results hold for a solar-type star with effective temperature $T_{\text{eff}} = 5772$ K. It should be kept in mind that this compensation effect depends on the spectral type of the target source, the slope of the adopted colour-effective temperature relation (and therefore the adopted photometric system), and the reddening law. For example, in Fig. F.2 we show that for an M3V star, neglecting absorption would produce an underestimation of the stellar radius by about 13%, neglecting the reddening would produce an overestimation of the stellar radius by about 11%, while neglecting both of them would underestimate the radius by about 4%. For a F5V star, neglecting absorption would produce an underestimation of the stellar radius by about 20%, neglecting the reddening would produce an overestimation of the stellar radius by about 38%, while neglecting both of them would overestimate the radius by about 11%.

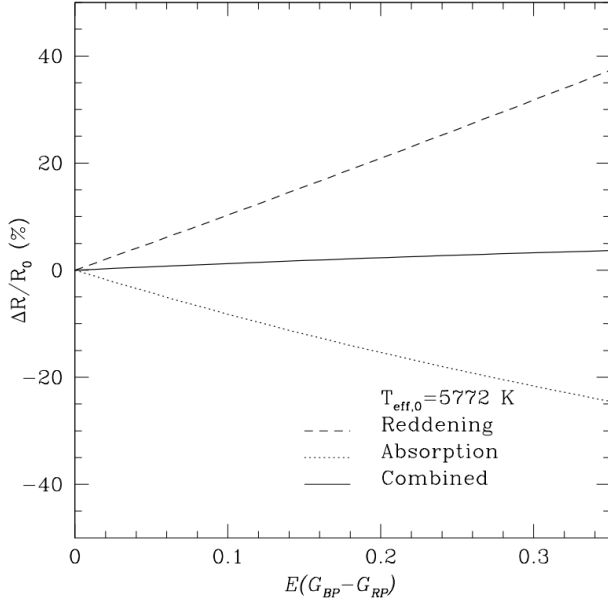


Fig. F.1. Percentage variation of the apparent radius R estimated neglecting the presence both of absorption and reddening, with respect to the true radius R_0 of a solar-type star (with effective temperature $T_{\text{eff}} = 5772$ K) as a function of reddening $E(G_{\text{BP}} - G_{\text{RP}})$ (solid line). The dashed and dotted lines represent the apparent radius variation resulting from neglecting either the reddening (dashed) or the absorption (dotted), respectively.

The presence of interstellar matter has a smaller impact on M-dwarfs given that these sources are intrinsically redder than earlier type stars, and absorption and reddening are more prominent at shorter wavelengths. Moreover, as we move towards

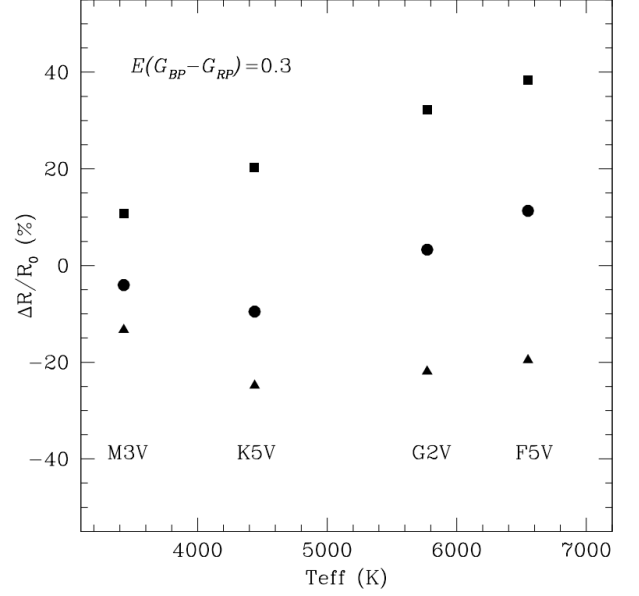


Fig. F.2. Percentage variation of the apparent radius R estimated neglecting the presence of both absorption and reddening (dots) for stars of different spectral types and for a reddening equal to $E(G_{\text{BP}} - G_{\text{RP}}) = 0.3$. Squares and triangles represent the percentage radius variation resulting from neglecting the reddening or the absorption, respectively.

F-type stars, the effect of reddening becomes dominant over the effect of absorption (because of the strong dependence on effective temperature in Eq. (F.1) and the steepening of the colour-effective temperature relation, Sect. 5.1) which generally leads to an overestimation of the stellar radius.